

## Low Mass Meson Spectroscopy and Its Problems

Tokyo Symposium, Feb. 24–26, 2003

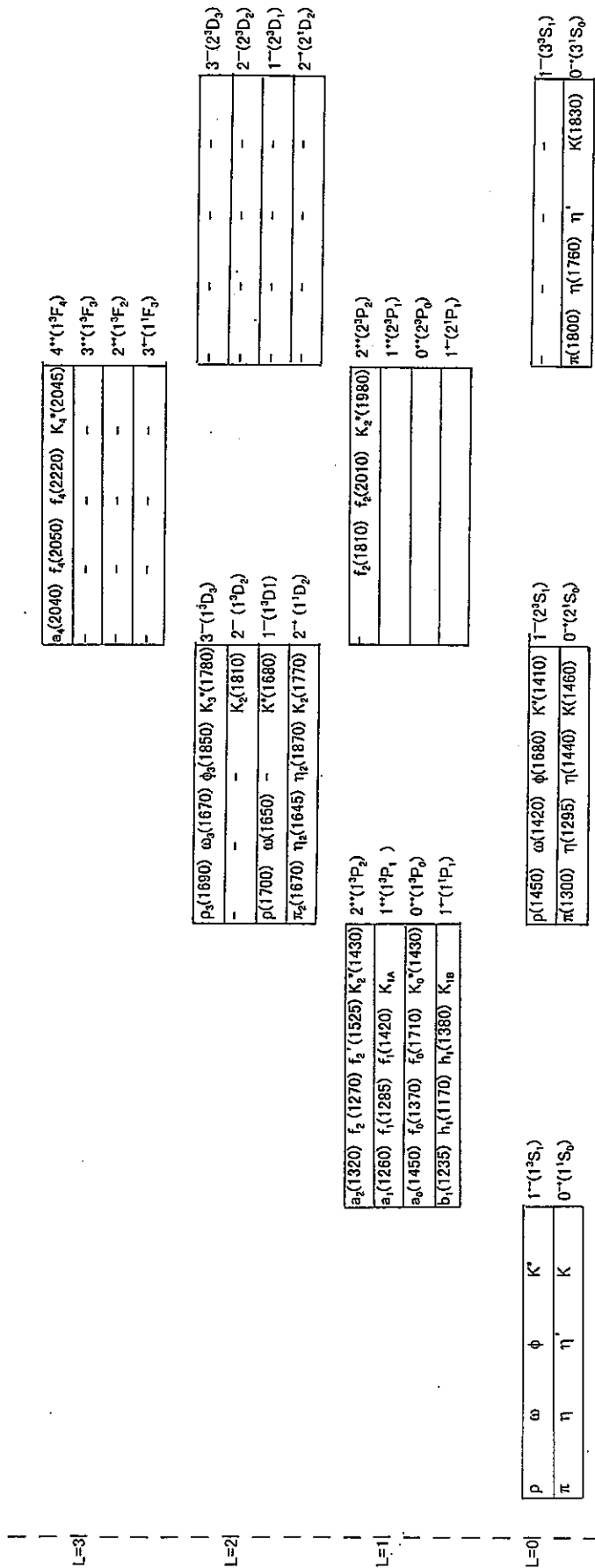
Nihon-U, Ichigaya, Tokyo

Kunio Takamatsu, KEK

### Contents.

1. Introduction
2. SU(3) qqbar nonets and extra states
  - 2.1 Pseudoscalars, radially excited states
  - 2.2 Vectors, radially and orbitally excited states
3. Axial-vector,  $a_1(1260)$
4. Scalars
5. Chiral Particles,  $\sigma(600)$ ,  $\kappa(900)$ ,  $a_1^*(1000)$
6. Summary

图6-7-2 Rosner Box Score



## 1. Introduction

SU(3) qqbar meson states, Non-Relativistic Description

LS coupling scheme, Quantum numbers:  $J^{PC}$ 

$$\mathbf{S} = \mathbf{S}_1 + \mathbf{S}_2, \mathbf{J} = \mathbf{S} + \mathbf{L},$$

$$P = (-1)^{L+1}, C = (-1)^{L+S}$$

$$l = 0 \text{ and } 1, \text{ No } l = 2$$

L=0	$^1S_0$	$0^{-+}$	$\pi$	$2^1S_0 \pi(1300)$
	$^3S_1$	$1^{-}$	$\rho$	$2^3S_1 \rho(1450)$
L=1	$^1P_1$	$1^{+-}$	$b_1(1235)$	
	$^1P_0$	$0^{++}$	$a_0(1450)$	
	$^3P_1$	$1^{++}$	$a_1(1260)$	
	$^3P_2$	$2^{++}$	$a_2(1320)$	
L=2	$^1D_2$	$2^{-+}$	$\pi(1670)$	
	$^3D_1$	$1^{-}$	$\rho(1700)$	
	$^3D_2$	$2^{-}$		
	$^3D_3$	$3^{-}$	$\rho(1690)$	

Rosner: Box Score, L vs  $M^2$ Godfrey and Isgur: SU(3) qqbar meson spectra from  $\pi$  to  $Y(6S)$  (1985)

Four relativised constituent quark masses, QCD potential

SU(3) qqbar so called exotic states,  $0^{+-}, 0^{-+}, 1^{+-}, \dots$ gluonic degrees of freedom glueballs, gg, ggg:  $0^{++}, 0^{-+}, 2^{++}, \dots$ hybrids, qqbarg:  $1^{+-}, \dots$ multiquark states possible  $l=2, 1^{+-}, \dots$ chiral symmetry partners of  $\pi, \rho, \dots, \bar{q}q: 0^{++}, 1^{++}, \dots$ 

Spontaneously breaking of chiral symmetry

Ground and first radial and orbital excited states:

below 1.7 GeV of mass and 0.1-0.3 GeV of width.

$^1S_0$	well established	$2^1S_0$	E/1 (in 80's)
$^3S_1$	well established	$2^3S_1$	$\rho(1250)$ A. Donnachie
$^3P_0$	controversial,		
$^3P_1$	$a_1(1260)$ mass and width		
$^3P_2$	well established,		

## 2. SU(3) qqbar nonets

## 2.1 Radially excited pseudoscalar nonet

$\pi(1300)$ ,  $\eta(1295)$ ,  $\eta(1410)$ ,  $K(1460)$ , and  $\eta(1470)$

decay modes:  $a_0(980)\pi$   $\eta(1410)$  ( $\eta\pi\pi/KK\pi$ )

$K^*(892)K$   $\eta(1470)$  ( $KK\pi$ )

E/ $\tau$  controversy and extra  $\eta$  on data with a desire for glueball

$E(1420)$   $\bar{p}p \rightarrow KK\pi\pi\pi, KK\pi, 0^{++}$ , 81cmHBC at CERN(1967)

$\pi^+p \rightarrow KsK\pi n, K^*K, 1^{++}$ , E(1420), 2mHBC at CERN(1980)

$pp$  collision:  $KK\pi 1^{++}$ , E(1420), Omega(1989)

$\eta(1420)$   $\pi^+p \rightarrow KK\pi n(1985), \eta\pi\pi n(1986), 0^{++}$

Discovery of  $J/\Psi$ , hope for glueball search in radiative decays.  $\gamma$

Radiative decays: gluon rich

$c$

$J/\Psi$

hadrons

Flavor tagging

$\bar{c}$

$J/\Psi \rightarrow \gamma KK\pi, \gamma\eta\pi\pi : \iota, \eta(1410)$  and  $\eta(1470)$

Mark II, a broad peak around 1440MeV  $0^+/1^{++}$

Mark III, resolving into two peaks; 1420MeV and 1490MeV

DM2, also  $4\pi\gamma$ ,

BES,

$\bar{p}p$  XB, OELIX

$\pi p$  CEX:  $\eta\pi\pi/KK\pi: \eta(1295), \eta(1410),$

$K^*(892)K \eta(1410), \eta(1470)$

Lattice calculation, scalar, pseudoscalar glueballs

PS glueball: around 2.3GeV, Scalar glueball: around 1.6GeV

Extra states, chiral symmetry,

S. Ishida and M. Ishida, PTP

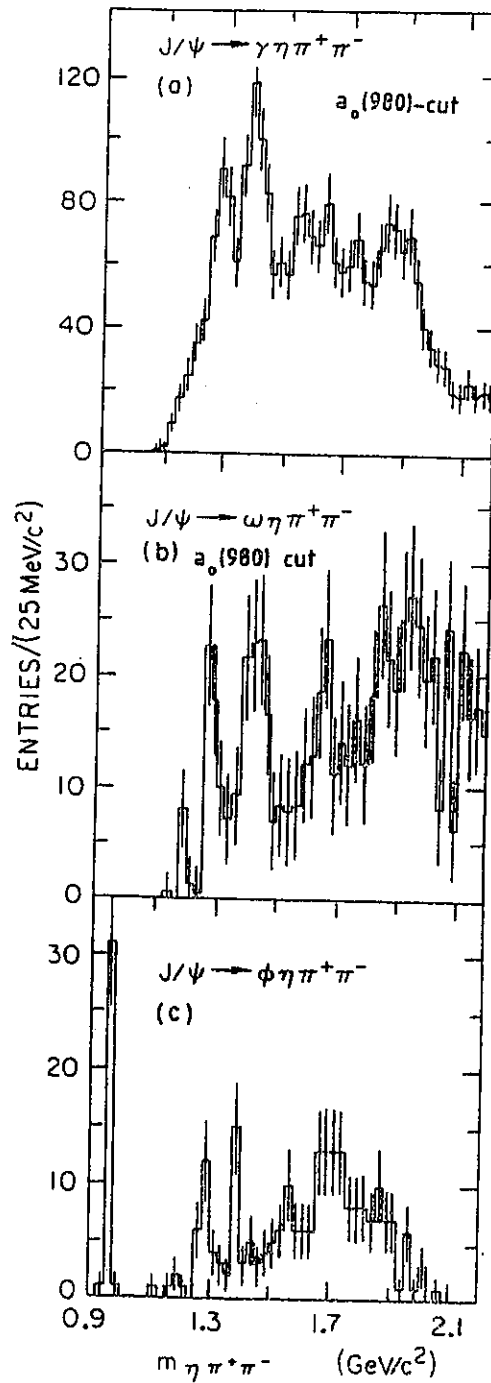
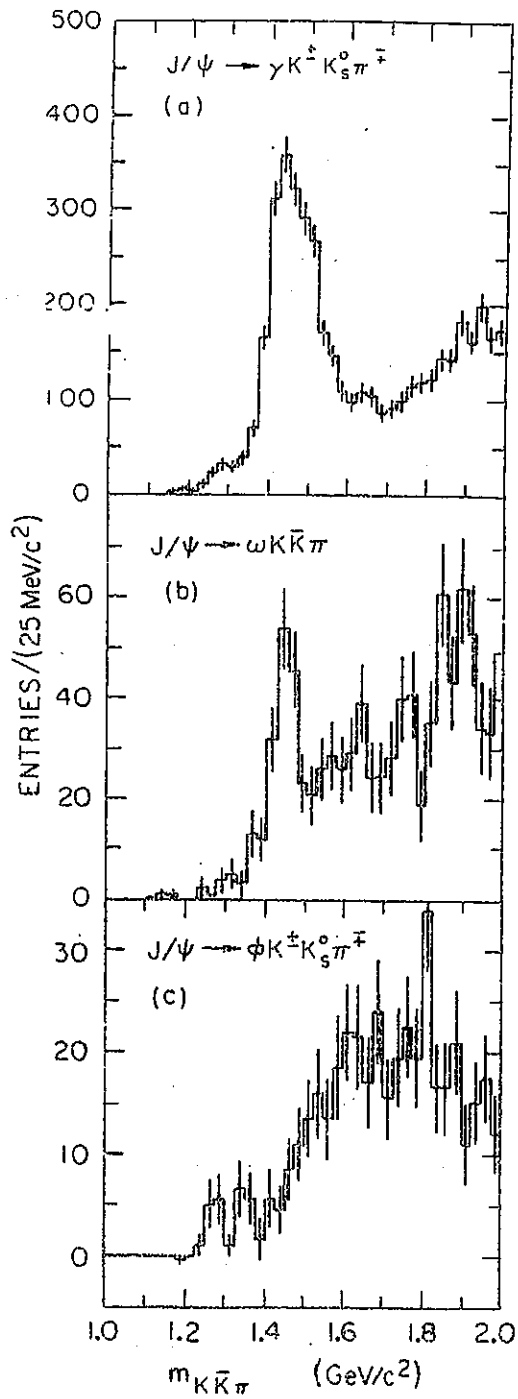
Expect for extra  $\eta$

$U\tilde{(12)}$  classification scheme for light quark mesons

$J/\Psi \rightarrow \gamma K K \pi$  and  $\gamma \eta \pi \pi$ , Mark III (1987)

$KK\pi$

$\eta\pi\pi$



## $J/\Psi \rightarrow \gamma KK\pi$ and $\gamma\eta\pi\pi$ , Mark III (1987)

$KK\pi$

$\eta\pi\pi$

- dominant 'iota' structure

- signals in  $f_1(1285)$ ,  $f_1(1420)$  region
- no 'iota' signal
- further structures between 1.5 and 2.0 GeV

- $f_1(1420)$  signal
- too narrow for 'iota'

- possible  $f_1(1285)$  and  $f_1(1420)$  signals
- no 'iota' signal

- no 'iota' signal
- no  $f_1(1420)$
- events in  $f_1(1285)$  region

- $f_1(1285)$  signal
- no 'iota' signal

## 2.2 Radially and orbitally excited vectors

 $\rho(1450)$ ,  $\rho(1700)$ , and  $\rho(1250)$  $\rho(1450)$  and  $\rho(1700)$ :  $\rho(1600)$  so far.

Comparison of  $\pi\pi$  data between photoproduction,  $\gamma p \rightarrow \pi\pi$ ,  
 and  $e^+e^-$  annihilation,  $e^+e^- \rightarrow \pi\pi$  (1985,1989),  $\rightarrow \omega\pi$  (1986) Donnachie(1987)

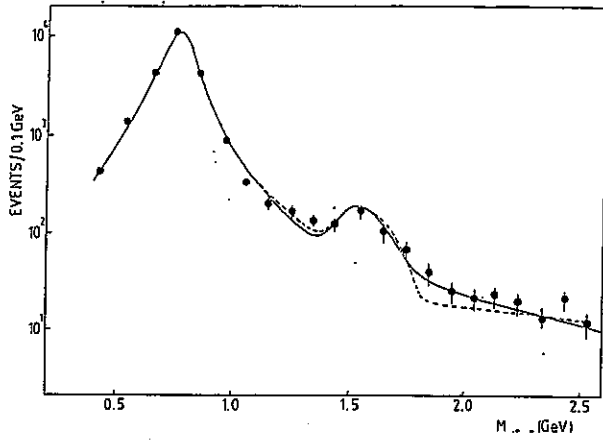
 $\pi p \rightarrow \eta\pi\pi$ ,  $e^+e^- \rightarrow \eta\pi\pi$  (KEK,DM2 1988) $e^+e^- \rightarrow \omega\pi$ ,  $4\pi$  (1991, 1981) $\tau \rightarrow 3\pi\nu$  (1997, 1998) $\bar{p}p$ ,  $\bar{p}n \rightarrow 4\pi$  in  $5\pi$ ,  $\omega\pi$  and  $a_1\pi\rho$ , XB(2001),  $\pi\pi$  OBELIX (1997)

Clegg and Donnachie (1990), Achasov(1997)

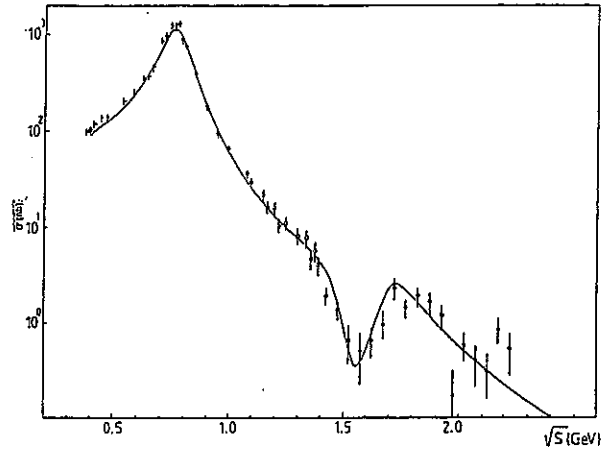
 $\pi^- p \rightarrow \phi\pi^0$ ,  $\rho(1480)$ (1987) $\rho(1250)$ : $\bar{p}p \rightarrow \omega\pi\pi$ ,  $\rho' \rightarrow \omega\pi$  1250MeV (1972)Photoproduction  $\gamma p \rightarrow \omega\pi^0\rho$ ,  $\rho'$  (1974);  $b_1(1235)$  (84, 88) $K^- p \rightarrow \pi\pi\Lambda$ , SLAC(1991),  $\rho'$  at 1300MeV $e^+e^- \rightarrow \pi\pi$  (1983)  $\rho'$  $e^+e^- \rightarrow 4\pi$ ,  $\rho' \rightarrow \rho\pi$  SND<sub>(2001)</sub> 1250MeV $\bar{p}d \rightarrow \omega\pi\pi$ ,  $\rho' \rightarrow \omega\pi$  at 1180MeV XB<sub>(2001)</sub>(( four quark states?:  $\rho(1200)$ ,  $\omega(1200)$ ,  $K^*(1410)$ ,  $C(1480)$  ))



$\gamma p \rightarrow \pi\pi$



$e^+e^- \rightarrow \pi\pi$



$e^+e^- \rightarrow \pi\pi, |F_\pi|^2$ , OLYA at VEPP2, 1985

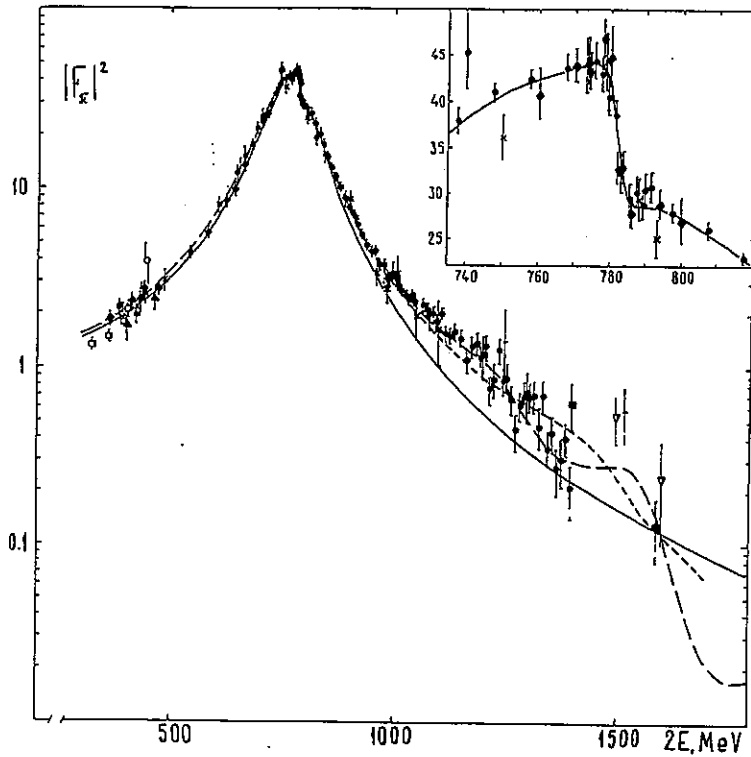


Fig. 7. Experimental values of  $|F_\pi|^2$ :  $\nabla$  [24], — [5],  $\blacksquare$  [6],  $\circ$  [11],  $\blacktriangle$  [20],  $\times$  [23],  $\blacklozenge$  [12],  $\blacktriangledown$  [22],  $\bullet$  [10],  $\square$  [27]. OLYA measurements are from refs. [10,11], CMD from ref. [12]. A solid curve corresponds to the Gounaris-Sakurai formula taking into account  $\rho - \omega$  interference, dotted line - parametrization with  $\omega, \rho, \rho'(1250), \rho''(1600)$  mesons.

$$m_\rho = 1290 \text{ MeV}, \Gamma_\rho = 200 \text{ MeV}$$

3.  ${}^3P_1$  Axialvector  $a_1(1260)$ 

Mass and width parameters of  $a_1(1260)$  in the PDG tables.

Wide spread for mass values in both hadronic production processes and  $\tau$  decays. Extraordinarily wide width,  $\Gamma \sim 600\text{MeV}$  in  $\tau$  decays.

Hadronic productions.

$\pi p$  diffractive process,

$m = 1240\text{MeV}$ ,  $\Gamma = 380\text{MeV}$ . ACCMOR (1980). Corrected for Deck effects.

Charge exchange processes,

$m = 1240\text{MeV}$ ,  $\Gamma = 380\text{MeV}$ . ANL (1981). Corrected for Deck effects.

$m = 1121\text{MeV}$ ,  $\Gamma = 239\text{MeV}$ . KEK (1992).

$m = 1145\text{MeV}$ ,  $\Gamma = 272\text{MeV}$ . Corrected for Deck effects.

Kp backward scattering,

$m = 1041\text{MeV}$ ,  $\Gamma = 230\text{MeV}$ . CERN (1977)

pp centrall. Collision

$m = 1240\text{MeV}$ ,  $\Gamma = 400\text{MeV}$ . CERN (1998).

$\tau$  decays,  $\tau \rightarrow 3\pi\nu$ .

DELCO(1986), MARKII(1986), ARGUS(1986), MAC(1987), ARGUS(1993), OPAL(1997), DELPHI(1998), CLEO(2000).

experimental data: peak value  $= 1000\text{--}1200\text{MeV}$ ,  $\Gamma = 400\text{--}500\text{MeV}$

mass dependent decay amplitude, running mass shift in analyses;

$\rightarrow 1260\text{MeV}$ ,  $\Gamma > \sim 600\text{MeV}$ .

N. Bowler (1986,88), N. Isgur, C.Morningstar and C.Reader(1989), J.M. Kuhn and A. Santamaria (1990), N. Törnqvist (1995).

CLEO(2002),  $m = 1330\text{MeV}$ ,  $\Gamma = 814\text{MeV}$  with  $a_1'(1700)$

two resonant states:  $m(a_1) = 1.11\text{GeV}$ ,  $\Gamma = 0.38\text{GeV}$  and  $m(a_1') = 1.5\text{GeV}$ . J.lizuka (1989).

$q\bar{q}q\bar{q}$ :  $m = 1100\text{MeV}$ . D. Peaslee (1987)

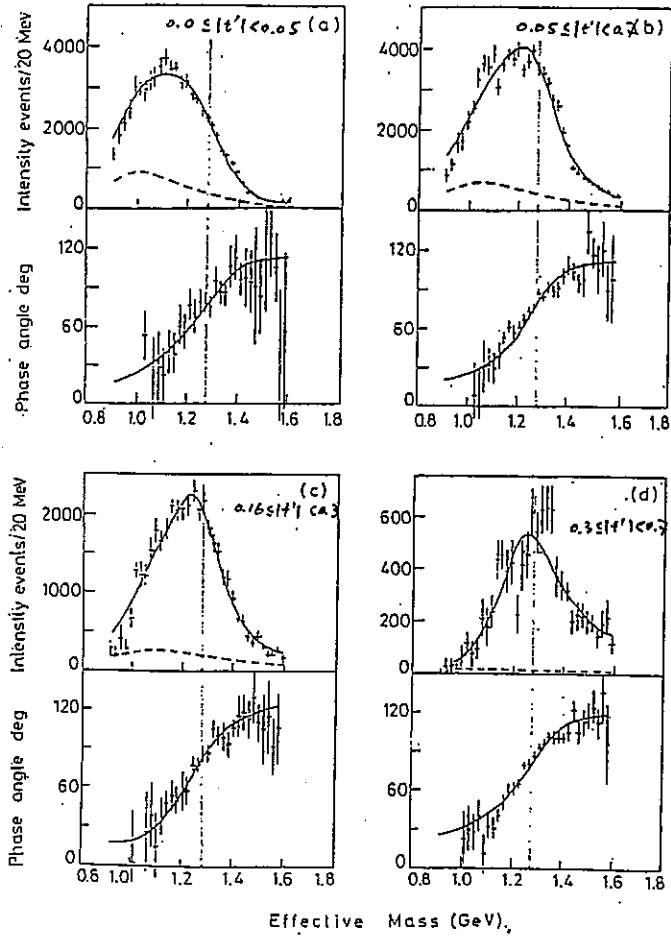
QCD sum rule:  $m = 1.15\text{GeV}$ . L.J.Reinders, S.Yazaki, H.R.Rubinstein (1982)

Lattice QCD:  $m = 1250\text{MeV}$ ,  $f_{a_1} = 0.3(\text{GeV})^2$  M.Wingate et al. (1995)

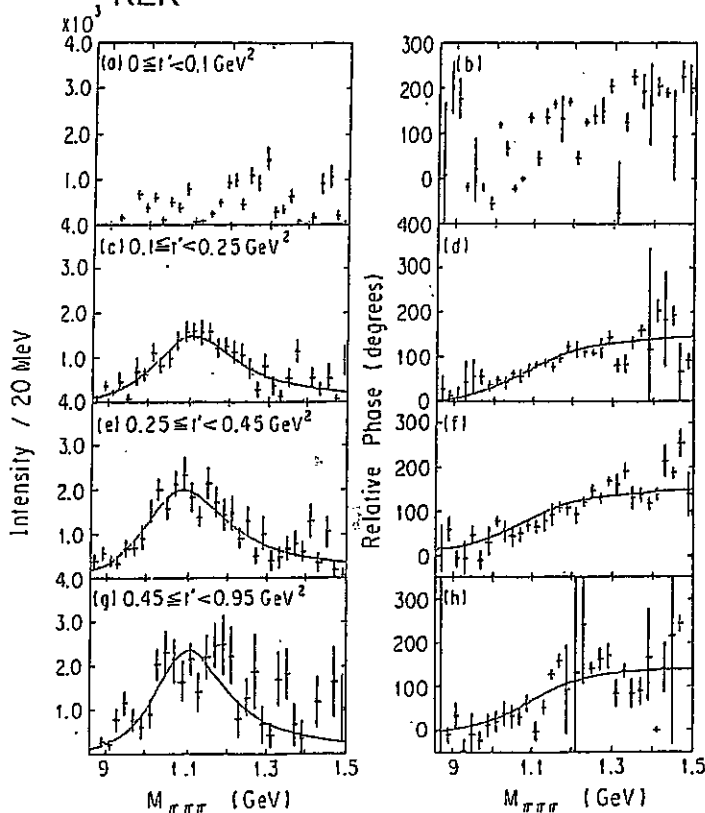
Resonance parameters of  $a_1(1260)$  are, still, uncertain.

$a_1^x$ : chiral partner of  $\rho$  meson: mass around  $1000\text{MeV}$

ACCMOR



KEK



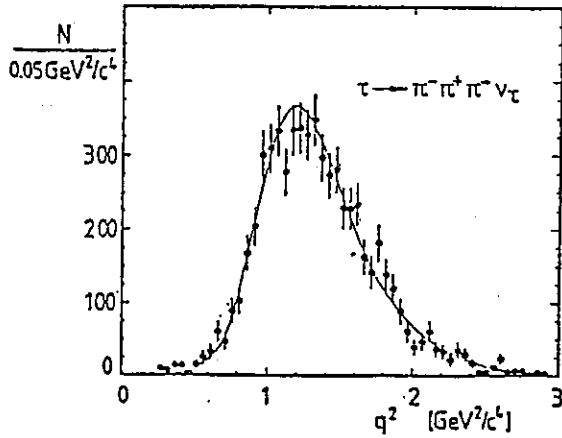
$a_1(1260)$		(MeV)			
$\Gamma$ (GeV/c) <sup>2</sup>	0.10 -0.25	0.25 -0.45	0.45 -0.95	Average	
Breit- M	1128±29	1112±54	1119±73	1122±48	
Wigner $\Gamma$	250±76	290±150	227±179	254±110	
Bowler M	1156±94	1145±42	1114±101	1143±62	
Model 1 $\Gamma$	230±135	281±54	240±163	272±83	
Bowler M	1156±70	1148±46	1115±135	1147±64	
Model 2 $\Gamma$	230±81	290±49	249±167	272±70	

M: mass,  $\Gamma$ : width

Table 2 Parameters of Breit-Wigner fitting to intensity distributions of  $11+\rho S1+$  wave.

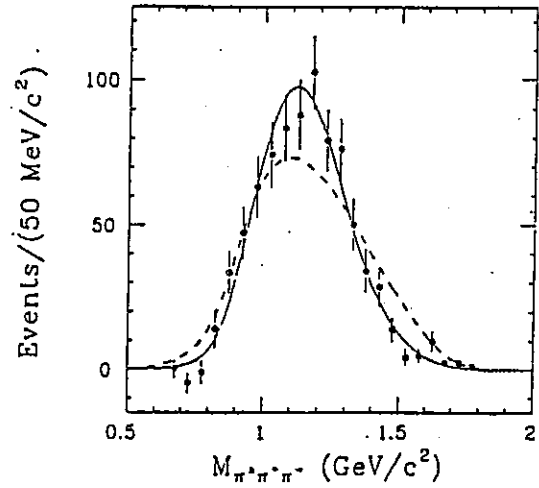
$\tau$  decays,  $\tau \rightarrow 3\pi\nu$

ARGAS(1986)



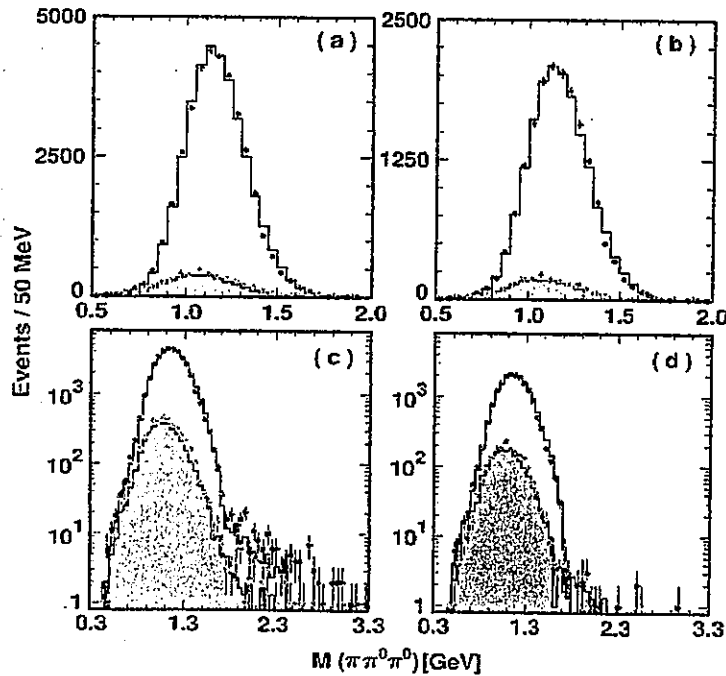
$m=1046\pm 11\text{MeV}$ ,  
 $\Gamma=521\pm 27\text{MeV}$ .

MARK II(1986)



$m=1194\pm 14\pm 10\text{MeV}$ ,  
 $\Gamma=462\pm 56\pm 30\text{MeV}$ .

CLEOII(2002)



Running mass,  $m=1331\pm 10\pm 3\text{MeV}$ ,  $\Gamma=814\pm 36\pm 13\text{MeV}$ .

$$m_{a_1}^2(s) = m_{0a_1}^2 + \delta^2(s), \quad \delta^2(s) = (1/\pi) \int_{s_{\text{min}}}^{\infty} (m_{0a_1} \Gamma_{\text{tot}}^{a_1}(s') / (s-s')) ds'$$

Possible  $a_1'$  contribution,  $m=1330\pm 11\text{MeV}$ ,  $\Gamma=814\pm 38\text{MeV}$ .

Axialvector meson,  $a_1^x$ , a chiral partner of  $\rho$  meson.

- S. Weinberg:  $m(a_1^x) = 2 m(\rho) \sim 1.1\text{GeV}$ ,  $\Gamma(a_1^x) \sim 50\text{MeV}$  ( $a_1^x \rightarrow \sigma\pi$ )  
Current algebra (1990).  
S. Ishida:  $m(a_1^x) \sim 1\text{GeV}$ , COQM (2000).  
M.D. Scadron: LQM and  $a_1^x \rightarrow \sigma\pi$ . (1991).

Hint for  $a_1^x$

$3\pi^0$  in the  $\pi^-p$  charge exchange process. GAMS, CERN (1984)

$\pi^-p \rightarrow \pi^0\pi^0\pi^0n$  at  $100\text{GeV}/c$

$\pi^0\pi^0\pi^0$  scatter plot on the  $2\pi^0$  vs  $3\pi^0$  invariant mass plane.

$2\pi^0$  and  $3\pi^0$  invariant mass spectra.

Rejection of  $K^*(892)$ :  $K_s^0$  cut in  $m_{2\pi}$

Resonances taken into the analysis.

$\eta$ ,  $\eta'$ (958),  $a_1^x$ ,  $a_1$ (1260) and  $\pi_2$ (1670)

$a_1^x \rightarrow \sigma\pi^0$ ;  $\sigma \rightarrow 2\pi^0$ ,

$a_1$ (1260)  $\rightarrow \sigma\pi^0$  and  $f_0$ (980)  $\pi^0$ ;  $\sigma$  and  $f_0$ (980)  $\rightarrow 2\pi^0$ ,

$\pi_2$ (1670)  $\rightarrow \sigma\pi^0$ ,  $f_0$ (980)  $\pi^0$  and  $f_2$ (1270)  $\pi^0$ ;  $\sigma$ ,  $f_0$  and  $f_2 \rightarrow 2\pi^0$ ,

$\eta'$ (1300)  $\rightarrow \sigma\pi^0$  and  $f_0$ (1370)  $\pi^0$ .

Fitting with the interfering BW amplitude method, the Variant Mass and Width (VMW) method.

$$M(s) = \sum_j r_j \exp(i\theta_j) \Delta_j$$

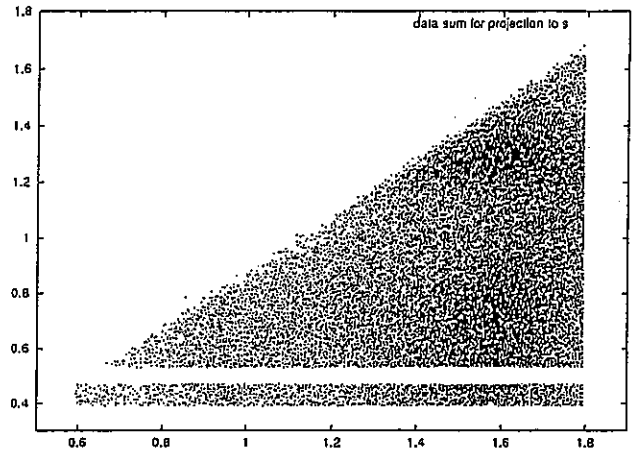
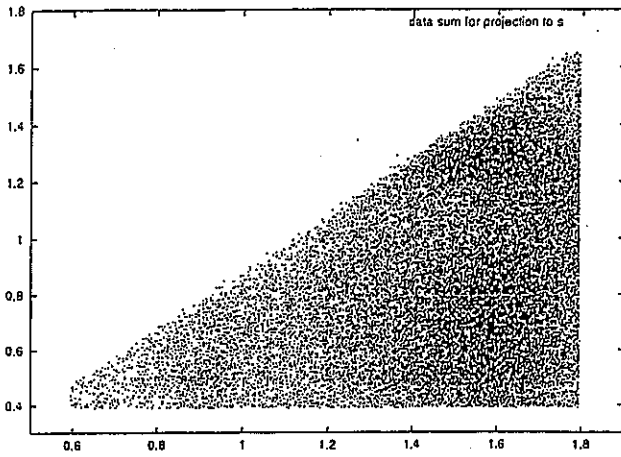
with

$$\Delta_j = \sqrt{s} \Gamma(s)_j / (s - m_j^2 + i\sqrt{s} \Gamma(s)_j),$$

where  $r_j$ : production amplitude of the  $j$ th resonance, and  $\theta_j$ : production phase of the  $j$ th resonance.

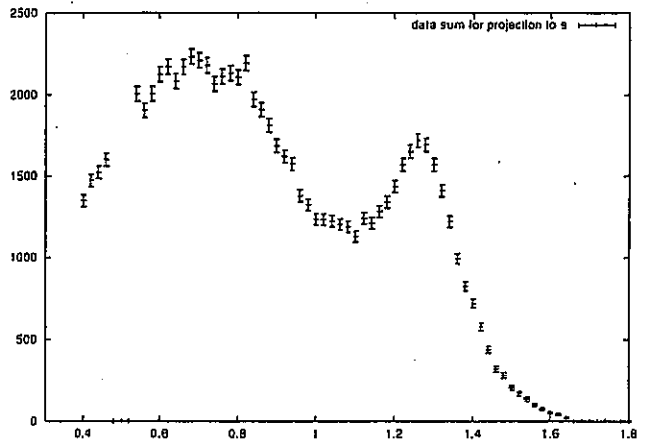
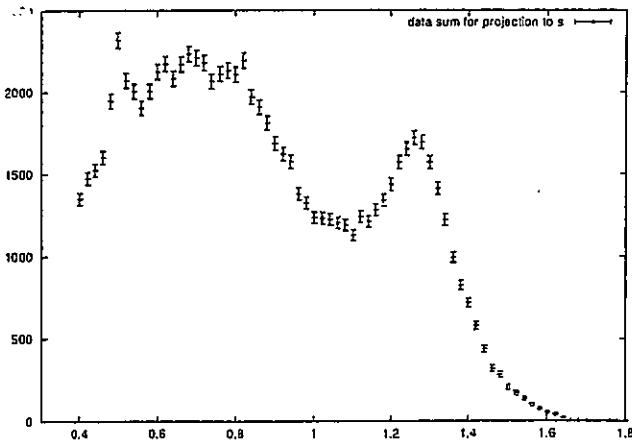
$\pi^0 \pi^0 \pi^0$  scatter plot on the  $2\pi^0$  vs  $3\pi^0$  invariant mass plane.

Rejection of  $K^*(892)$ :  $K_s^0$  cut in  $m_{2\pi}$

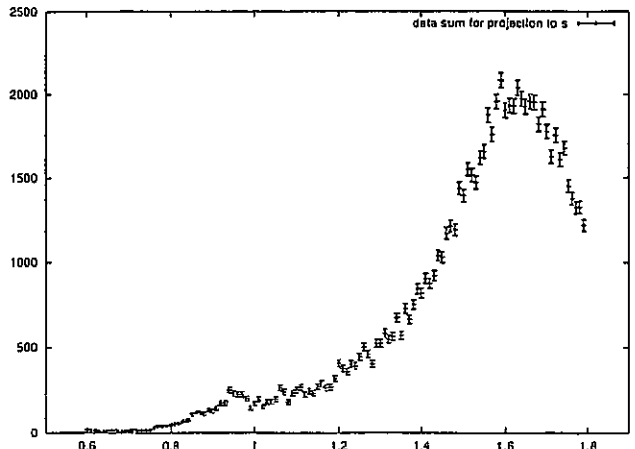
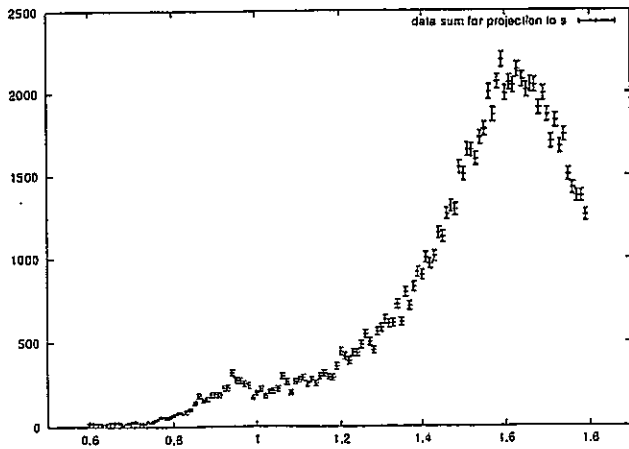


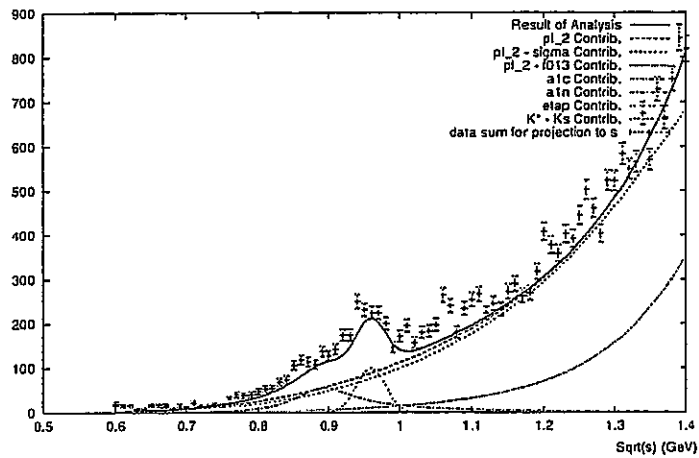
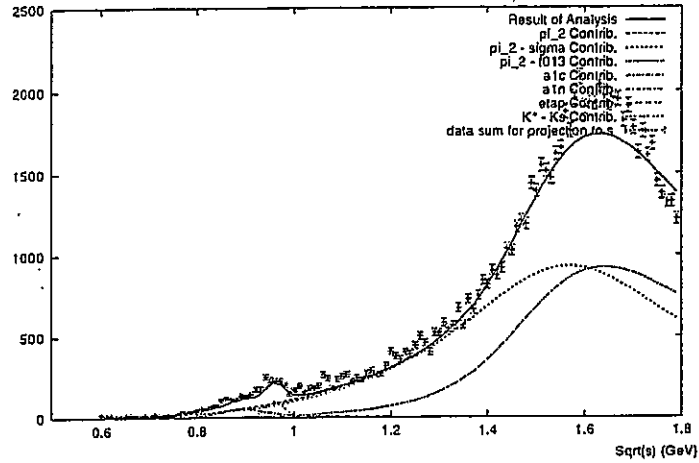
$2\pi^0$  and  $3\pi^0$  invariant mass spectra.

$2\pi^0$



$3\pi^0$



Fitting results with  $a_1^{\chi}$ .

Parameters obtained:

$$m(a_1^{\chi})=884\text{MeV}, \quad \Gamma(a_1^{\chi})=107\text{MeV}, \quad r(a_1^{\chi})=49$$

$$m(\pi_2)=1521\text{MeV}, \quad \Gamma(\pi_2)=615\text{MeV}, \quad r(\pi_2)=85$$

$$m(\sigma)=667\text{MeV}, \quad \Gamma(\sigma)=600\text{MeV}, \quad r(\pi_2 \rightarrow \sigma)=0.59$$

$$m(f_2)=1290\text{MeV}, \quad \Gamma(f_2)=200\text{MeV}, \quad r(\pi_2 \rightarrow f_2)=1.0, \quad \theta(\pi_2 \rightarrow f_2)=-24$$

$$m(\eta')=957\text{MeV}, \quad \Gamma(\eta')=20\text{MeV}, \quad r(\eta')=85$$

## 4. Scalars,

$I=0$   $f_0(600)$  or  $\sigma$ ,  $f_0(980)$ ,  $\underline{f_0(1370)}$ ,  $f_0(1500)$ ,  $\underline{f_0(1710)}$ ,  $f_0(2020)$

$I=1$   $a_0(980)$ ,  $\underline{a_0(1450)}$ ,

- Scalar glueball: around 1.6GeV by the QCD lattice calculations:

1.55(UKQCD), 1.63GeV(Morningstar), 1.74GeV(IBM), \*\*\*GeV(A. Nakamura)

-  $G(1590)$  (GAMS)/ $f_0(1500)$  (XB)/ $f_{0,2}(1710)$  ( $J/\Psi$  radiative decays) have been controversial.

-  $\sigma(555)$ : produced in pp central collision process, WA102, Hadron '95

Unitarity and universality arguments, Hadron '95, M. Pennington

$$SS = 1, \text{ and } F_{\pi\pi} = \alpha(s)T_{\pi\pi}$$

$\alpha(s)$ : slowly varying real function of  $s$ .

$F_{\pi\pi}$  and  $T_{\pi\pi}$  have the same phases and the same structures.

No threshold suppression (Adler 0) exists in production processes, experimentally.

- Relation between scattering and production amplitudes, Hadron '97,

S. Ishida and M. Ishida

- Final state interaction in production processes. Various final state interactions including resonance particles should be taken into account.

-The production process is to be studied independent of the scattering process. The importance and necessity of studies on resonant states in various production processes.

1976-1994: no  $\sigma$  in PDG tables, no low mass scalar below 1 GeV in the analysis of the scattering phase shifts data.

1996-2000:  $f_0(400-1200)$  or  $\sigma$ , re-analyses of the scattering phase shifts data.

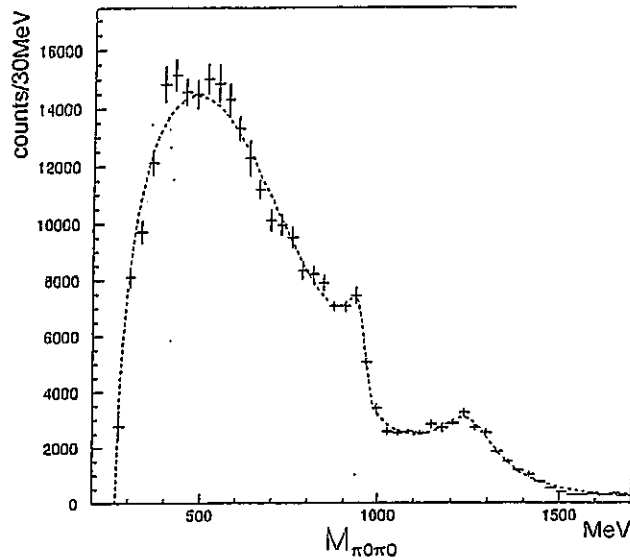
2002:  $f_0(600)$  or  $\sigma$ , evidence for  $\sigma$  in various production processes: pp central collision,  $\bar{p}p$  annihilation,  $J/\Psi$  decays,  $Y$  decays,  $D$  decays in 1995-2003.



$\pi^0\pi^0$  states produced in the pp central collision process.

$pp \rightarrow p_f X^0 p_s, X^0 \rightarrow \pi^0\pi^0$  at 450 GeV/c, GAMS4000, CERN,

D. Alde et al., Phys. Lett. B397 (1997) 350

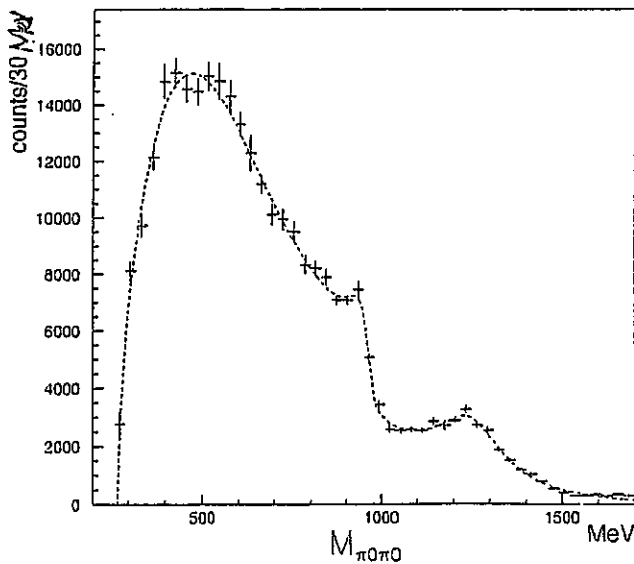


$$M_\sigma = 590 \pm 10 \text{ MeV}$$

$$\Gamma_\sigma = 710 \pm 30 \text{ MeV}$$

$$\begin{aligned} \mathcal{M} &= \mathcal{M}_0 + \mathcal{M}_c \\ &= \frac{\xi_{f_0}}{(s - m_{f_0}^2) + im_{f_0}\Gamma_{f_0}^{tot}(s)} + \frac{\xi_c e^{i\theta}}{(s - m_c^2) + im_c\Gamma_c^{tot}(s)}. \end{aligned}$$

Analyzed with an conventional background (exponential type).



$$\begin{aligned} \mathcal{M} &= \mathcal{M}_0 + \mathcal{M}_{exp} \\ &= \frac{\xi_{f_0}}{(s - m_{f_0}^2) + im_{f_0}\Gamma_{f_0}^{tot}(s)} + a \cdot e^{i\theta} (\sqrt{s}/m_\pi - b)^c \cdot \exp(-d\sqrt{s}/m_\pi). \end{aligned}$$

### 5. Chiral particles

$\sigma(600)$ : a chiral partner of  $\pi$  meson as a Nambu-Goldstone boson.

Low energy effective theory of QCD.

L $\sigma$ M of the Nambu-Jona-Lasinio Type.

M.D.Scadron(1984, 1985),

T.Hatsuda-T.Kunihiro(1985,1988),

S.Ishida-M.Ishida(1996,1997,2000)

$\sigma$  in re-analyses of  $\pi\pi$  scattering phase shifts.

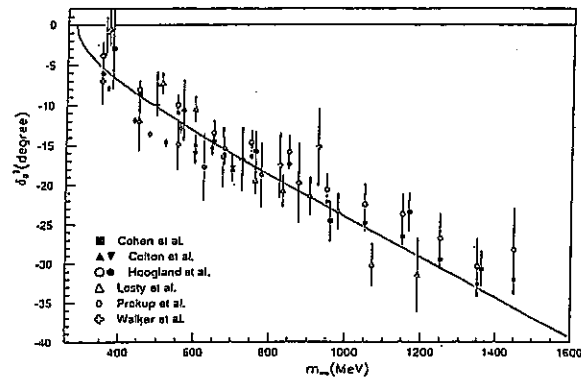
N.N.Achasov and G.N.Shestakov(1994), B.S.Zou and D.V.Bugg(1994), R.Kaminski et al.(1994,99), N.A.Tornqvist and M.Roos((1996), M.Harada et al.(1996), S.Ishida et al.(1996,97), V.V.Anisovich et al.(1997), J.A.Oller and E.Oset(1997), M.P.Locher et al.(1998), K.Igi and K.Hikasa(1999).

Negative phase shifts,

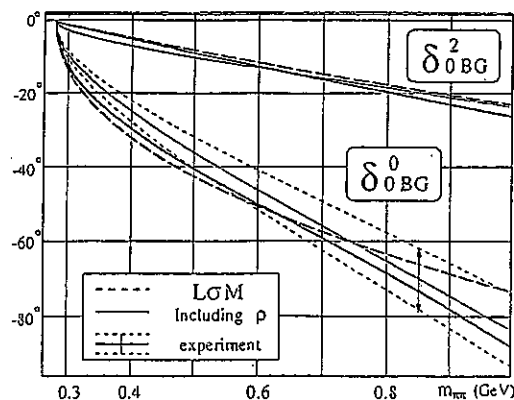
I=2, S wave  $\pi\pi$  phase shifts,  $\delta_0^2$

Compensating  $\lambda\phi^4$  contact interaction in L $\sigma$ M

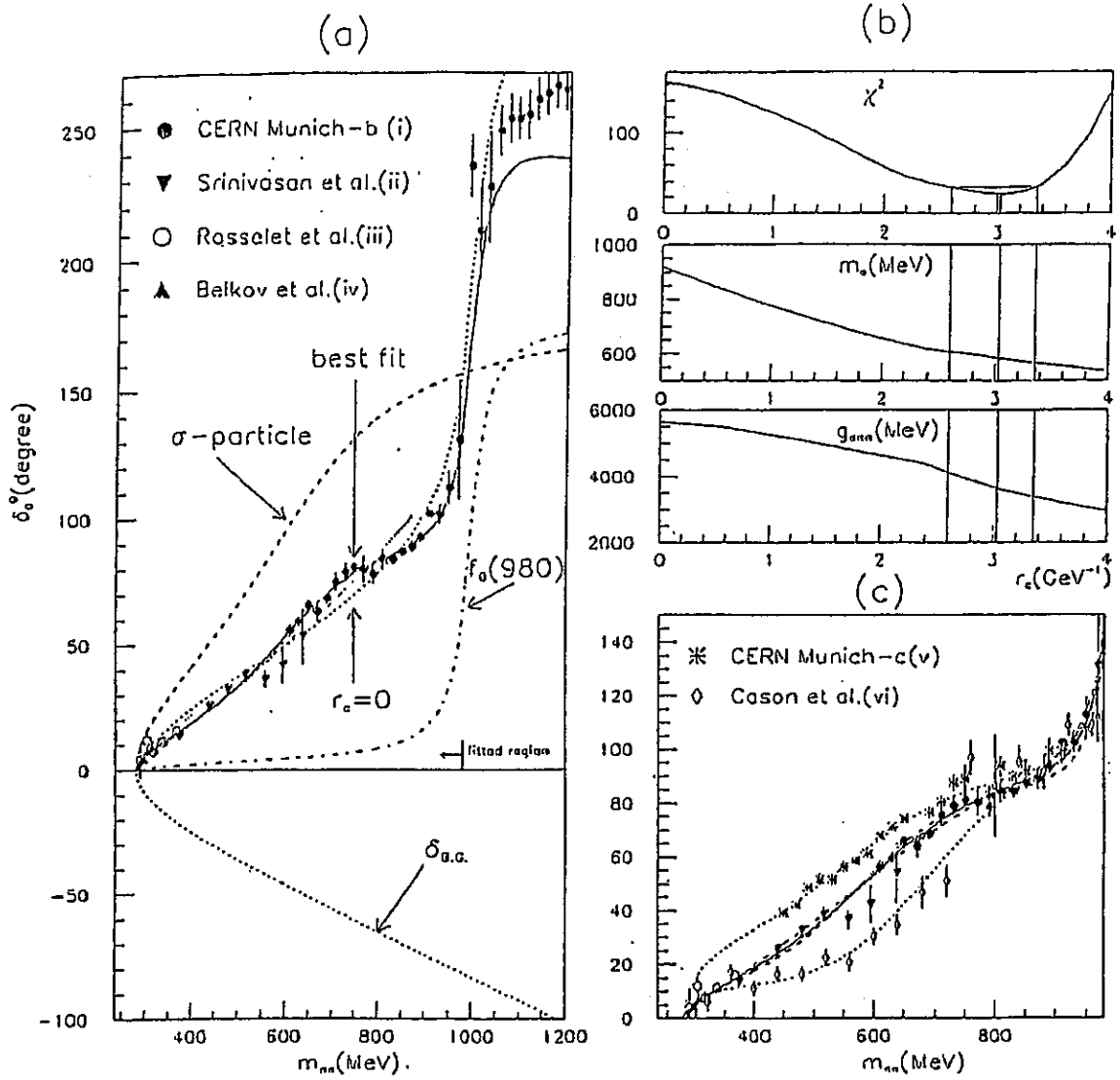
I=2 S wave  $\pi\pi$  phaseshifts.



Repulsive core corresponds compensating  $\lambda\phi^4$  contact interaction.



Analysis on  $l=0$ , S wave  $\pi\pi$  phase shifts  
Interfering amplitude method.



$$M_\sigma = 585 \pm 20 \text{ MeV}, \Gamma_\sigma = 340 \pm 45 \text{ MeV}, r_c = 3.03 \text{ GeV}^{-1}$$

$\sigma$  in production processes

## Studies in production processes

No cancellation mechanism works, necessarily.

Analysis: VMW

pp central coll.;  $J/\Psi \rightarrow \omega\pi\pi$  (DM2 data);  $p\bar{p} \rightarrow 3\pi^0$  (XB);  $Y(2S)$ ,  $Y(2S)$  decays.

mass and width in our observation and analysis.

$\pi\pi$ scattering( $\pi^+\pi^-$ ):	$m_\sigma = 585 \pm 20\text{MeV}$ ,	$\Gamma_\sigma = 385 \pm 70\text{MeV}$
$\pi\pi$ scattering( $\pi^0\pi^0$ ):	$m_\sigma = 588 \pm 12\text{MeV}$ ,	$\Gamma_\sigma = 281 \pm 725\text{MeV}$
pp central coll.:	$m_\sigma = 590 \pm 10\text{MeV}$ ,	$\Gamma_\sigma = 710 \pm 30\text{MeV}$
$J/\Psi \rightarrow \omega\pi\pi$ (DM2):	$m_\sigma = 480 \pm 5\text{MeV}$ ,	$\Gamma_\sigma = 325 \pm 10\text{MeV}$
$p\bar{p} \rightarrow 3\pi^0$ :	$m_\sigma = 540^{+36-29}\text{MeV}$ ,	$\Gamma_\sigma = 385^{+64-80}\text{MeV}$
Y decays:	$m_\sigma = 526^{+48-37}\text{MeV}$ ,	$\Gamma_\sigma = 301^{+145-100}\text{MeV}$

$\tau$  decay (CLEO)  $m_\sigma = 555\text{MeV}$ ,  $\Gamma_\sigma = 540\text{MeV}$

$J/\Psi \rightarrow \omega\pi\pi$ (BESI)  $m_\sigma = 390^{+60-36}\text{MeV}$ ,  $\Gamma_\sigma = 282^{+77-50}\text{MeV}$

$J/\Psi \rightarrow \omega\pi\pi$ (BESII) ( $m_\sigma$ ,  $\Gamma_\sigma$ )

Next talk by Wu Ning

$D \rightarrow 3\pi$  decays (E791)  $m_\sigma = 483^{+25-26}\text{MeV}$ ,  $\Gamma_\sigma = 338^{+45-42}\text{MeV}$

$\sigma(600)$ : mass in a range 500–600MeV, width in 300–400MeV.

Y decays,  $Y(nS) \rightarrow Y(mS)\pi\pi$ .

T. Komada et al., P.L. B 508 (2001) 31; M. Ishida et al., P.L. B 518 (2001) 47.

Chiral symmetric effective amplitude in linear  $\sigma$  model.

- Non derivative coupling.

$$F_{\sigma+2\pi} = F_{\sigma} + F_{2\pi} \equiv (\text{polarization vectors})(\text{coupling})(m_{\pi}^2 - s)/(m_{\sigma}^2 - s).$$

Adler 0 occurs at  $s = -(p_1^2 + p_2^2) = m_{\pi}^2$ , leading to the threshold suppression.

- Derivative coupling.

$$F_{2\pi}^G \equiv (\text{polarization vectors})(\text{coupling})m_{\gamma} \cdot m_{\gamma} p_{10} p_{20} \\ \equiv (\text{polarization vectors})(\text{coupling})m_{\gamma} \cdot m_{\gamma} [(m_{\gamma} + m_{\gamma})/2]^2 = \text{constant}.$$

$p_{10}$  and  $p_{20}$  are energy of pions. Not small near the threshold.

$F_{\sigma+2\pi}$  and  $F_{2\pi}^G$

Fitting results

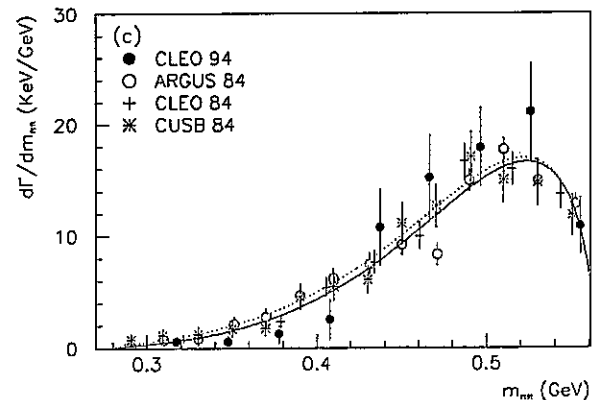
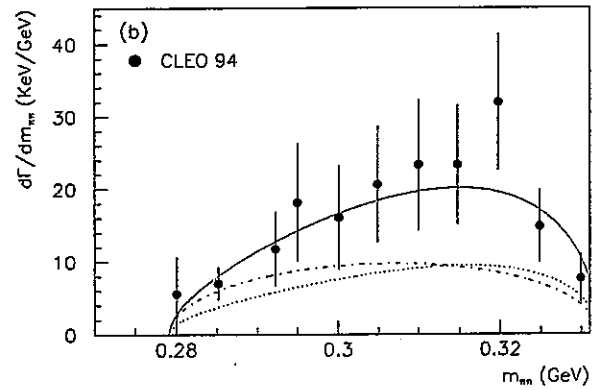
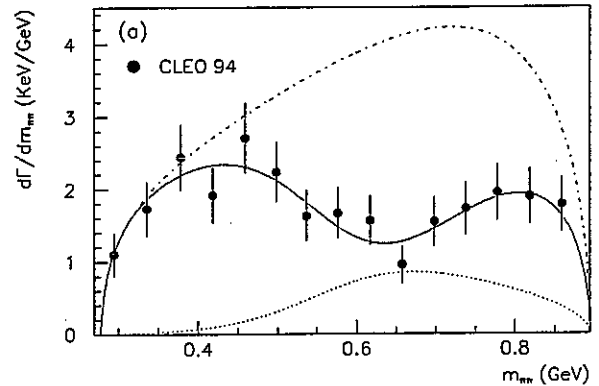
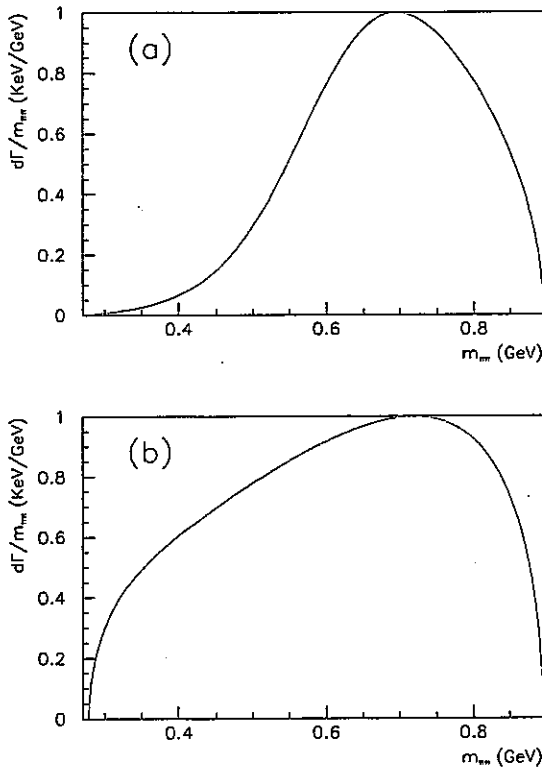


Fig. 1. The  $\pi\pi$  mass spectrum obtained (a) from  $\mathcal{F}_{\sigma+2\pi}^{\text{phen}}$  in Eq. (3.1) and (b) from  $\mathcal{F}_{2\pi}^G$  in Eq. (2.11) in case of  $\Upsilon(3S) \rightarrow \Upsilon(1S)\pi\pi$  decay. The peaks are normalized to unity. Both amplitudes satisfy the derivative coupling property of  $\pi$ , and in (a) the spectrum is suppressed in the  $\pi\pi$  threshold region, while (b) shows a steep increase from the threshold.

Fig. 2. Fit to the  $\pi\pi$  invariant mass spectrum by the amplitude Eq. (3.1) explicitly consistent with chiral constraint: (a)  $\Upsilon(3S) \rightarrow \Upsilon(1S)\pi\pi$ , (b)  $\Upsilon(3S) \rightarrow \Upsilon(2S)\pi\pi$  and (c)  $\Upsilon(2S) \rightarrow \Upsilon(1S)\pi\pi$ . The respective contributions from  $\mathcal{F}_{\sigma+2\pi}^{\text{phen}}$  and from  $\mathcal{F}_{2\pi}^G$  are shown by dotted and dot-dashed lines. The  $\mathcal{F}_{2\pi}^G$  contribution becomes dominant in the case (a)  $\Upsilon(3S) \rightarrow \Upsilon(1S)$  transition in comparison with the case (c)  $\Upsilon(2S) \rightarrow \Upsilon(1S)$  transition.

Existence of  $\kappa(900)$  in the analysis of  $K\pi$  scattering phase shifts.

$$m_\kappa = 905^{+63}_{-30} \text{MeV}, \quad \Gamma_\kappa = 545^{+235}_{-110} \text{MeV}.$$

S. Ishida et al, Prog. Theor. Phys. 98(1997)621

$$m_\kappa = 897 \text{MeV}, \quad \Gamma_\kappa = 322 \text{MeV}. \quad \text{D.Black et al. (1998).}$$

$$m_\kappa = 727 \text{MeV}, \quad \Gamma_\kappa = 526 \text{MeV}. \quad \text{E.VanBeveren et al. (1996)}$$

$$m_\kappa = 900 \text{MeV}, \quad \text{J.A.Ollerand E.Oset (1999)}$$

## in production processes

D decays,  $D \rightarrow K\pi\pi$  (E791 FNAL)

C. Goebel

 $D \rightarrow K\pi\mu\nu$  (FOCUS FNAL)

interference with S-wave res.

 $J/\Psi \rightarrow K^*K\pi$  (BES)

Wu Ning, and T. Komada

 $\kappa(900)$ : mass in a range 800–900MeV, width in 300–600MeV.New scalar nonet:  $\sigma(600)$ ,  $\kappa(900)$ ,  $a_0(980)$ ,  $f_0(980)$ .

a chiral partner of the ground state pseudoscalar nonet..

S. Ishida and M.Ishida(1997), M.D.Scadron(1982)

Mass and width by SU(3) L $\sigma$ M (undelined values: inputs)

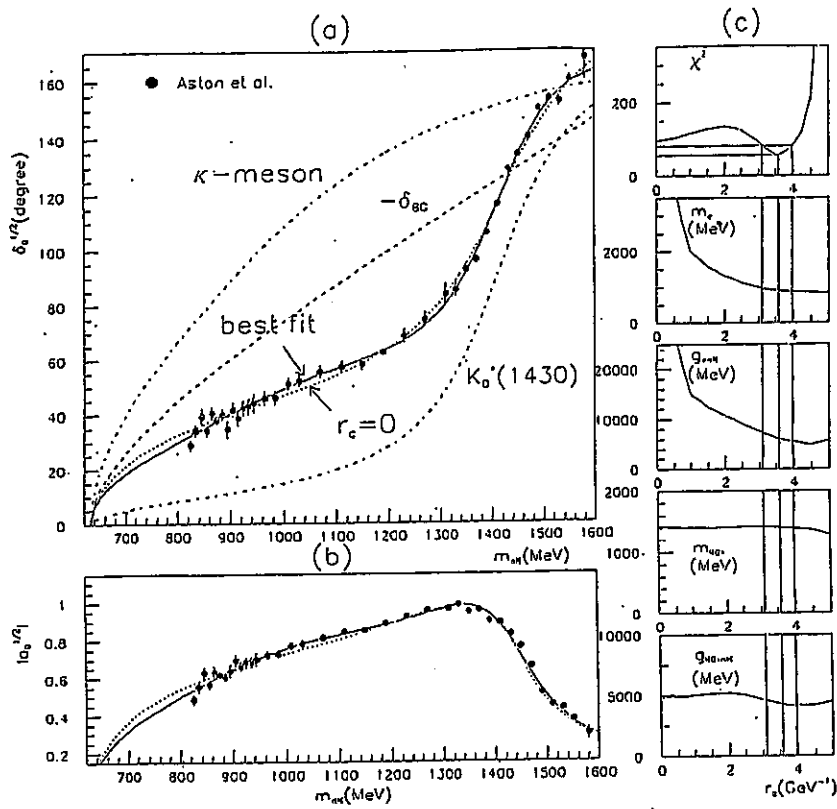
	M (MeV)	m(exp)(MeV)	$\Gamma$ (MeV)	$\Gamma$ (exp)(MeV)
$\sigma$	<u>535–650</u>	<u>535–650</u>	400–800	385+–70
$\kappa$	<u>905+65–30</u>	<u>905+65–30</u>	300–600	545+235–110
$\delta = a_0(980)$	900–930	982.7+–2.0	110–170	95+–14
$\sigma' = f_0(980)$	1030–1200	993.2+–9.5	0–300	67.9+–9.4

M. Ishida, Sigma WS, Kyoto(2000).

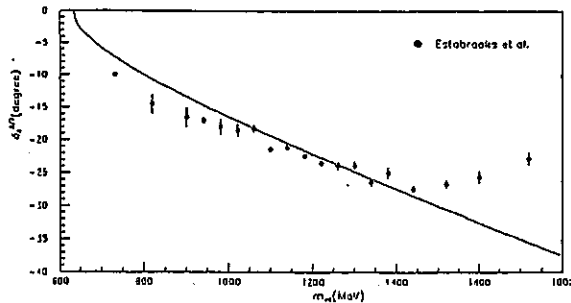
Analysis of  $K\pi$  scattering phase shift data

Data from LASS at SLAC,  $Kp \rightarrow K^-\pi^+n$  at 11GeV/c, D. Aston et al.(1988)

$I=1/2$   $K\pi$  scatt. S wave phase shifts and analysis



$I=3/2$   $K\pi$  scattering phase shifts, P. Esterbrook et al. (1978)



New axial-vector nonet, a chiral partner of the ground state vector nonet.

S Ishida et al., Prog. Theor. Phys. 95 (1996) 745.

S. Ishida and M. Ishida, Hadron '97.

$$a_1^\chi, K_1^\chi, f_1^\chi, f_1^{\chi'}$$

Search for chiral particles

Possible channels for  $\sigma$ ,  $\kappa$  and  $a_1^\chi$  studies in  $J/\Psi$  and  $\Psi'$  decays.

$\sigma$  in various  $J/\Psi$  and  $\chi_{cJ}$  ( $\Psi' \rightarrow \gamma\chi_{cJ}$ ) decays.

$\kappa$  in  $\chi_{cJ}$  decays.  $\chi_{cJ} \rightarrow \pi^+\pi^-K^+K^-$  ( $K\pi K$ ).

$a_1^\chi$  in  $J/\Psi \rightarrow \rho a_1^\chi$ ,  $J/\Psi \rightarrow \gamma\pi a_1^\chi$ , and  $\chi_{c0} \rightarrow \pi a_1^\chi$ .

Possible channels for  $\sigma$ ,  $\kappa$  and  $a_1^\chi$  studies in  $J/\Psi$  and  $\Psi'$  decays

For  $\kappa$ :  $\Psi' \rightarrow \gamma\chi_{cJ}$ ,  $\chi_{cJ} \rightarrow \pi^+\pi^-K^+K^-$   $K\pi K$

For  $a_1^\chi$ :  $J/\Psi \rightarrow X$ ,  $X \rightarrow VA^+, A^+A^-$   $\rho a_1^\chi$ ,

$J/\Psi \rightarrow \gamma X$ ,  $X \rightarrow PA^+, A^+A^+$   $\pi a_1^\chi$ ,

$\Psi' \rightarrow \gamma\chi_{cJ}$ ,  $\chi_{c0} \rightarrow PA^+, TA^+, A^+A^+$   $\pi a_1^\chi$ ,

$\chi_{c1} \rightarrow PA^+, SA^+, TA^+, A^+A^+$   $\pi a_1^\chi$ ,

$\chi_{c2} \rightarrow PA^+, TA^+, A^+A^+$   $\pi a_1^\chi$ .

S: scalar, P: pseudoscalar, V: vector, T: tensor, A+: axialvector with positive C parity and A-: axialvector with negative C parity.



$U\tilde{(12)}$ , classification of mesons

	Non relativistic view		relativistic view			
model	non relativistic model		NLJ model			
	NRQM		COQM + chiral symmetry			
symmetry	LS symmetry		chiral symmetry			
evidence	LS coupling		$\pi$ -nonet: NG-boson			
ground states	$0^+$	$\pi$ -nonet	$0^+$	$0^{++}$	$0^{+-}$	$0^+$
	$1^-$	$\rho$ -nonet	$1^-$	$1^{++}$	$1^{+-}$	$1^+$
		$I=1$	$I=0$		$I=1/2$	
ground states	$0^+$	$\pi$	$\eta$	$\eta'$	$K$	
	$1^-$	$\rho$	$\omega$	$\phi$	$K^*$	
1 <sup>st</sup> exc. states	$1^{++}$	$b_1(1235)$	$h_1(1170)$	$h_1(1380)$	$K_1(1270)$	
	$0^{++}$	$a_0(1450)$	$f_0(1370)$	$f_0(1710)$	$K_0^*(1430)$	
	$1^{++}$	$a_1(1230)$	$f_1(1285)$	$f_1(1420)$	$K_1(1400)$	
	$2^{++}$	$a_2(1320)$	$f_2(1275)$	$f_2(1525)$	$K_2^*(1430)$	

Exotic quantum numbers;  $0^{+-}$ ?  $1^{--}$ ?  $0^{--}$ ? Naturally accommodated in  $U\tilde{(12)}$

Where do we find them?

$J/\Psi$  decays

$1^- \rightarrow 1^- + 0^{++}$  S wave decay

$1^- \rightarrow 1^+ + 0^{+-}$  S wave decay

$\chi_{c1}(1^{++}), \chi_{c0}(0^{++})$  decays

$1^{++} \rightarrow 1^+ + 0^{+-}$  S wave decay

$1^{++} \rightarrow 1^- + 0^+$  P wave decay

$0^{++} \rightarrow 0^{++} + 0^+, 0^{+-} + 0^+$  S wave decay