

Numerical Calculations of Excited Baryon

Masses in Lattice QCD

Shoichi Sasaki (Univ.of Tokyo)

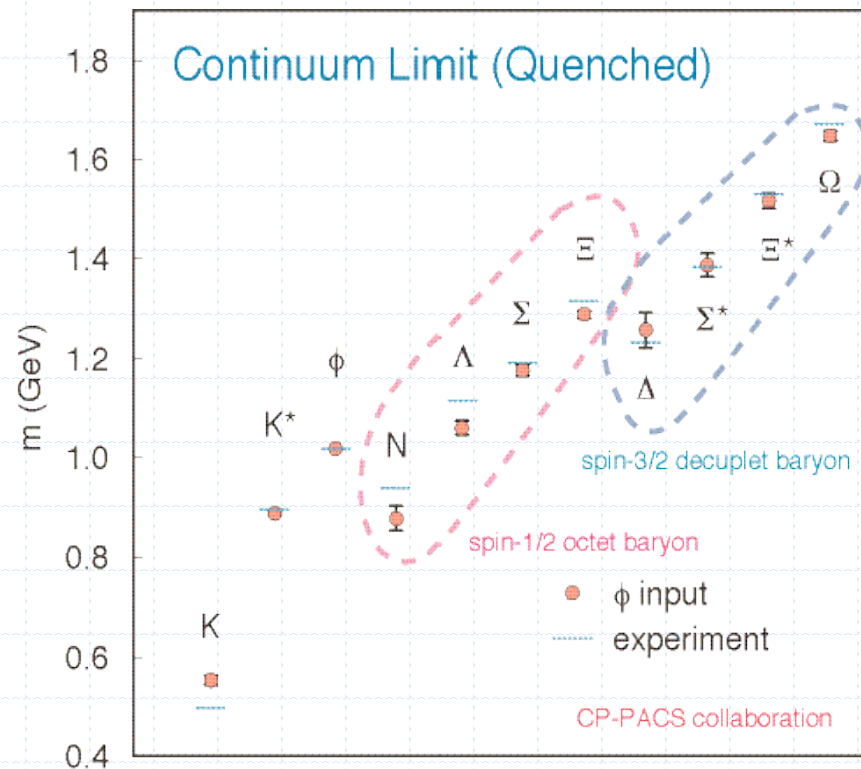
References:

S. Sasaki, T. Blum and S. Ohta, PRD65 (02) 074503

S. Sasaki, K. Sasaki, T. Hatsuda and M. Asakawa, hep-lat/0209059

Hadron Mass Spectrum from Lattice QCD

Experimental results are reproduced to within about **5-10%** errors.



CP-PACS collaboration,
Phys. Rev. Lett. 84 (2002) 238

➡ But this success is restricted to **ground states**, and does not fully apply to excited states

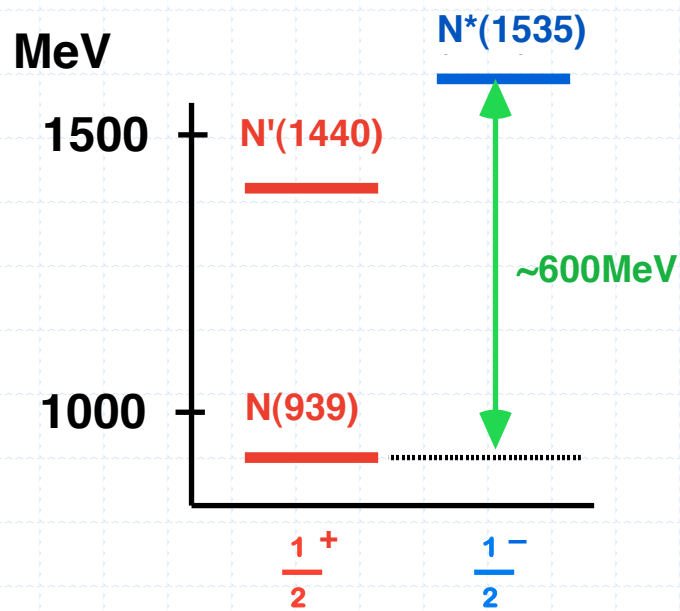
Lattice studies of N^* spectrum

- ✓ **Lee-Leinweber**
D \square 34 action hep-lat/9809095, D234 action, hep-lat/0011060, 0110164.
- ✓ **Sasaki-Blum-Ohta (RIKEN-BNL)**
Domain wall fermion, hep-lat/9909093, *Phys. Rev. D* **65** (2002) 074503.
- ✓ **Richards et al (UKQCD-QCDSF-LHPC)**
Clover fermion, hep-lat/0011025, *Phys. Lett. B* **532** (2002) 63.
- ✓ **Melnitchouk et al (Adelaide)**
Fat-link clover fermion, hep-lat/0202022.
- ✓ **Nakajima-Matsufuru-Nemoto-Suganuma**
Clover fermion & anisotropic action, hep-lat/0204014.
- ✓



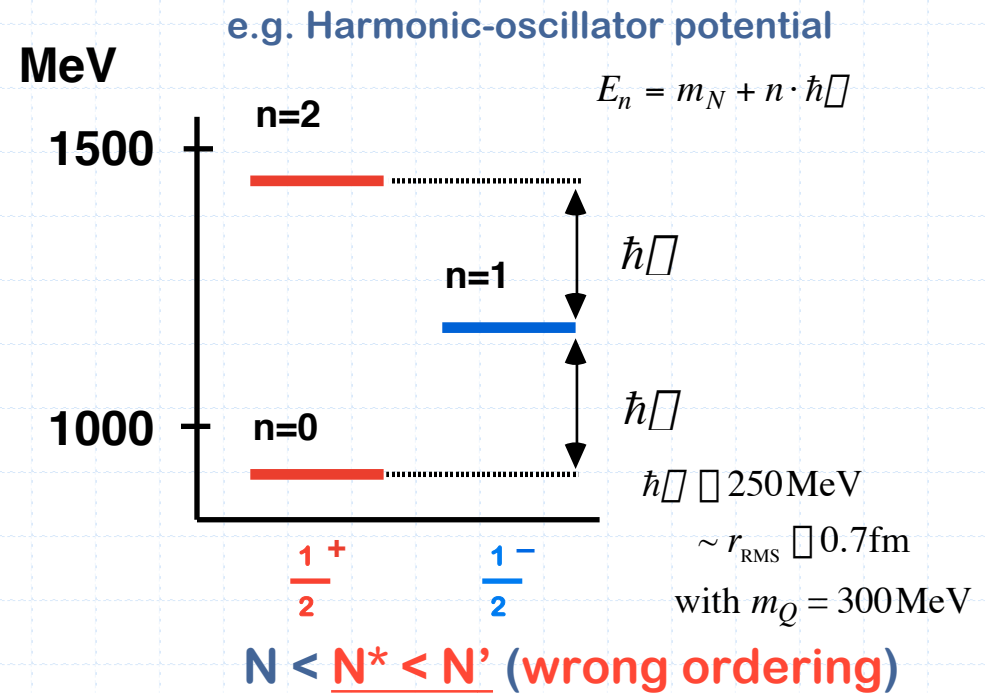
A puzzle in nucleon excited states

Experiment



Potential Models

Naive quark model
Bag model



👉 Need lattice QCD calculation to resolve it.

How to compute the nucleon mass

Consider time-ordered, 2pt function:

$$\langle N(t)\bar{N}(0) \rangle_{\square\square} = \int_{\vec{x}} \langle 0 | T \{ N_{\square}(\vec{x}, t) \bar{N}_{\square}(\vec{0}, 0) \} | 0 \rangle$$

interpolating operator: $N = \int_{abc} q_a q_b q_c$ with $J=1/2, I=1/2$

The lowest-lying state dominates for large t

$$\langle N(t)\bar{N}(0) \rangle_{\square\square} = \underbrace{(1 + \int_4)_{\square\square} \int(t) A_N e^{\int M_N t}}_{\text{nucleon}} + \underbrace{(1 - \int_4)_{\square\square} \int(\int t) A_N e^{+M_N t}}_{\text{anti-nucleon}}$$

Particle contribution:

$$\text{Tr}\{P_+ \langle N(t)\bar{N}(0) \rangle\} \sim e^{\int M_N t} \quad \text{with the projection operator } P_+ = (1 + \int_4)/2$$

Spin-1/2 interpolating operators

Baryon operators under $SU(2)_L \times SU(2)_R$

✓ belong to the $(1/2,0)+(0,1/2)$ chiral multiplet

✓ without derivative

$J^P=1/2^+$ operator:

$$N_1^+ = \epsilon_{abc} (u_a^T C \epsilon_{5} d_b) u_c \quad \text{and} \quad N_2^+ = \epsilon_{abc} (u_a^T C d_b) \epsilon_{5} u_c$$

$J^P=1/2^-$ operator: $N_i^\square = \epsilon_{5} N_i^+$

$$N_1^\square = \epsilon_{abc} (u_a^T C \epsilon_{5} d_b) \epsilon_{5} u_c \quad \text{and} \quad N_2^\square = \epsilon_{abc} (u_a^T C d_b) u_c$$

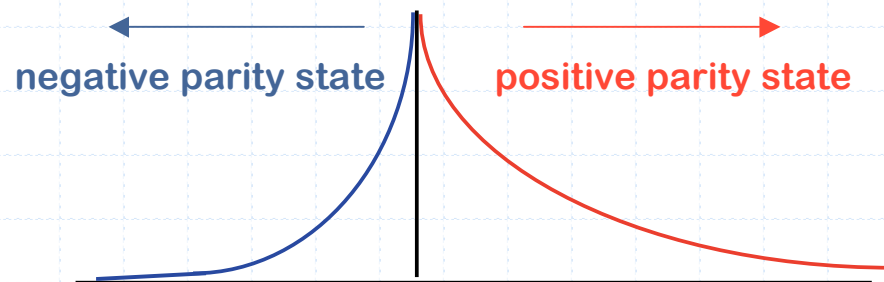
It turns out that $\langle N^\square(t) \bar{N}^\square(0) \rangle = \epsilon_{5} \langle N^+(t) \bar{N}^+(0) \rangle \epsilon_{5}$

$$\langle N^+(t) \bar{N}^+(0) \rangle = (1 + \gamma_4) \Gamma(t) A_{N_+} e^{+M_{N_+} t} + (1 - \gamma_4) \Gamma(\square t) A_{N_+} e^{+M_{N_+} t} \quad (\text{positive parity})$$

$$\epsilon_{5} (1 + \gamma_4) \Gamma(\square t) A_{N_\square} e^{+M_{N_\square} t} \quad \epsilon_{5} (1 - \gamma_4) \Gamma(t) A_{N_\square} e^{+M_{N_\square} t} \quad (\text{negative parity})$$

Wrap-around effect

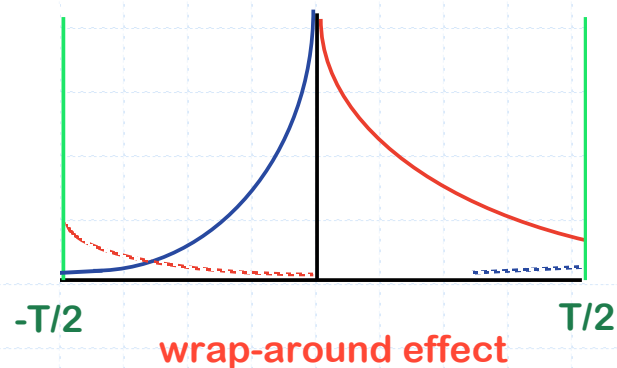
After projection with operator P_+



$$\text{Tr}\{P_+ \langle N^+(t) \bar{N}^+(0) \rangle\}$$

$$\sim \theta(t) A_{N_+} e^{\square M_{N_+} t} + \theta(-t) A_{N_-} e^{+M_{N_-} t}$$

Finite extent T with (anti-) periodic boundary condition



$$\text{Tr}\{P_+ \langle N^+(t) \bar{N}^+(0) \rangle\}$$

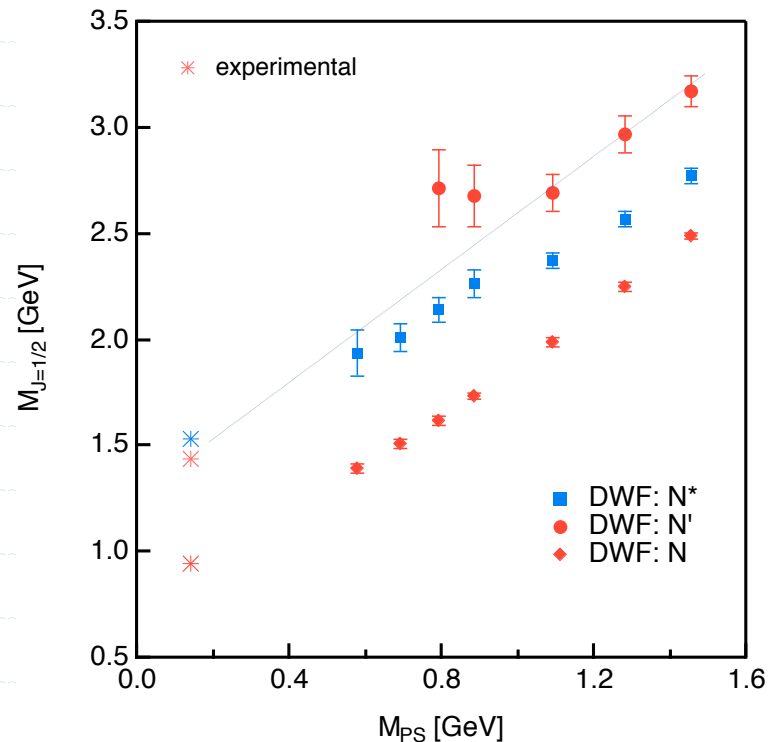
$$b = \begin{cases} +1 & \text{(PBC)} \\ -1 & \text{(APBC)} \end{cases}$$

$$\sim \begin{cases} A_{N_+} e^{\square M_{N_+} t} + b \cdot A_{N_-} e^{+M_{N_-} (t-T)} & \text{for } T/2 > t > 0 \\ A_{N_-} e^{+M_{N_-} t} + b \cdot A_{N_+} e^{\square M_{N_+} (t+T)} & \text{for } -T/2 < t < 0 \end{cases}$$

☛ Necessary to take an appropriate boundary condition

Nucleon excited states in lattice QCD

DWF calculation $16^3 \times 32 \times 16$ at $\beta = 6.0$



S. Sasaki, T. Blum, S. Ohta, PRD65 (02) 074503

□ Large mass splitting of $N(939)$ - $N^*(1535)$ is observed.

• $M_N \approx 1.0 \text{ GeV}$ and $M_{N^*} \approx 1.6 \text{ GeV}$.

• good agreement with the experimental value.

□ $N'(1440)$ state is heavier than N^* state (in the heavy quark region).

• inverted level order ?

• missing the Roper $N'(1440)$?

→ Possible systematic errors

✓ quenching

✓ finite lattice spacing

✓ non-negligible higher states contribution

✓ finite volume effect

To resolve the remaining puzzle

Need information in the **lighter** quark mass region

I. Analysis to take account of higher lying states

☞ **Maximum Entropy Method**

II. Finite size effect might be serious for excited states

☞ **Large physical volume** $> (2.0 \text{ fm})^3$

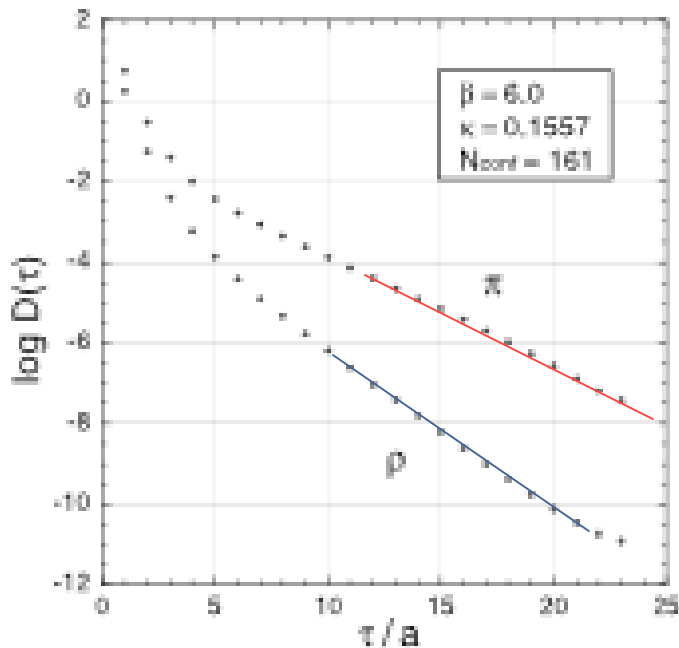
III. Similarity of the level pattern is found in the \square channel

☞ Not only spin-1/2 baryon but also **spin-3/2 baryon**



Measure the hadron masses in lattice QCD

Traditional approach: **Asymptotic analysis**

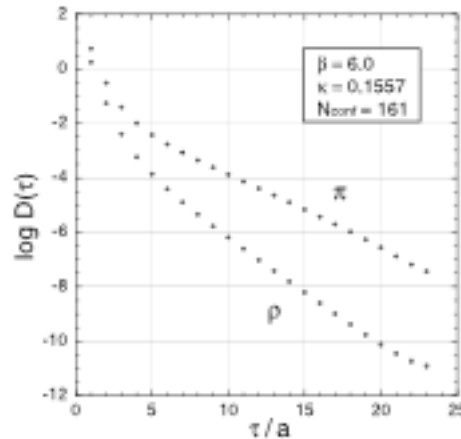


$$\begin{aligned}
 D(\tau) &= \int d^3x \langle h(\vec{x}, \tau) h^\dagger(0, 0) \rangle \\
 &= \sum_n |\langle 0 | h | n \rangle|^2 e^{-M_n \tau} \sum_{\vec{x}} e^{-M_0 \tau}
 \end{aligned}$$

✓ Allow us to fit only the large- τ behavior of a hadron correlator

- Data truncation
- Discard information for excited states

Spectral functions in lattice QCD



Nakahara, Asakawa, Hatsuda, PRD60 (99) 091503

PPNP46 (01) 459

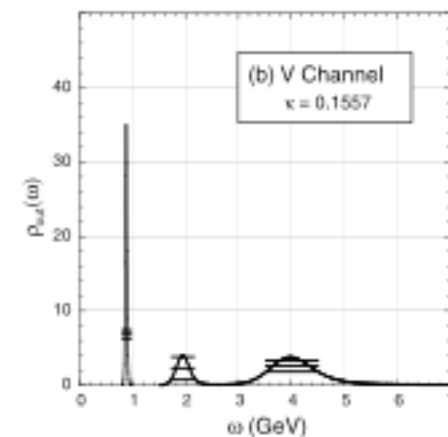
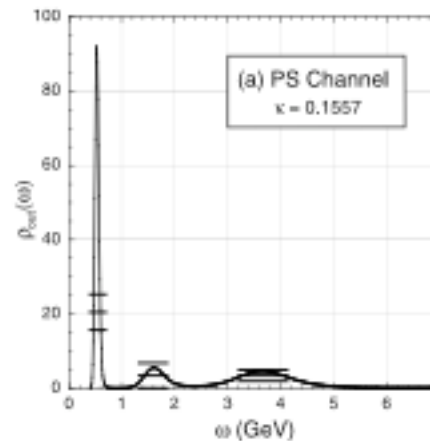
pseudo-scalar, vector, scalar, axial-vector

Yamazaki (CP-PACS collaboration), PRD65 (02) 014501

pseudo-scalar, vector (continuum limit)

Application only for mesons !!

MEM analysis



Maximum Entropy Method (MEM)

- Spectral function $A(\omega)$ gives us much information

$$D(\omega) = \int_0^{\infty} d\tau A(\tau) e^{j\omega\tau}$$

Input: $N_{\tau} \sim O(10)$

Output: $N_{\omega} \sim O(100 - 1000)$

But, the reconstruction of $A(\omega)$ from $D(\omega)$ is not unique !

Assume : the probability of $A(\omega)$ can be assigned for given $D(\omega)$

→ The most probable $A(\omega)$ is inferred from $P[A|D]$

$$A_{\text{imag}}(\omega) = \int [dA] A(\omega) \cdot P[A|D]$$

Note: the posterior probability $P[A|D]$ is not the conditional probability of $D(\omega)$ given $A(\omega)$; $P[D|A]$ (the likelihood function)

More about MEM

Bayesian statistics

posterior probability

$$P[A | D] = \frac{P[D | A]P[A]}{P[D]}$$

likelihood function

$$P[D | A]P[A]$$

evidence

prior function

- Central limiting theorem

$$P[D | A] \sim e^{-\chi^2/2}$$

- Information entropy (Shanon-Jaynes)

$m(\square)$: default model

α : the entropy strength

$$P[A] \sim e^{\alpha S} = \exp \left\{ \alpha \int d\square \left[A(\square) m(\square) A(\square) \ln \frac{A(\square)}{m(\square)} \right] \right\}$$

[Bryan's method]

To find the maximum of $P[A|D]$ with respect to $A(\square)$ at fixed $D(\square)$,

• Minimize $Q = \chi^2/2 - \alpha S$ to use singular value decomposition (SVD)

Details of the simulation

Gauge: Standard plaquette action

$\beta=6.0$, $a^{-1}=2.0$ GeV

three lattice sizes $32^3 \times 32$, $V=(3.0 \text{ fm})^3$, 200 configs

$24^3 \times 32$, $V=(2.2 \text{ fm})^3$, 300 configs

$16^3 \times 32$, $V=(1.5 \text{ fm})^3$, 444 configs

Fermion: Wilson fermions

4 quark masses ($M_{\square} / M_{\square} = 0.69 - 0.91$)

Point source - Point sink

P.B.C. + A.P.B.C. for the temporal direction

Basic result:

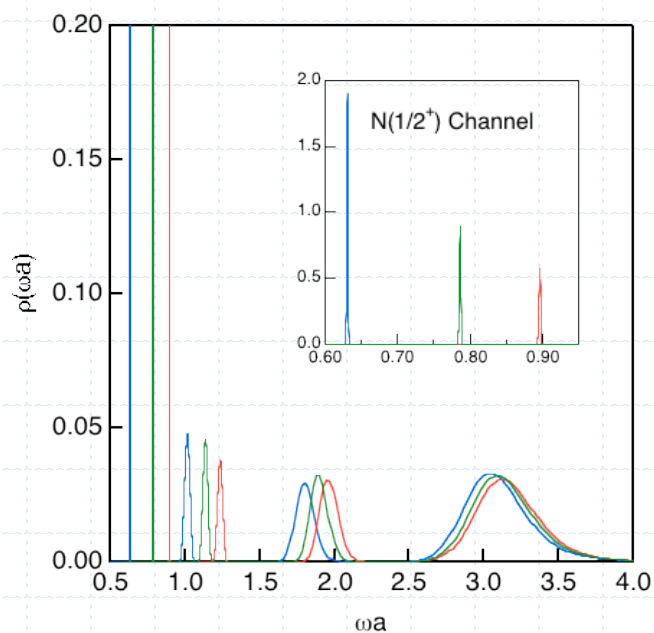
$K_c=0.1568(15)$

Nucleon and Delta Spectral functions

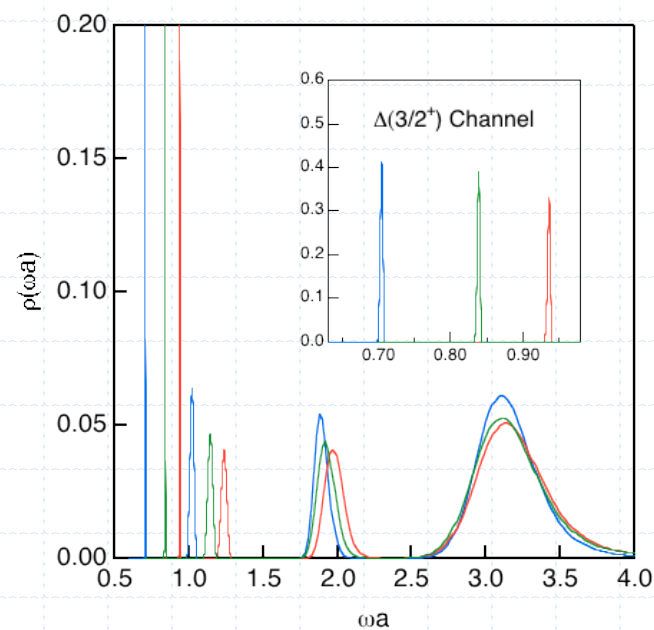
MEM analysis

beta=6.0, $V=24^3 \times 32$ (Wilson fermions)

$\beta_{\max} a \sim 2\beta$, $N_{\beta} \sim O(500)$, $m_0 a = 4 \times 10^{-3}$ (N), 8×10^{-3} (Δ)



Nucleon channel

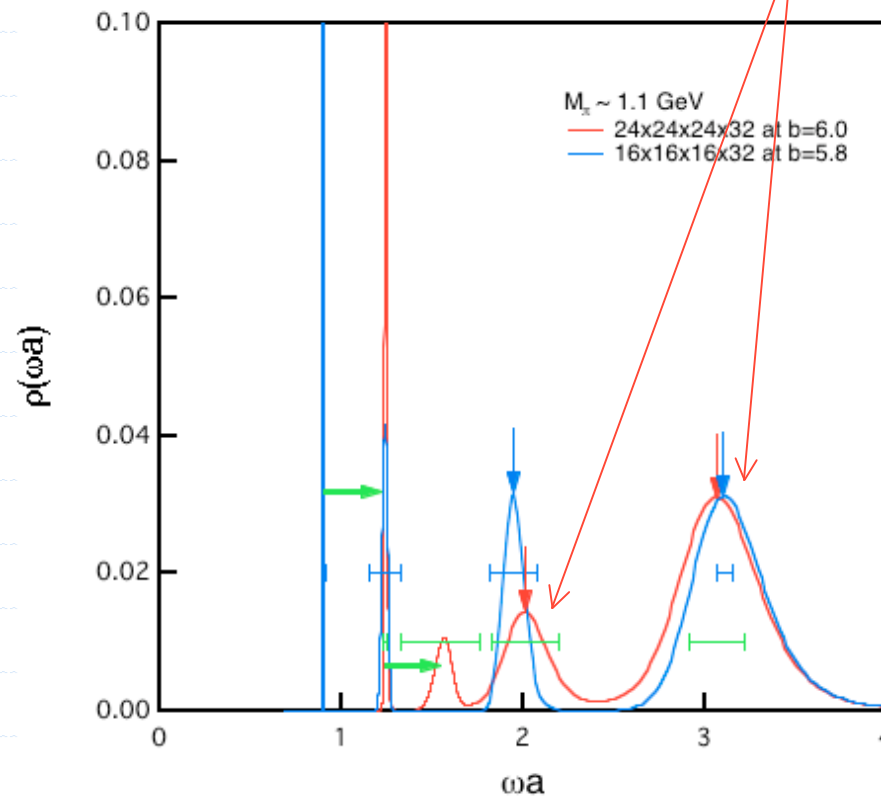


Delta channel

Baryonic bound state of doublers ?

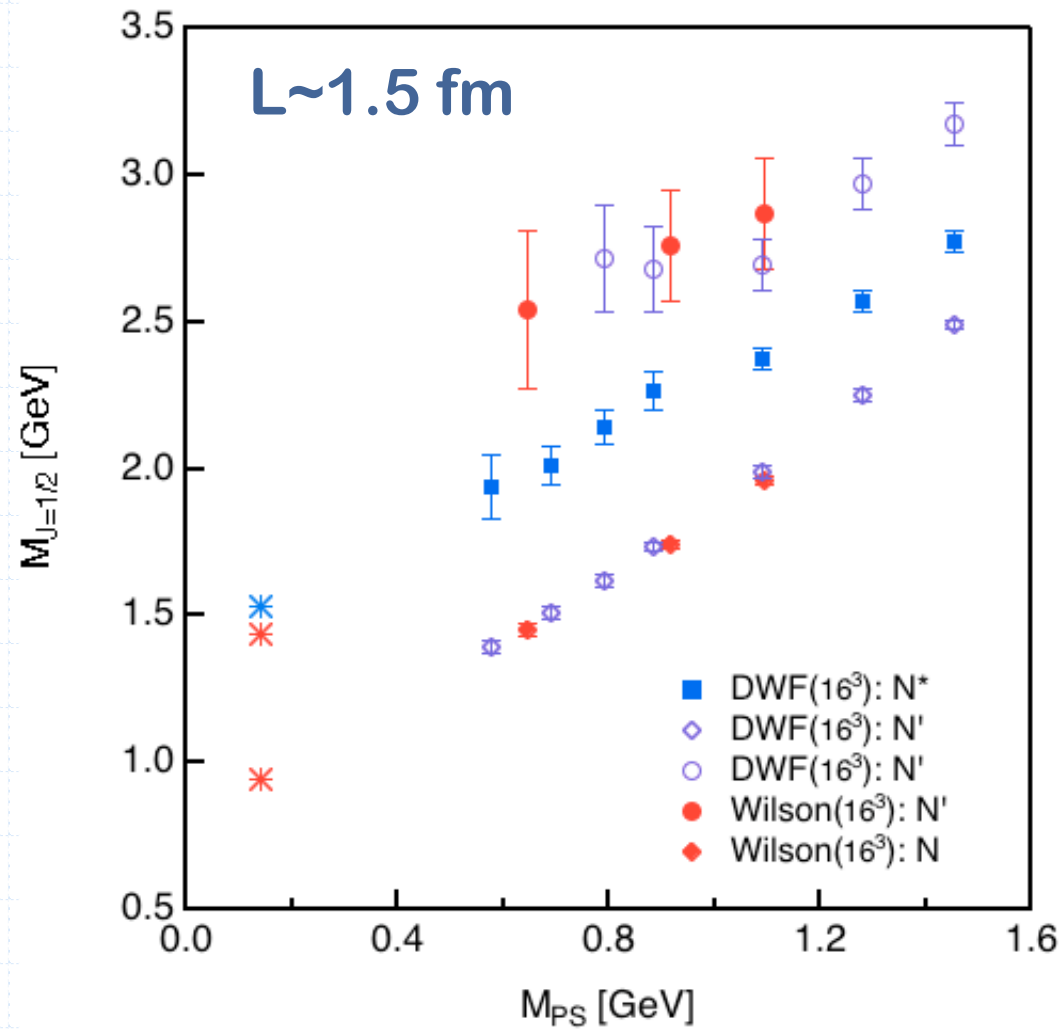
unphysical states

infinitely heavy in the continuum limit ($a \rightarrow 0$)



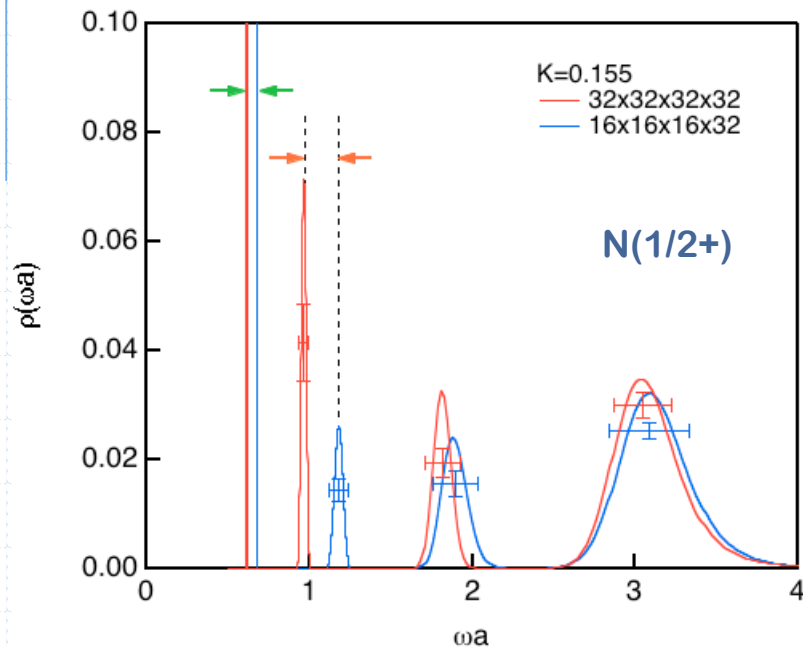
CP-PAC collaboration already found the **mesonic** bound state of doublers (Yamazaki et al., PRD65 (02) 014501)

Results for $16^3 \times 32$ lattice at $\beta=6.0$

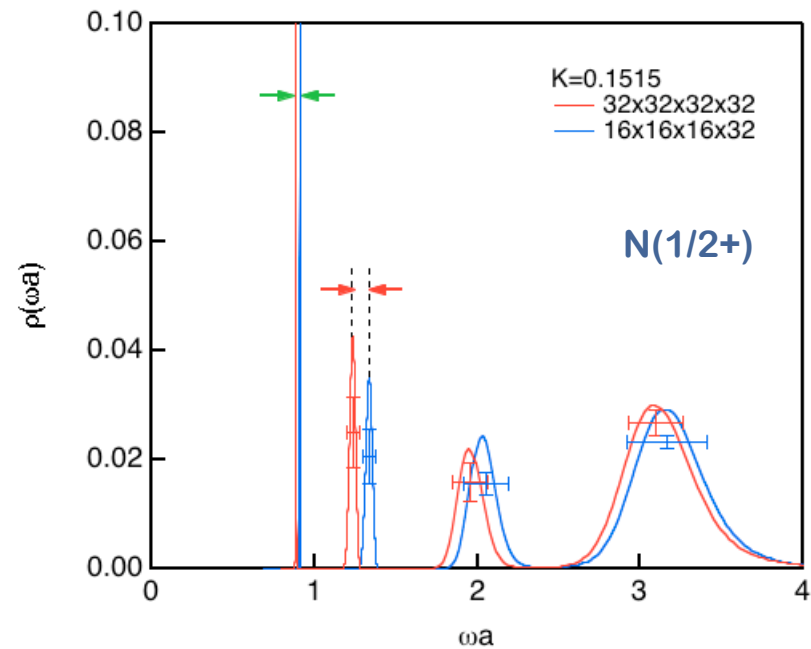


Large finite-size effect on the N' state

The observed size effect is **much larger** for the (physical) **excited state** than for the **ground state**



Light ($m_\pi \sim 0.6 \text{ GeV}$)

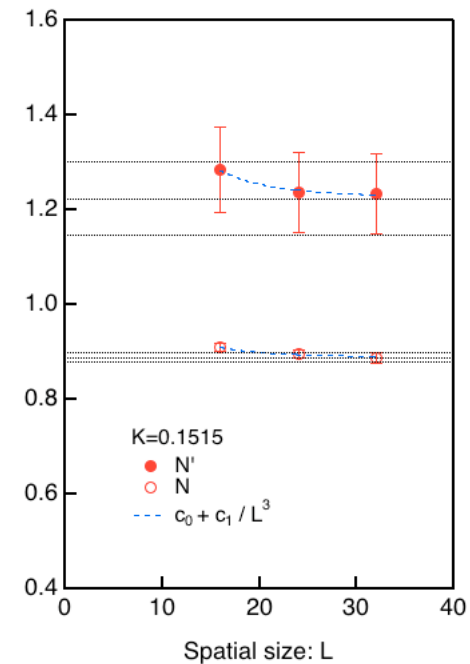
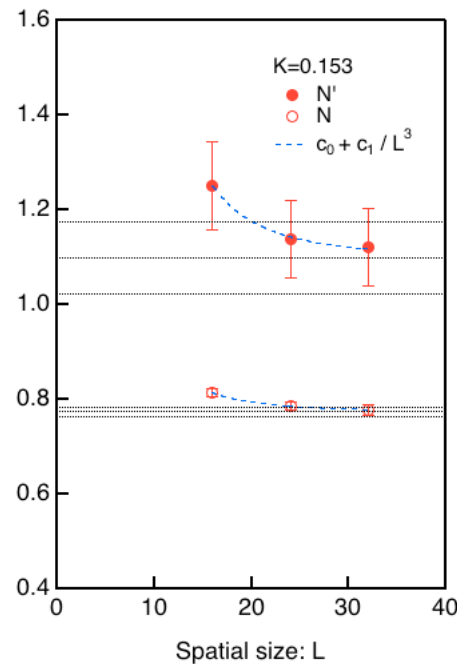
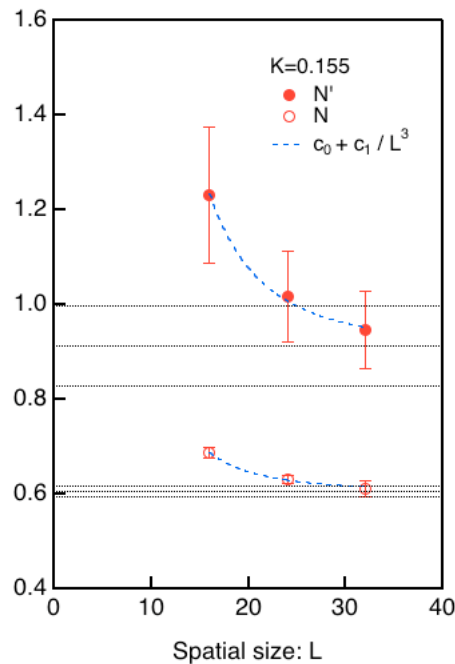


Heavy ($m_\pi \sim 1.1 \text{ GeV}$)



Infinite volume limit

- All data is well fitted by the $1 / L^3$ curve



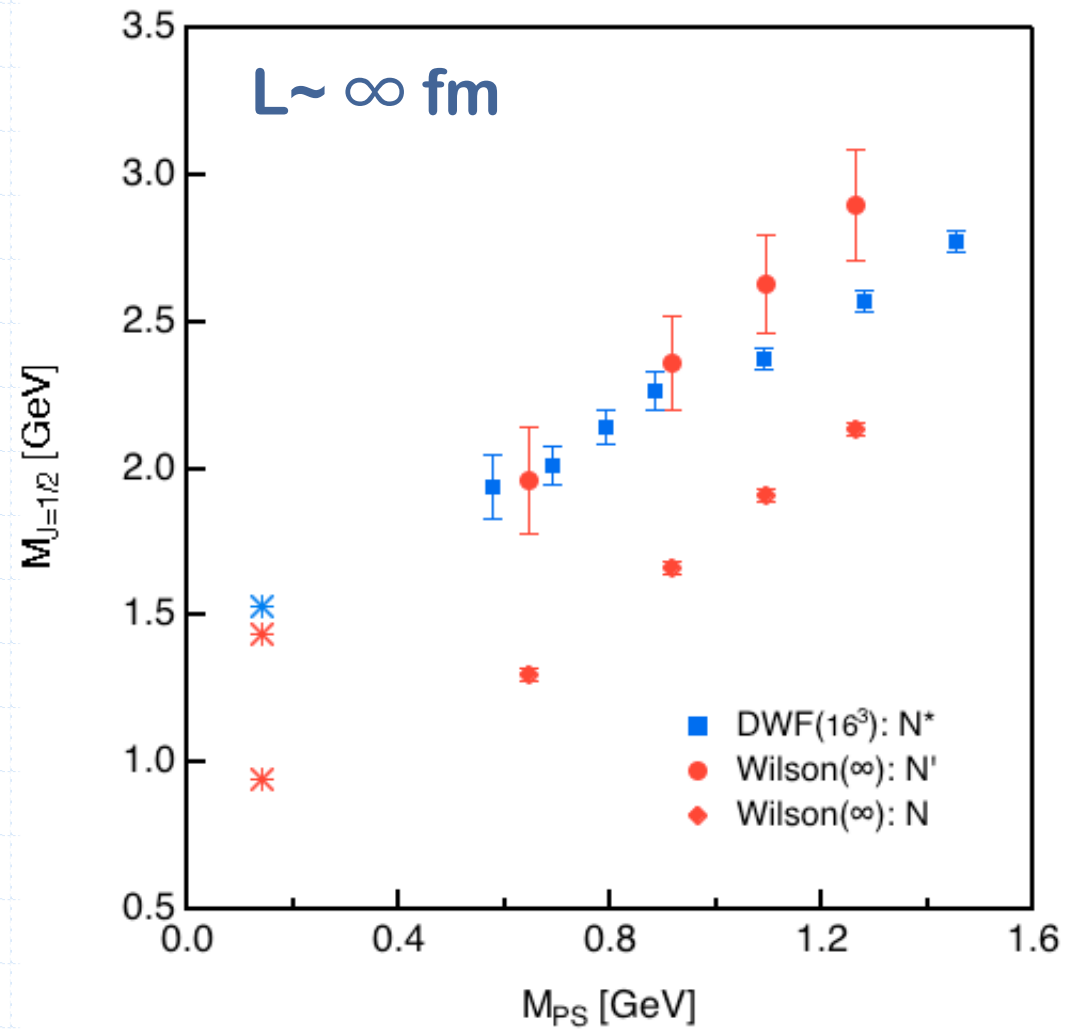
light



heavy




➡ The finite-size effect is rather severe in the lighter quark mass region

Level switching between N' and N^*



Summary

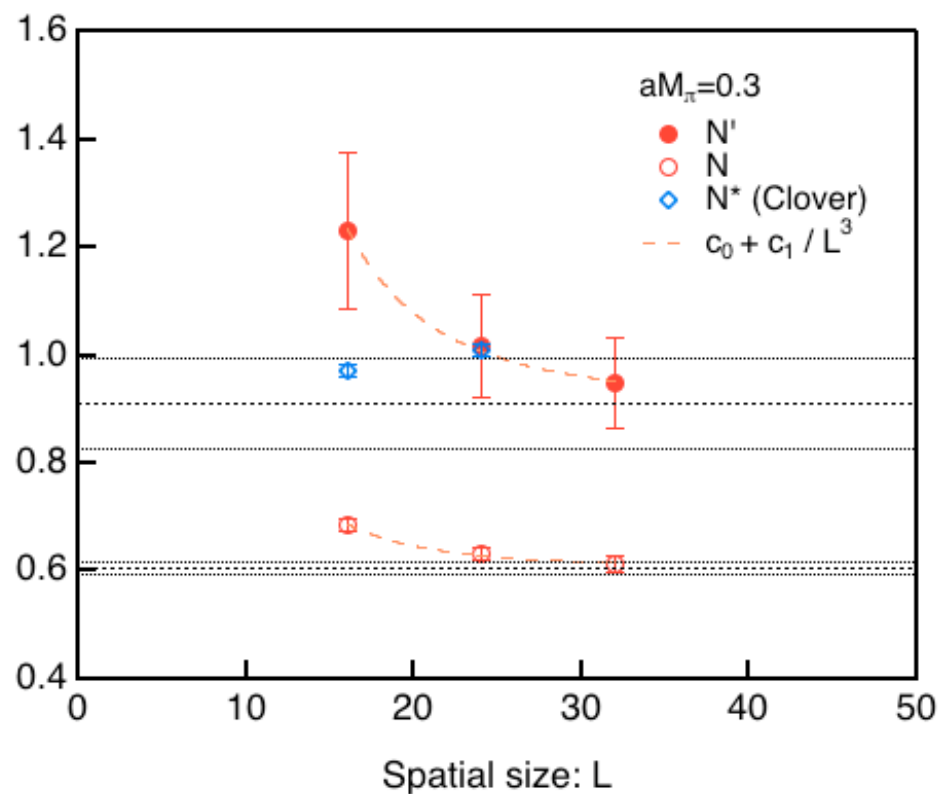
Maximum entropy method is applied to quenched baryon spectrum with Wilson fermions

- ❑ Succeed in extracting N and Δ spectral functions
 - ✓ Can extract the mass of first excited states
 - ✓ Find unphysical states around $\Delta a \sim 2.0$ and 3.0 
- ❑ Confirm the large finite-size effect on the N' state in the light quark mass region
 - ✓ Spatial lattice size $L a \sim 2.2$ fm is not large enough
 - ➡ The level switching between N' and N^* should happen in lattice simulations with large spatial size $L a > 3.0$ fm 
 - ➡ The Roper resonance can be described by the simple three quark excitation of sizable extent 



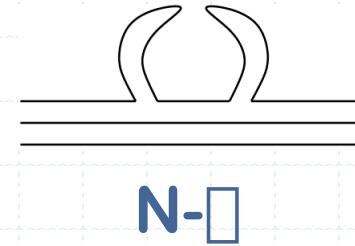
Opposite finite-size effect on N^* to N and N'

The N^* mass on the larger lattice ($L=24$) is higher than on the smaller lattice ($L=16$) by an amount of order 5%. (JLab/UKQCD, 2002)

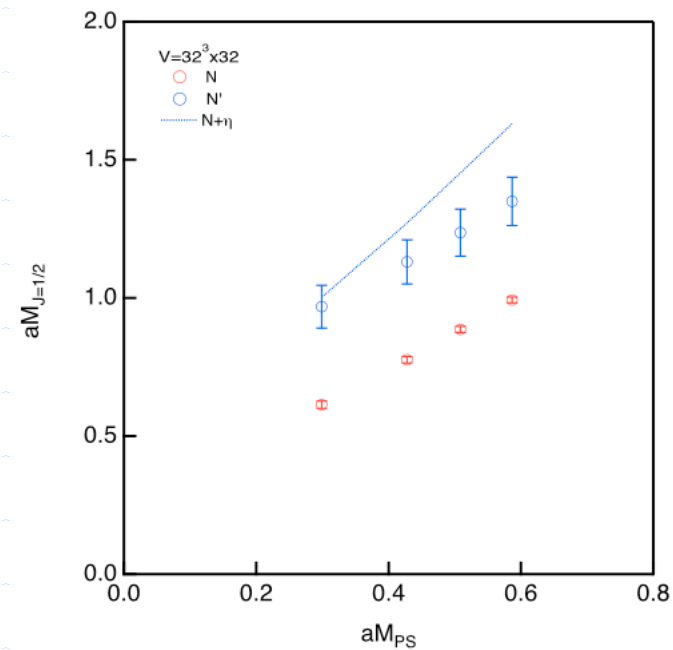
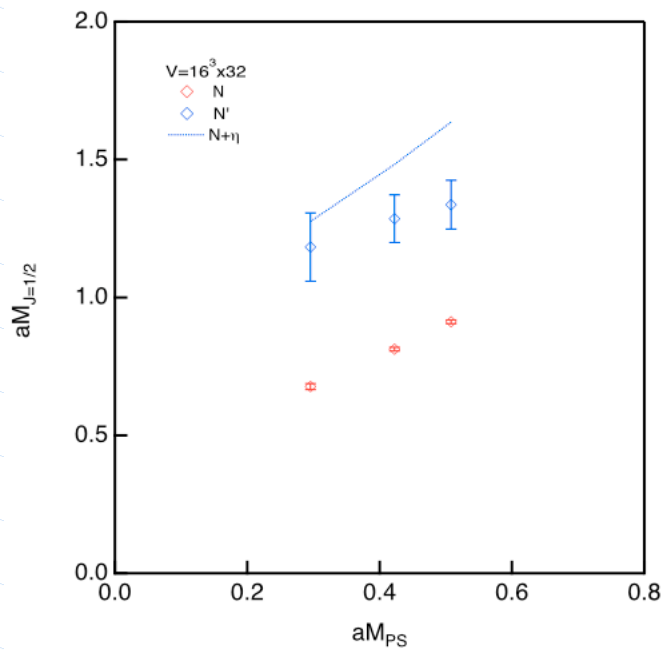


Data for N^* (clover fermions at $\beta=6.0$) is taken from Phys. Lett. B532 (2002) 63

Multi-hadron state???



No, it can't be !



$$\underline{M_{N\Box} < \sqrt{M_N^2 + p_{\min}^2} + \sqrt{M_{\Box}^2 + p_{\min}^2}} \quad \text{where } p_{\min} = \frac{2\Box}{aL} \quad \text{and} \quad M_{\Box} = M_{\Box}$$