

Ground State Baryons in Covariant Level Classification Scheme

Muneyuki Ishida (Tokyo Inst. Tech.)

I. Introduction

II. Spinor-Flavor WF

(Constituent Dirac Spinor)

(Baryon Spinor)

(Composition rule of $|F\rangle, |p\rangle, |s\rangle$ WF)

(WF in general frame)

III. Properties of Ground State Baryons

(Interactions among hadrons)

(Magnetic Moment)

(WF mixing and Parity of Chiral States)

(Mechanism of Pion Emission)

($56^{\oplus} \rightarrow 56^{\oplus} + \pi$)

(Candidates for 56^{\oplus} Chiral States)

(Properties of 56^{\oplus})

($56_E - 56_F$ mixing)

IV. Concluding Remarks

I. Introduction

Covariant Level Classification Scheme of Hadrons presented by Shin Ishida

$$\tilde{U}(12)_{SF} \otimes O(3,1)_L \iff SU(6)_{SF} \otimes O(3)_L$$

NRQM

Hadron WF : Tensors in $\tilde{U}(12)_{SF} \otimes O(3,1)_L$ space

$$\left(\begin{array}{l} \text{Meson } \Phi_A^B(x, y) \sim \psi_A(x) \bar{\psi}^B(y) \\ \text{Baryon } \Phi_{A_1 A_2 A_3}(x_1, x_2, x_3) \sim \psi_{A_1}(x_1) \psi_{A_2}(x_2) \psi_{A_3}(x_3) \end{array} \right. \quad A = (\overset{\text{spinor}}{\alpha}, \overset{\text{flavor}}{a}) \text{ etc.}$$

\implies Existence of

Chiral Particle/States Out of Non-Rel. Scheme

Ground State Baryons

$$\tilde{U}(12)_{SF}$$

$$SU(6)_{SF}$$

$$(12 \times 12 \times 12)_{sym} = 364 \iff (6 \times 6 \times 6)_{sym} = 56$$

$$= 182_B \oplus 182_{\bar{B} \text{ or } "8_S B"}$$

182 ₈	<table style="border-collapse: collapse; width: 100%;"> <tr> <td style="padding: 5px;">56^{\oplus}</td> <td style="padding: 5px;">$\Delta_{\frac{3}{2}}^{\oplus}$</td> <td style="padding: 5px;">$N_{\frac{1}{2}}^{\oplus}$</td> <td style="padding: 5px;"></td> <td style="padding: 5px;">56^{\oplus}</td> <td style="padding: 5px;">$\Delta_{\frac{3}{2}}^{\oplus}$</td> <td style="padding: 5px;">$N_{\frac{1}{2}}^{\oplus}$</td> </tr> <tr> <td style="padding: 5px;">70</td> <td style="padding: 5px;">$\Delta_{\frac{1}{2}}$</td> <td style="padding: 5px;">$N_{\frac{3}{2}}$</td> <td style="padding: 5px;">$N_{\frac{1}{2}}$</td> <td style="padding: 5px;">$\Lambda_{\frac{1}{2}}$</td> <td style="padding: 5px;"></td> <td style="padding: 5px;"></td> </tr> <tr> <td style="padding: 5px;">$56'$</td> <td style="padding: 5px;">$\Delta'_{\frac{3}{2}}$</td> <td style="padding: 5px;">$N'_{\frac{1}{2}}$</td> <td style="padding: 5px;"></td> <td style="padding: 5px;"></td> <td style="padding: 5px;"></td> <td style="padding: 5px;"></td> </tr> </table>	56^{\oplus}	$\Delta_{\frac{3}{2}}^{\oplus}$	$N_{\frac{1}{2}}^{\oplus}$		56^{\oplus}	$\Delta_{\frac{3}{2}}^{\oplus}$	$N_{\frac{1}{2}}^{\oplus}$	70	$\Delta_{\frac{1}{2}}$	$N_{\frac{3}{2}}$	$N_{\frac{1}{2}}$	$\Lambda_{\frac{1}{2}}$			$56'$	$\Delta'_{\frac{3}{2}}$	$N'_{\frac{1}{2}}$					Parity \swarrow
56^{\oplus}	$\Delta_{\frac{3}{2}}^{\oplus}$	$N_{\frac{1}{2}}^{\oplus}$		56^{\oplus}	$\Delta_{\frac{3}{2}}^{\oplus}$	$N_{\frac{1}{2}}^{\oplus}$																	
70	$\Delta_{\frac{1}{2}}$	$N_{\frac{3}{2}}$	$N_{\frac{1}{2}}$	$\Lambda_{\frac{1}{2}}$																			
$56'$	$\Delta'_{\frac{3}{2}}$	$N'_{\frac{1}{2}}$																					

Chiral States : Properties ?

Magnetic Moment

One Pion Emission Decay

II. Spinor-Flavor Wave Function

(Constituent Dirac Spinor)

$$\Psi_{q,\alpha}(X) = \sum_{\mathbb{P}} \left[e^{i\mathbb{P}\cdot X} \underbrace{W_{q,\alpha}^{(+)}(P)}_{u_{\alpha}(P)} + e^{-i\mathbb{P}\cdot X} \underbrace{W_{q,\alpha}^{(-)}(P)}_{u_{\alpha}(-P)} \right]$$

$$\bar{\Psi}_{\bar{q}}^{\alpha}(X) = \sum_{\mathbb{P}} \left[e^{i\mathbb{P}\cdot X} \underbrace{\bar{W}_{\bar{q}}^{(+)\alpha}(P)}_{\bar{v}_{\bar{q}}^{\alpha}(P)} + e^{-i\mathbb{P}\cdot X} \underbrace{\bar{W}_{\bar{q}}^{(-)\alpha}(P)}_{\bar{v}_{\bar{q}}^{\alpha}(-P)} \right]$$

$$\mathbb{P}=0 \quad W_q^{(+)} = \begin{pmatrix} \chi \\ 0 \end{pmatrix} \rho_3 = +1$$

$$W_q^{(-)} = \begin{pmatrix} 0 \\ \chi \end{pmatrix} \rho_3 = -1$$

$$\bar{W}_{\bar{q}}^{(+)} = (0, -\chi^{\dagger}) \bar{\rho}_3 = +1$$

$$\bar{W}_{\bar{q}}^{(-)} = (\chi^{\dagger}, 0) \bar{\rho}_3 = -1$$

(Baryon Spinor)

$$W_{\alpha\beta\gamma}^{(B)}(P) = W_{q,\alpha}(P) W_{q,\beta}(P) W_{q,\gamma}(P)$$

$$\mathbb{P}=0 \quad (\rho_3^{(1)}, \rho_3^{(2)}, \rho_3^{(3)}) = \begin{pmatrix} + & + & + \end{pmatrix}$$

$$\begin{pmatrix} + & + & - \end{pmatrix}$$

$$\begin{pmatrix} + & - & - \end{pmatrix}$$

} chiral states

$$W_{\alpha\beta\gamma}^{(\bar{B})}(P) = W^{(B)} \left\{ W_{q,\alpha} \rightarrow W_{\bar{q},\alpha} \right\}$$

(Composition rule of $|F\rangle, |\rho\rangle, |\sigma\rangle$ WF)

WF is fully-symmetric in $|F\rangle \otimes |\rho\rangle \otimes |\sigma\rangle$ space
flavor Dirac spinor

$$|F_1 F_2\rangle_S = \begin{cases} |F_1\rangle_S |F_2\rangle_S \\ |F_1\rangle_A |F_2\rangle_A \\ \frac{1}{\sqrt{2}} (|F_1\rangle_\alpha |F_2\rangle_\alpha + |F_1\rangle_\beta |F_2\rangle_\beta) \end{cases}$$

$|F\rangle_S \sim uuu, \frac{1}{\sqrt{3}}(uud+udu+duu), \text{etc.} \sim \Delta\text{-decouplet}$

$|F\rangle_A \sim -\frac{1}{\sqrt{6}} \epsilon_{abc} \quad a=u,d,s, \dots \sim \Lambda\text{-singlet}$

$|F\rangle_\alpha \sim -\frac{1}{\sqrt{6}} (-N_b^d \epsilon_{dca} + N_c^d \epsilon_{dab}) : (bc) \sim N\text{-octet}$

$|F\rangle_\beta \sim \frac{1}{\sqrt{2}} N_a^d \epsilon_{dcb} \quad : [bc] \sim N\text{-octet}$

$|\sigma\rangle_S : \text{spin } \frac{3}{2}$

$|\sigma\rangle_{\alpha, \beta} : \text{spin } \frac{1}{2}$

No $|\sigma\rangle_A$

$$N_a^d = \begin{pmatrix} \frac{\Sigma^0 + \Lambda}{\sqrt{2}} & \Sigma^+ & p \\ \Sigma^- & -\frac{\Sigma^0 + \Lambda}{\sqrt{2}} & n \\ \Xi^- & \Xi^0 & -\frac{2}{\sqrt{6}} \Lambda \end{pmatrix}$$

$|\rho, \frac{3}{2}\rangle_S = |+++ \rangle_\rho : \text{NR states}$

$|\rho, -\frac{3}{2}\rangle_S = |--- \rangle_\rho$

$|\rho, -\frac{1}{2}\rangle_S = \frac{1}{\sqrt{3}} (|+-- \rangle_\rho + |-+- \rangle_\rho + |--+ \rangle_\rho)$

$|\rho, \frac{1}{2}\rangle_S = \frac{1}{\sqrt{3}} (|++- \rangle_\rho + |+-+ \rangle_\rho + |-++ \rangle_\rho)$

$|\rho, -\frac{1}{2}\rangle_\alpha = \frac{1}{\sqrt{6}} (-|--+ \rangle_\rho - |-+- \rangle_\rho + 2|+-+ \rangle_\rho)$

$|\rho, \frac{1}{2}\rangle_\alpha = \frac{1}{\sqrt{6}} (|++- \rangle_\rho + |+-+ \rangle_\rho - 2|--+ \rangle_\rho)$

$|\rho, -\frac{1}{2}\rangle_\beta = \frac{1}{\sqrt{2}} (-|--+ \rangle_\rho + |-+- \rangle_\rho)$

$|\rho, \frac{1}{2}\rangle_\beta = \frac{1}{\sqrt{2}} (|++- \rangle_\rho - |+-+ \rangle_\rho)$

Positive Parity

Negative Parity

$$364 = 182 (B) \oplus 182 (\gamma_5 B \sim \bar{B})$$

Positive Parity

Negative Parity

	E^\oplus	E^\ominus	
56 _E	$\Delta_{\frac{3}{2}}$	$ p-\frac{3}{2}\rangle_s \Delta\rangle_s \sigma\rangle_s$	$ p-\frac{3}{2}\rangle_s \Delta\rangle_s \sigma\rangle_s$
	$N_{\frac{1}{2}}$	$ p-\frac{3}{2}\rangle_s \frac{1}{\sqrt{2}}(N\rangle_\alpha \sigma\rangle_\alpha + N\rangle_\beta \sigma\rangle_\beta)$	$ p-\frac{3}{2}\rangle_s \frac{1}{\sqrt{2}}(N\rangle_\alpha \sigma\rangle_\alpha + N\rangle_\beta \sigma\rangle_\beta)$
	F^\oplus	F^\ominus	
56 _F	$\Delta'_{\frac{3}{2}}$	$ p-\frac{1}{2}\rangle_s \Delta\rangle_s \sigma\rangle_s$	$ p-\frac{1}{2}\rangle_s \Delta\rangle_s \sigma\rangle_s$
	$N'_{\frac{1}{2}}$	$ p-\frac{1}{2}\rangle_s \frac{1}{\sqrt{2}}(N\rangle_\alpha \sigma\rangle_\alpha + N\rangle_\beta \sigma\rangle_\beta)$	$ p-\frac{1}{2}\rangle_s \frac{1}{\sqrt{2}}(N\rangle_\alpha \sigma\rangle_\alpha + N\rangle_\beta \sigma\rangle_\beta)$
	G^\oplus	G^\ominus	
70 _G	$\Delta_{\frac{1}{2}}$	$ \Delta\rangle_s \frac{1}{\sqrt{2}}(p-\frac{1}{2}\rangle_\alpha \sigma\rangle_\alpha + p-\frac{1}{2}\rangle_\beta \sigma\rangle_\beta)$	$ \Delta\rangle_s \frac{(-1)}{\sqrt{2}}(p-\frac{1}{2}\rangle_\alpha \sigma\rangle_\alpha + p-\frac{1}{2}\rangle_\beta \sigma\rangle_\beta)$
	$N_{\frac{3}{2}}$	$\frac{1}{\sqrt{2}}(p-\frac{1}{2}\rangle_\alpha N\rangle_\alpha + p-\frac{1}{2}\rangle_\beta N\rangle_\beta) \sigma\rangle_s$	$\frac{(-1)}{\sqrt{2}}(p-\frac{1}{2}\rangle_\alpha N\rangle_\alpha + p-\frac{1}{2}\rangle_\beta N\rangle_\beta) \sigma\rangle_s$
	$N_{\frac{1}{2}G}$	$\frac{1}{2} \left[p-\frac{1}{2}\rangle_\alpha (- N\rangle_\alpha \sigma\rangle_\alpha + N\rangle_\beta \sigma\rangle_\beta) \right. \\ \left. + p-\frac{1}{2}\rangle_\beta (N\rangle_\alpha \sigma\rangle_\beta + N\rangle_\beta \sigma\rangle_\alpha) \right]$	$\frac{(-1)}{2} \left[p-\frac{1}{2}\rangle_\alpha (- N\rangle_\alpha \sigma\rangle_\alpha + N\rangle_\beta \sigma\rangle_\beta) \right. \\ \left. + p-\frac{1}{2}\rangle_\beta (N\rangle_\alpha \sigma\rangle_\beta + N\rangle_\beta \sigma\rangle_\alpha) \right]$
	$\Lambda_{\frac{1}{2}}$	$- \Lambda\rangle_A \frac{1}{\sqrt{2}}(- p-\frac{1}{2}\rangle_\alpha \sigma\rangle_\beta + p-\frac{1}{2}\rangle_\beta \sigma\rangle_\alpha)$	$ \Lambda\rangle_A \frac{(+1)}{\sqrt{2}}(p-\frac{1}{2}\rangle_\alpha \sigma\rangle_\beta + p-\frac{1}{2}\rangle_\beta \sigma\rangle_\alpha)$

WF at rest frame $P=0$: $|F\rho\sigma\rangle = |F\rho\sigma(v_0)\rangle$

$$\frac{P_0}{M} = v_{0\mu} = (000:i)$$

(WF in general frame)

$$|F\rho\sigma(v)\rangle = B(v) |F\rho\sigma\rangle$$

$B(v) = B(v)^{(1)} B(v)^{(2)} B(v)^{(3)}$: Lorentz booster

$$B(v)^{(i)} = \cosh\theta + \rho_i^{(i)} \sigma_z^{(i)} \sinh\theta \quad \cosh\theta = \sqrt{\frac{E+m}{2m}} = \sqrt{\frac{\omega+1}{2}}$$

$$v_\mu = (00\omega_3:i\omega) \quad \sinh\theta = \sqrt{\frac{E-m}{2m}} = \sqrt{\frac{\omega-1}{2}}$$

$$2 \cosh\theta \sinh\theta = \frac{|\rho|}{m}$$

$$\begin{pmatrix} \chi \\ 0 \end{pmatrix} \rightarrow B(v) \begin{pmatrix} \chi \\ 0 \end{pmatrix} = \begin{pmatrix} \cosh\theta \chi \\ \sinh\theta \sigma_z \chi \end{pmatrix} = u(v)$$

(Magnetic Moment)

$$3 \langle F\rho\sigma(\uparrow) | \mu_z^{(1)} \sigma_z^{(1)} | F\rho\sigma(\uparrow) \rangle$$

$$\mu_B = \left(\begin{matrix} \mu_u \\ \mu_d \\ \mu_s \end{matrix} \right) \frac{e\hbar}{2m_u}$$

8

M. Ishida and S. Ishida

NR

chiral

	$\bar{p}p$	$\bar{n}n$	$\Lambda\Lambda$	$\Lambda\Sigma^0$	$\Sigma^+\Sigma^+$	$\Sigma^-\Sigma^-$	$\Sigma^0\Sigma^0$	$\Xi^-\Xi^-$	$\Xi^0\Xi^0$
$ N\rangle_\alpha$	$\frac{\mu_u + 2\mu_d}{3}$	$\frac{2\mu_u + \mu_d}{3}$	$\frac{\mu_u + \mu_d}{2}$	$\frac{\mu_u - \mu_d}{-2\sqrt{3}}$	$\frac{\mu_u + 2\mu_s}{3}$	$\frac{\mu_d + 2\mu_s}{3}$	$\frac{\mu_u + \mu_d + 4\mu_s}{6}$	$\frac{2\mu_d + \mu_s}{3}$	$\frac{2\mu_u + \mu_s}{3}$
$ N\rangle_\beta$	μ_u	μ_d	$\frac{\mu_u + \mu_d + 4\mu_s}{6}$	$\frac{\mu_u - \mu_d}{2\sqrt{3}}$	μ_u	μ_d	$\frac{\mu_u + \mu_d}{2}$	μ_s	μ_s
$ N_{\frac{56}{8}}^{56}\rangle$	$\frac{4\mu_u - \mu_d}{3}$	$\frac{-\mu_u + 4\mu_d}{3}$	μ_s	$\frac{\mu_u - \mu_d}{\sqrt{3}}$	$\frac{4\mu_u - \mu_s}{3}$	$\frac{4\mu_d - \mu_s}{3}$	$\frac{2\mu_u + 2\mu_d - \mu_s}{3}$	$\frac{-\mu_d + 4\mu_s}{3}$	$\frac{-\mu_u + 4\mu_s}{3}$
Theor.	<u>2.793</u>	-1.862	<u>0.613</u>	1.61	2.69	-1.04	0.825	-0.507	-1.44
Exp.	2.793	-1.913	-0.613(4)	1.61(8)	2.46(1)	-1.16(3)		-0.651(3)	-1.25(1)
$ N_{\frac{56}{8}'}^{56}\rangle$	$\frac{4\mu_u - \mu_d}{3}$	$\frac{-\mu_u + 4\mu_d}{3}$	μ_s	$\frac{\mu_u - \mu_d}{\sqrt{3}}$	$\frac{4\mu_u - \mu_s}{3}$	$\frac{4\mu_d - \mu_s}{3}$	$\frac{2\mu_u + 2\mu_d - \mu_s}{3}$	$\frac{-\mu_d + 4\mu_s}{3}$	$\frac{-\mu_u + 4\mu_s}{3}$
Theor.	2.793	-1.862	-0.613	1.61	2.69	-1.04	0.825	-0.507	-1.44
$ N_{\frac{70}{8}}^{70}\rangle$	$\frac{2\mu_u + \mu_d}{3}$	$\frac{\mu_u + 2\mu_d}{3}$	$\frac{\mu_u + \mu_d + \mu_s}{3}$	0	$\frac{2\mu_u + \mu_s}{3}$	$\frac{2\mu_d + \mu_s}{3}$	$\frac{\mu_u + \mu_d + \mu_s}{3}$	$\frac{\mu_d + 2\mu_s}{3}$	$\frac{\mu_u + 2\mu_s}{3}$
Theor.	0.931	0	0.106	0	1.037	-0.825	0.106	-0.719	0.212
$ N_{\frac{70}{8}'}^{70}\rangle$	$2\mu_u + \mu_d$	$\mu_u + 2\mu_d$	$\mu_u + \mu_d + \mu_s$	0	$2\mu_u + \mu_s$	$2\mu_d + \mu_s$	$\mu_u + \mu_d + \mu_s$	$\mu_d + 2\mu_s$	$\mu_u + 2\mu_s$
Theor.	<u>2.793</u>	0	0.318	0	3.11	-2.48	0.318	-2.16	0.636

the same

Table VI. Magnetic moment of nucleon octets: Unit is the nuclear magneton μ_N . $\mu_d = -\mu_u/2$ is assumed. The values with underlines, which are used as inputs, give $\mu_u = 1.862\mu_N$ and $\mu_s = -0.613\mu_N$.

NR

chiral

	$\Delta^{++}\Delta^{++}$	$\Delta^+\Delta^+$	$\Delta^0\Delta^0$	$\Delta^-\Delta^-$	$\Sigma^+\Sigma^+$	$\Sigma^-\Sigma^-$	$\Sigma^0\Sigma^0$	$\Xi^-\Xi^-$	$\Xi^0\Xi^0$	$\Omega^-\Omega^-$
$ \Delta_{\frac{56}{8}E}^{56}\rangle$	$3\mu_u$	$2\mu_u + \mu_d$	$\mu_u + 2\mu_d$	$3\mu_d$	$2\mu_u + \mu_s$	$2\mu_d + \mu_s$	$\mu_u + \mu_d + \mu_s$	$\mu_d + 2\mu_s$	$\mu_u + 2\mu_s$	$3\mu_s$
Theor.	5.59	2.79	0	-2.79	3.11	-2.48	0.318	-2.16	0.636	-1.84
Exp.	3.7 ~ 7.5									
$ \Delta_{\frac{56}{8}F}^{56}\rangle$	$3\mu_u$	$2\mu_u + \mu_d$	$\mu_u + 2\mu_d$	$3\mu_d$	$2\mu_u + \mu_s$	$2\mu_d + \mu_s$	$\mu_u + \mu_d + \mu_s$	$\mu_d + 2\mu_s$	$\mu_u + 2\mu_s$	$3\mu_s$
Theor.	5.59	2.79	0	-2.79	3.11	-2.48	0.318	-2.16	0.636	-1.84
$ \Delta_{\frac{70}{8}G}^{70}\rangle$	μ_u	$\frac{2\mu_u + \mu_d}{3}$	$\frac{\mu_u + 2\mu_d}{3}$	μ_d	$\frac{2\mu_u + \mu_s}{3}$	$\frac{2\mu_d + \mu_s}{3}$	$\frac{\mu_u + \mu_d + \mu_s}{3}$	$\frac{\mu_d + 2\mu_s}{3}$	$\frac{\mu_u + 2\mu_s}{3}$	μ_s
Theor.	1.86	0.931	0	-0.931	1.037	-0.825	0.106	-0.719	0.212	-0.613

the same

Table VII. Magnetic moment of Δ decuplets: Unit is the nuclear magneton μ_N . $\mu_d = -\mu_u/2$ is assumed. The $\mu_u = 1.862\mu_N$ and $\mu_s = -0.613\mu_N$ are used.

the same as the prediction by NRQM $SU(6)_{SF}$

Chiral $56'$ has the same M.M. as 56^\oplus .

70_G different from

Ground State Baryons in Covariant ...

9

Chiral

	$\Lambda^-\Lambda^-$
$ \Lambda_{\frac{70}{8}G}^{70}\rangle$	$\frac{\mu_u + \mu_d + \mu_s}{3}$
Theor.	0.106

Table VIII. Magnetic moment of Λ singlet

(WF mixing and Parity of Chiral States)

$SU(6)_{SF}$ good \leftarrow (Magnetic Moment
F/D ratio of N-octet)

$$56^{\oplus} = c_1 56_E^{\oplus} + s_1 56_F^{\oplus} \quad |c_1|^2 + |s_1|^2 = 1$$

$$56' = -s_1^* 56_E + c_1^* 56_F \quad \left. \vphantom{56'} \right\} \text{chiral states}$$

$$70 = 70_G$$

Candidates for Chiral States

$$N^{\oplus}(1440) \quad \Lambda^{\ominus}(1405)$$

\implies Two Possibilities

(Possibility A)

$$\begin{array}{cc} 56^{\oplus} & , & 70^{\ominus} \\ \text{U} & & \text{U} \\ N(1440) & & \Lambda(1405) \\ & & \text{singlet} \\ & & \text{~~octet~~} \end{array}$$



Promising!

(Possibility B)

$$\begin{array}{cc} 56^{\ominus} & , & 70^{\oplus} \\ \text{U} & & \text{U} \\ \Lambda(1405) & & N(1440) \\ \text{octet} & & \end{array}$$

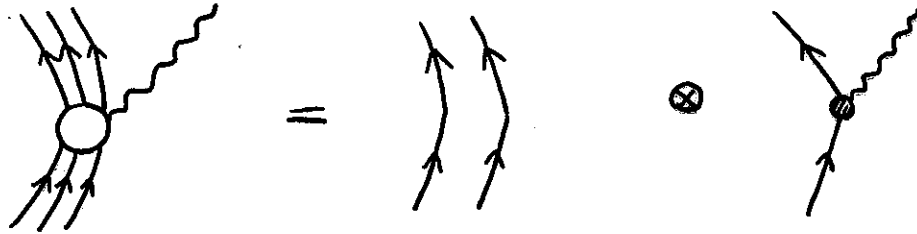
difficulties:

- Negative parity Octet in low mass, but No N^{\ominus} below 1.5 GeV in exp.
- $\Gamma_{\Lambda(1405) \rightarrow \Sigma\pi}^{\text{theor.}} = 0$ but $\Gamma^{\text{exp}} = 50 \text{ MeV}$

III. Properties of Ground State Baryons

(Interactions among hadrons) by Shin Ishida

Interaction vertex = ^(Spectator) Overlapping interaction \otimes Elementary quark interaction



$$\bar{B}^{A_1 A_2 D}(P_2) B_{A_1 A_2 C}(P_1) \bar{M}_D^C(q)$$

$$= (\bar{\Psi}^{A_1} \Psi_{A_1}) (\bar{\Psi}^{A_2} \Psi_{A_2}) \otimes \bar{\Psi}^D \Psi_C \bar{M}_D^C(q)$$

Non-symmetric

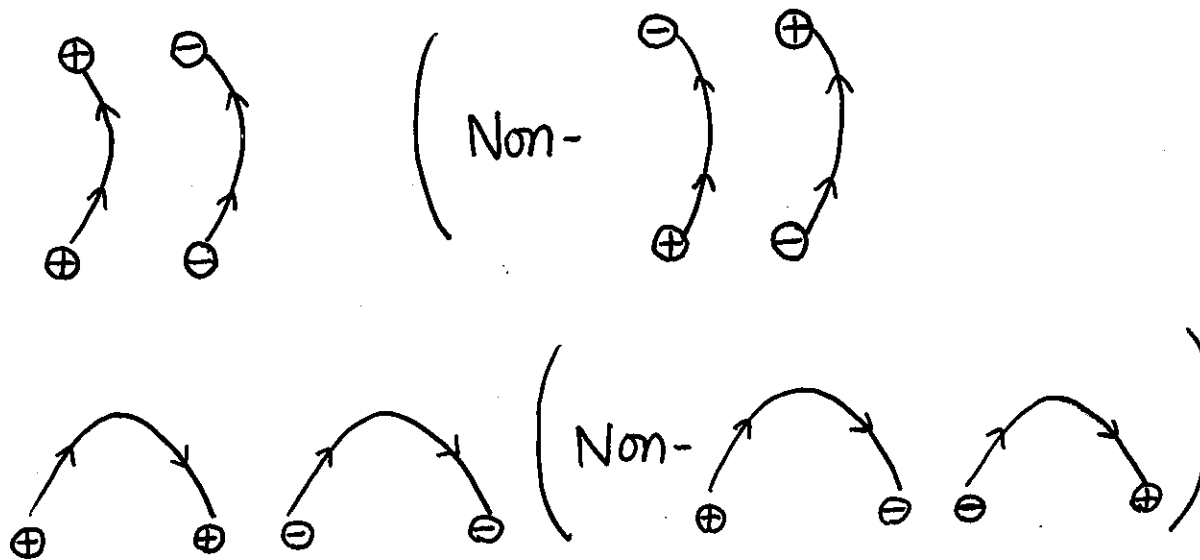
Chiral symmetric

$$(\bar{\Psi}^{A_1} \Psi_{A_1}) = \bar{\Psi}(v_2) \sum_{\lambda}^{(0)} \Psi(v_1) ; \quad \gamma = \frac{-i\hat{v}_2 \cdot \gamma - i\hat{v}_1 \cdot \gamma}{2}$$

$$\hat{v}_\mu = \begin{cases} v_\mu & (\text{for } +) \\ -v_\mu & (\text{for } -) \end{cases}$$

\Rightarrow ρ_3 - line rule

ρ_3 at respective rest frame
conserves on spectator quark line



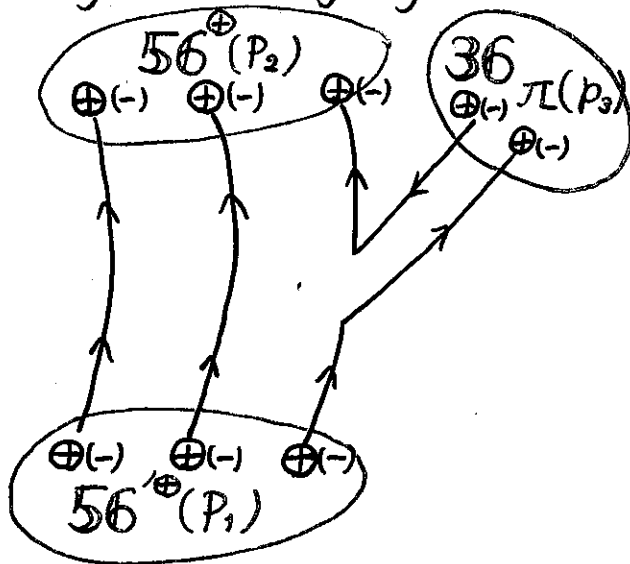
(Mechanism of Pion Emission)

Strong Interaction : is presumably

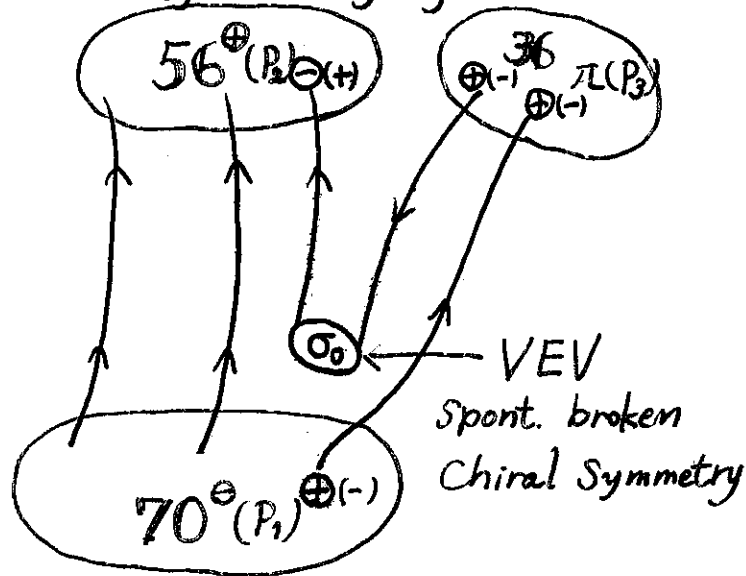
spontaneously broken
Chiral Symmetric Overlapping Interactions

⇒ Concrete form of effective One Pion Emission vertex derived!

Parity - Unchanging



Parity - Changing



$$3g \langle F\rho\sigma(v_2) | (\gamma^0 i \gamma_5 \phi \gamma^0)^{(1)} (\gamma^0)^{(2)} (\gamma^0)^{(3)} | F\rho\sigma(v_1) \rangle$$

$$= 3g \langle F\rho\sigma | (\sigma_z p_3 i \phi)^{(1)} | F\rho\sigma \rangle$$

$$\times \text{ch}^2 \theta_2 \text{sh}(\theta_2 + \theta_3) \text{ch} \theta_3$$

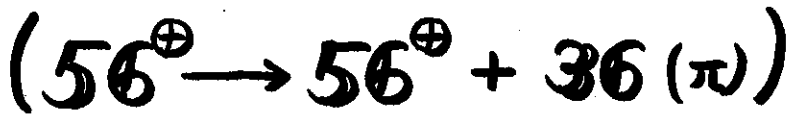
P-wave decay

$$3g \langle F\rho\sigma(v_2) | (\gamma^0 s \gamma^0 i \gamma_5 \phi \gamma^0)^{(1)} (\gamma^0)^{(2)} (\gamma^0)^{(3)} | F\rho\sigma \rangle$$

$$= 3g \langle F\rho\sigma | (i \gamma_5 s \phi)^{(1)} | F\rho\sigma \rangle$$

$$\times \text{ch}^3 \theta_2 \text{ch}^2 \theta_3$$

S-wave decay



πNN -vertex

$$56 = c_1 56_E + s_1 56_F$$

$$\left(\begin{aligned} \mathcal{M} &= 3ig \left(|c_1|^2 - \frac{1}{3} |s_1|^2 \right) s_x \left[-\frac{1}{3} \alpha + \beta \right] \frac{ch^2 \theta_2 \, sh(\theta_2 + \theta_3) \, ch \theta_3}{\left(\frac{ch^2 \theta_2}{2m_\pi} + \frac{ch \theta_2 \, ch^2 \theta_3}{2M_2} \right) \parallel p} \\ \alpha &= \langle \bar{N} | \phi^{(+) | N \rangle_\alpha} \\ \beta &= \langle \bar{N} | \phi^{(+) | N \rangle_\beta} \end{aligned} \right) \Rightarrow \frac{5\sqrt{2}}{9} \bar{\psi} \cdot \pi \psi \left(\frac{ch^2 \theta_2}{2m_\pi} + \frac{ch \theta_2 \, ch^2 \theta_3}{2M_2} \right) \parallel p$$

$$\mathcal{L}_{\pi NN}^{\text{exp}} = g_{\pi NN} \bar{u}(p_2) i \gamma_5 u(p_1) \bar{\psi} \cdot \pi \psi \frac{\parallel p}{2M_2}$$

$$= i 2s_x g_{\pi NN} \bar{\psi} \cdot \pi \psi \, sh \theta_2 \frac{\parallel p}{2M_2}$$

$$\Rightarrow g_{\pi NN} = 13.45 = \underbrace{g \left(|c_1|^2 - \frac{1}{3} |s_1|^2 \right)}_{1.462 : \text{common coupling constant}} \frac{5}{3\sqrt{2}} \left(\frac{M_N}{m_\pi} + 1 \right)$$

$$\Gamma_{B' \rightarrow B\pi} = \boxed{f^2} \frac{\parallel p M_B}{2\pi M_{B'}} \left| g \left(|c_1|^2 - \frac{1}{3} |s_1|^2 \right) \right|^2 \left(ch^2 \theta_2 \, sh(\theta_2 + \theta_3) \, ch \theta_3 \right)^2$$

	f	$\Gamma_{\text{theor}} / \text{MeV}$	$\Gamma_{\text{exp.}} / \text{MeV}$
$\Delta \rightarrow N\pi$	$\sqrt{\frac{4}{3}}$	81.7	115 ~ 125
$\Sigma^*(1385) \rightarrow \Lambda\pi$	$\sqrt{\frac{2}{3}}$	30.7	33
$\rightarrow \Sigma\pi$	$\frac{2}{3}$	4.4	4
$\Xi^*(1530) \rightarrow \Xi\pi$	$\sqrt{\frac{2}{3}}$	11.6	9.1 ± 0.5

\uparrow
 $SU(6)_{SF}$

the same as the predictions by NRQM.

(Candidates for 56^{\oplus} chiral states) PDG'02

$N^{\oplus}(939)$	$\Delta^{\oplus}(1232)$	$\Lambda^{\oplus}(1116)$	$\Sigma^{\oplus}_{\frac{1}{2}}(1191)$	$\Xi^{\oplus}_{\frac{1}{2}}(1318)$
$N^{\oplus}_{\frac{1}{2}}(1440)$	$\Delta^{\oplus}_{\frac{3}{2}}(1600)$	$\Lambda^{\ominus}(1405)$	$\Sigma^{*\oplus}_{\frac{3}{2}}(1385)$	$\Xi^{*\oplus}_{\frac{3}{2}}(1530)$
$N^{\ominus}_{\frac{3}{2}}(1520)$	$\Delta^{\ominus}_{\frac{1}{2}}(1620)$	$\Lambda^{\ominus}_{\frac{3}{2}}(1520)$	$\Sigma^{\oplus}_{\frac{1}{2}}(1660)$	$\Xi(1690)$
$N^{\ominus}_{\frac{1}{2}}(1535)$	$\Delta^{\ominus}_{\frac{1}{2}}(1700)$	$\Lambda^{\oplus}_{\frac{1}{2}}(1600)$	$\Sigma^{\ominus}_{\frac{3}{2}}(1670)$	$\Xi^{\ominus}_{\frac{3}{2}}(1820)$
$N^{\ominus}_{\frac{1}{2}}(1650)$	$(\Delta^{\oplus}_{\frac{1}{2}}(1750))$ $(\Delta^{\ominus}_{\frac{1}{2}}(1900))$	$\Lambda^{\ominus}_{\frac{1}{2}}(1670)$	$\Sigma^{\ominus}_{\frac{1}{2}}(1750)$	$\Xi(1950)$
$N^{\ominus}_{\frac{3}{2}}(1675)$	$\Delta^{\oplus}_{\frac{5}{2}}(1905)$	$\Lambda^{\ominus}_{\frac{3}{2}}(1690)$	$\Sigma^{\ominus}_{\frac{5}{2}}(1775)$	$\Xi(2030)$
$N^{\oplus}_{\frac{5}{2}}(1680)$	$\Delta^{\oplus}_{\frac{1}{2}}(1910)$	$\Lambda^{\ominus}_{\frac{1}{2}}(1800)$	$(\Sigma^{\ominus}_{\frac{3}{2}}(1580))$ $(\Sigma^{\ominus}_{\frac{1}{2}}(1620))$	
$N^{\ominus}_{\frac{3}{2}}(1700)$	$\Delta^{\oplus}_{\frac{3}{2}}(1920)$	$\Lambda^{\oplus}_{\frac{1}{2}}(1810)$		$\Omega^{\oplus}_{\frac{3}{2}}(1672)$
$N^{\oplus}_{\frac{1}{2}}(1710)$	$\Delta^{\oplus}_{\frac{5}{2}}(1930)$ $(\Delta^{\oplus}_{\frac{3}{2}}(1940))$	$\Lambda^{\oplus}_{\frac{5}{2}}(1820)$	$\Sigma^{\oplus}_{\frac{5}{2}}(1915)$	$\Omega(2250)$
$N^{\oplus}_{\frac{3}{2}}(1720)$	$\Delta^{\oplus}_{\frac{7}{2}}(1950)$	$\Lambda^{\oplus}_{\frac{7}{2}}(1830)$	$\Sigma^{\oplus}_{\frac{3}{2}}(1940)$	
$N^{\ominus}_{\frac{7}{2}}(2190)$		$\Lambda^{\oplus}_{\frac{3}{2}}(1890)$ $(\Lambda^{\oplus}_{\frac{7}{2}}(2020))$ $\Lambda^{\ominus}_{\frac{7}{2}}(2100)$	$\Sigma^{\oplus}_{\frac{7}{2}}(2030)$	

Mass formula for 56^{\oplus} chiral states

$$1.44^2 - 0.939^2 \approx (1.6^2 - 1.232^2) \approx (1.6^2 - 1.116^2) \approx (1.66^2 - 1.191^2)$$

$$= 1.19 \quad (= 1.04) \quad (= 1.32) \quad (= 1.337)$$

$$\ll 2\Omega = 2.3 \text{ GeV}^2$$

$$(\Omega = 1.15 \text{ GeV}^2 \leftarrow (\text{Regge slope})^{-1})$$

(Properties of 56^{\oplus} Chiral States)

$$56^{\oplus}(B') \rightarrow 56^{\oplus}(B) + 3F(M)$$

$$\Gamma_{B' \rightarrow B\pi} = [f^2] \frac{19M_B}{2\pi M_{B'}} |g c_i^* s_i^*|^2 (ch^2 \theta_2 \overset{\text{P-wave decay}}{sh(\theta_2 + \theta_3)} ch \theta_3)^2$$

B'	BM	f	$\Gamma_{\text{theor}} / \text{MeV}$	$\Gamma_{\text{exp.}} / \text{MeV}$
$N^{\oplus}(1440)$	$N\pi$	$\frac{10\sqrt{2}}{3\sqrt{3}}$	227.5	~ 227.5 ((250~450) x (60~70)%)
	$\Delta\pi$	$\frac{16}{3\sqrt{3}}$	13	~ 87.5 (x (20~30)%)
$\Lambda(1600)$ (1560 ~ 1700)	$\Sigma\pi$	$\frac{4\sqrt{2}}{3}$	61.8 (45~120)	~ 50 ((50~250) x (10~60)%)
	$\Sigma^*\pi$	$-\frac{8}{3}$	11.4 (3~54)	
	NK	$-\frac{4}{\sqrt{3}}$	10.3 (7~23)	
$\Sigma(1660)$ (1630 ~ 1690)	$\Lambda\pi$	$\frac{4\sqrt{2}}{3\sqrt{3}}$	48.1 (41~56)	~ 20 ((40~200) x (10~30)%)
	$\Sigma\pi$	$\frac{16}{9}$	83.5 (68~101)	
	$\Sigma^*\pi$	$-\frac{8\sqrt{2}}{9}$	7.3 (5~11)	
	NK	$\frac{4}{9}$	0.6 (0.5~0.8)	
$\Xi(1710)$ (1810)	$\Xi\pi$	$\frac{2\sqrt{2}}{3\sqrt{3}}$	4.6 (~9)	
	ΛK	$\frac{2\sqrt{2}}{3\sqrt{3}}$	0.3 (~0.8)	
	ΣK	$\frac{10\sqrt{2}}{3\sqrt{3}}$	0.7 (~9)	
	$\Xi^*\pi$	$-\frac{8}{3\sqrt{3}}$	1.3 (~12)	
$\Delta(1600)$ (1550 ~ 1700)	$N\pi$	$-\frac{8}{3\sqrt{3}}$	163 (130~244)	~ 60 ((250~450) x (10~25)%)
	$\Delta\pi$	$\frac{10\sqrt{2}}{3\sqrt{3}}$	92.5 (57~196)	
$\Sigma^*(1760)$ (~1860)	$\Sigma\pi$	$\frac{8}{9}$	36.9 (~59)	
	$\Lambda\pi$	$\frac{4\sqrt{2}}{3\sqrt{3}}$	78.2 (~118)	
	NK	$-\frac{8}{9}$	4.8 (~7.9)	
	$\Sigma^*\pi$	$\frac{8\sqrt{2}}{9}$	53.5 (~112)	
	ΔK	$\frac{2\sqrt{2}}{9}$	0.6 (~5.5)	
$\Xi^*(1880)$ (~1980)	$\Xi\pi$	$\frac{4\sqrt{2}}{3\sqrt{3}}$	54.6 (~87)	
	ΛK	$\frac{4\sqrt{2}}{3\sqrt{3}}$	5.3 (~9.2)	
	ΣK	$\frac{4\sqrt{2}}{3\sqrt{3}}$	3.1 (~6.2)	
	$\Xi^*\pi$	$\frac{2\sqrt{2}}{3\sqrt{3}}$	15.9 (~26)	
$\Omega(2000)$ (~2100)	ΞK	$\frac{8\sqrt{2}}{3\sqrt{3}}$	11.7 (~24)	
	$\Xi^* K$	$\frac{4\sqrt{2}}{3\sqrt{3}}$	— (~3.3)	

56E - 56F Mixing)

$$|c_1|^2 + |s_1|^2 = 1$$

$$56 = c_1 56_E + s_1 56_F \quad \supset \quad N(939), \Delta(1232)$$

$$56' = -s_1^* 56_E + c_1^* 56_F \quad \supset \quad N(1440), \Delta(1600)$$

$$\begin{cases} g_{\pi NN} = 13.45 & \longrightarrow g (|c_1|^2 - \frac{1}{3}|s_1|^2) = 1.462 \\ \Gamma_{N(1440) \rightarrow N\pi} \approx 227.5 \text{ MeV} & \longrightarrow |g c_1^* s_1^*| = 0.4383 \end{cases}$$

(Two Solutions)

(A1) $N \sim N_E$ $N(1440) \sim N_F$

$$|c_1|^2 = 0.922 \gg |s_1|^2 = 0.0782 \quad g = 1.63$$

(A2) $N \sim N_F$ $N(1440) \sim N_E$

$$|c_1|^2 = 0.00934 \ll |s_1|^2 = 0.9907 \quad g = 4.56$$

$$\Delta(1600) \rightarrow N(1440)\pi$$

$$\Gamma = \left(\sqrt{\frac{4}{3}}\right)^2 \frac{1^p M_{1440}}{2\pi M_{1600}} |g (|s_1|^2 - \frac{1}{3}|c_1|^2)|^2 (ch^2\theta_2 sh(\theta_2 + \theta_3) ch\theta_3)^2$$

B'	Bπ	f	Γ _{theor} / MeV	Γ _{exp.} / MeV
Δ(1600) (1550~1700)	Nπ	$-\frac{8}{3\sqrt{3}}$	163 (130~244)	~ 60 ((250~450) × (10~25)%)
	Δπ	$\frac{10\sqrt{2}}{3\sqrt{3}}$	92.5 (57~196)	~ 190 (× (40~70)%)
	N(1440)π	$\sqrt{\frac{4}{3}}$	(A1) 0.19 (-~3.6) (A2) 27 (-~521)	~ 80 (× (10~35)%)

• Small mixing of $56_E \leftrightarrow 56_F$

• Sizable $\Delta(1600) \rightarrow N(1440)\pi$ seems to suggest (A2).

$$\begin{pmatrix} N(939), \Delta(1232) \sim |p - \frac{1}{2}\rangle_s & \text{chiral states} \\ N(1440), \Delta(1600) \sim |p \frac{3}{2}\rangle_s & \text{NR states} \end{pmatrix}$$

IV. Concluding Remarks

- Covariant Level Scheme based on

$$\tilde{U}(12)_{SF} \otimes O(3,1)_L$$

Ground-State B and \bar{B} are classified as 364 ,

$$56^{\oplus}$$

56^{\oplus}	$N(1440)$	$\Lambda(1600)$	$\Sigma(1660)$	$;\Delta(1600)$
70^{\ominus}	$\Lambda(1405)$			

Chiral States

- Mass formula

$$M_{N(1440)}^2 - M_{N(939)}^2 \approx M_{\Lambda(1600)}^2 - M_{\Lambda(1116)}^2 \approx M_{\Sigma(1660)}^2 - M_{\Sigma(1191)}^2 \\ \approx M_{\Delta(1600)}^2 - M_{\Delta(1232)}^2$$

- Magnetic M. of 56^{\oplus} is the same as 56^{\oplus}

- Decay properties of 56^{\oplus} predicted by $SU(6)_{SF}$ seems to be consistent with exp. data.

⇒ Covariant Level Scheme : Promising!

More accurate exp. data Necessary for confirmation.

- Property of 70_G^{\ominus} :

Mixing with ordinary P-wave states 70_1^{\ominus}

- Excited States of Chiral States.

candidates $\Lambda_{\frac{5}{2}}^{\ominus}(1830)?$, $N_{\frac{5}{2}}^{\oplus}(1680), \dots$

(Assignment of Baryons) in $U(12)_{SF} \times O(3,1)_L$

$N=0$

56_0^+	$N(939) \Lambda(1116) \Sigma(1192) \Xi(1318)$
$\Delta_{\frac{3}{2}}^+$	$\Delta(1232) \Sigma^*(1385) \Sigma^*(1385) \Omega(1672)$

56_{G1}^+	$N(1440) \Lambda(1600) \Sigma(1660)$
$\Delta_{\frac{3}{2}}^+$	$\Delta(1600)$

$N=1$

70_1^+	$\Lambda(1520)$
$N_{\frac{1}{2}}^+$	$N(1535) \Lambda_{\frac{1}{2}}^+(1670) \Sigma_{\frac{1}{2}}^+(1750)$
$N_{\frac{3}{2}}^+$	$N(1520) \Lambda_{\frac{3}{2}}^+(1670) \Sigma_{\frac{3}{2}}^+(1670) \Xi_{\frac{3}{2}}^+(1820)$
$N_{\frac{5}{2}}^+$	$N(1650) \Lambda_{\frac{5}{2}}^+(1800)?$
$N_{\frac{7}{2}}^+$	$N(1700) \Sigma_{\frac{7}{2}}^+(1940)$
$N_{\frac{9}{2}}^+$	$N(1675) \Lambda_{\frac{9}{2}}^+(1830)? \Sigma_{\frac{9}{2}}^+(1775)$
$\Delta_{\frac{1}{2}}^+$	$\Delta(1420)$
$\Delta_{\frac{3}{2}}^+$	$\Delta(1700)$

70_{G1}^+	$\Lambda(1405)?$
$N_{\frac{1}{2}}^+$	$(\Sigma_{\frac{1}{2}}^+(1580))$
$\Delta_{\frac{1}{2}}^+$	$(\Sigma_{\frac{1}{2}}^+(1620))$

70_1^+	$\Lambda_{\frac{1}{2}}^+(1830)?$
70_{G1}^+	$N_{\frac{1}{2}}^+(1710) N_{\frac{3}{2}}^+(1720) N_{\frac{5}{2}}^+(1680)$
	$\Delta_{\frac{1}{2}}^+(1750) \Lambda_{\frac{1}{2}}^+(1810)$
56_{G1}^+	
20_{G1}^+	

$N=2$

56_2^+	$N_{\frac{1}{2}}^+(1810) \Lambda_{\frac{1}{2}}^+(1890)$
$N_{\frac{3}{2}}^+$	$\Lambda_{\frac{3}{2}}^+(1820) \Sigma_{\frac{3}{2}}^+(1915) \Xi_{\frac{3}{2}}^+(2030)$
$\Delta_{\frac{1}{2}}^+$	$\Delta_{\frac{1}{2}}^+(1910)$
$\Delta_{\frac{3}{2}}^+$	$\Delta_{\frac{3}{2}}^+(1920)$
$\Delta_{\frac{5}{2}}^+$	$\Delta_{\frac{5}{2}}^+(1905)$
$\Delta_{\frac{7}{2}}^+$	$\Delta_{\frac{7}{2}}^+(1950) \Sigma_{\frac{7}{2}}^+(2030)$

56_0^+	
70_2^+	$(\Lambda_{\frac{3}{2}}^+(2010))$
70_0^+	
20_1^+	

$N=3$

$\Delta_{\frac{5}{2}}^+$	(1930)	$\Lambda_{\frac{7}{2}}^+(2100)$
$(\Delta_{\frac{3}{2}}^+)$	(1940)	
$(\Delta_{\frac{1}{2}}^+)$	(1900)	



Chiral States