

# Scalar Mesons in Lattice QCD Calculations

Scalar Collaboration

International Symposium on

Hadron Spectroscopy, Chiral Symmetry and  
Relativistic Description of Bound Systems

24-26 Feb. 2003 Nihon University

# SCALAR Collaboration

- T. Kunihiro, YITP, Kyoto Univ.
- S. Muroya, Tokuyama Women's Coll.
- A. Nakamura, RIISE, Hiroshima Univ.
- C. Nonaka, Dept. Phys., Duke Univ.
- M. Sekiguchi, Fac. of Eng. Kokushikan Univ.
- H. Wada, Fac. of Sci. and Eng., Nihon Univ.

# Objective of Scalar Collaboration

- ◆ Confidence level of Sigma Meson (and other scalar meson,  $\kappa$ ) has been increasing.
- ◆ Using Lattice QCD, we have been (and will be) addressing the following Question about scalar mesons:

**Are you a Pole in QCD ?**

- ◆ We study its features in QCD.

# Lattice QCD Calculation

- ◆ Relativistic Formulation
  - Quarks are described by Dirac Fermions
- ◆ Not a Model
- ◆ Apart from numerical limitations, there is no approximation.
- ◆ No bound state Calculation
  - No potential
  - No B-S
  - It measures the mass gap in a given channel.

# Lattice QCD Calculation (cont'd)

## ◆ Euclidean Path Integral

$$Z = \int dU d\bar{\psi} d\psi e^{-\bar{\psi} D\psi - S_G}$$

$$G(x, y) = \frac{1}{Z} \int dU d\bar{\psi} d\psi H(y) H^\dagger(x) e^{-\bar{\psi} D\psi - S_G}$$

$$\longrightarrow e^{-m|x-y|}$$

$H(x)$ : Hadron Operator. For  $\sigma$  meson  $H(x) = \bar{\psi}(x)\psi(x)$

$H^\dagger(x) |0\rangle$  State with Quantum Numbers specified by H

# You should trust Lattice QCD

because it is the First Principle  
Calculation!

# You should not trust Lattice QCD

until the following conditions are satisfied:

- **Enough Statistic**

- ◆ Gauge configurations are generated by Monte Carlo, and there are statistical errors like Experiments.

- **Continuum Limit**

- ◆ Lattice spacing

- **Infinite Volume Limit**  $a \rightarrow 0$

- ◆ Lattice Volume is large enough to include hadron.

- **Chiral Extrapolation**  $(La)^4$

- ◆ u and d quark masses on the lattice are large, and extrapolated to zero.

# Operator for $\sigma$ Meson

◆  $I=0$ , scalar

$$\sigma(x) \equiv \sum_{c=1}^3 \bar{\psi}_c(x) \psi^c(x)$$

$$= \sum_{c=1}^3 \sum_{\alpha=1}^4 \frac{\bar{u}_\alpha^c(x) u_\alpha^c(x) + \bar{d}_\alpha^c(x) d_\alpha^c(x)}{\sqrt{2}}$$

$\left( \begin{array}{l} c = 1, 2, 3 \quad \dots \text{ color} \\ \alpha = 1, 2, 3, 4 \quad \dots \text{ Dirac spin} \end{array} \right)$

$$\left( \psi \equiv \begin{pmatrix} u \\ d \end{pmatrix} \right)$$



# Propagator

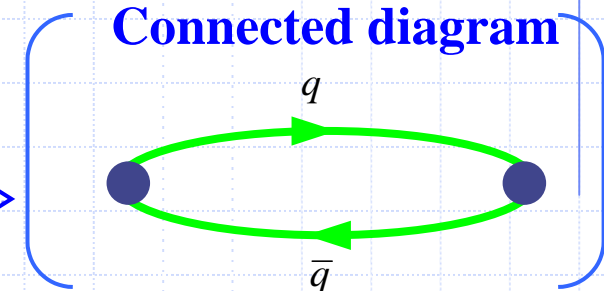
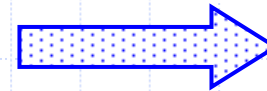
$$G(y, x) = \langle \sigma(y) \sigma(x)^\dagger \rangle$$

$$\begin{aligned}
 &= \frac{1}{Z} \int DUD\bar{u}DuD\bar{d}Dd \sum_{a,b=1}^3 \sum_{\alpha,\beta=1}^4 \frac{\bar{u}_\beta^b(y)u_\beta^b(y) + \bar{d}_\beta^b(y)d_\beta^b(y)}{\sqrt{2}} \\
 &\quad \left( \frac{\bar{u}_\alpha^a(x)u_\alpha^a(x) + \bar{d}_\alpha^a(x)d_\alpha^a(x)}{\sqrt{2}} \right)^\dagger e^{-S_g - \bar{u}Wu - \bar{d}Wd} \\
 &= \frac{1}{Z} \int DUD\bar{u}DuD\bar{d}Dd \sum_{a,b=1}^3 \sum_{\alpha,\beta=1}^4 \frac{1}{2} \left[ \bar{u}_\beta^b(y)u_\beta^b(y)\bar{u}_\alpha^a(x)u_\alpha^a(x) \right. \\
 &\quad + \bar{d}_\beta^b(y)d_\beta^b(y)\bar{d}_\alpha^a(x)d_\alpha^a(x) + \bar{u}_\beta^b(y)u_\beta^b(y)\bar{d}_\alpha^a(x)d_\alpha^a(x) \\
 &\quad \left. + \bar{d}_\beta^b(y)d_\beta^b(y)\bar{u}_\alpha^a(x)u_\alpha^a(x) \right] e^{-S_g - \bar{u}Wu - \bar{d}Wd}
 \end{aligned}$$

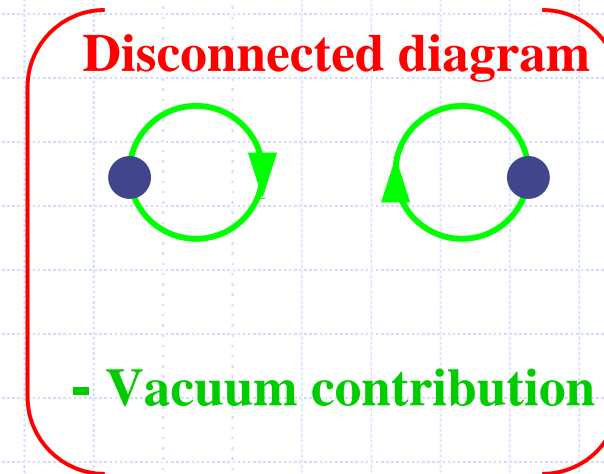
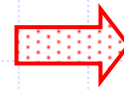
# Propagator (cont'd)

$$G(x, y)$$

$$= - \left\langle \text{Tr} \left( W^{-1}(x, y) W^{-1}(y, x) \right) \right\rangle$$



$$+ 2 \left\langle \text{Tr} \left( W^{-1}(y, y) \right) \text{Tr} \left( W^{-1}(x, x) \right) \right\rangle$$



$$- 2 \left\langle \text{Tr} \left( W^{-1}(y, y) \right) \right\rangle \left\langle \text{Tr} \left( W^{-1}(x, x) \right) \right\rangle$$



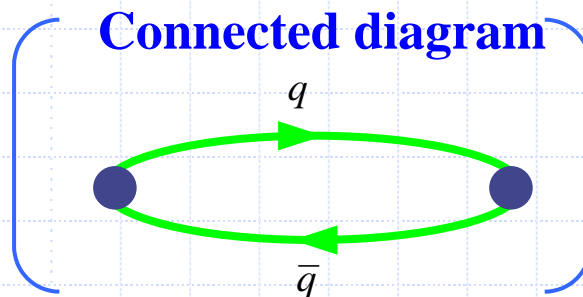
$W^{-1}(x, y)$ : Inverse of Fermion Matrix, i.e., Quark Propagators

# Propagator (cont'd)

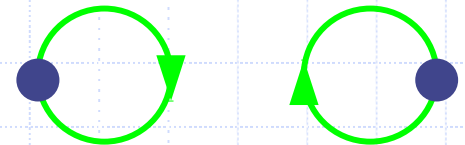
$$G(x, y) = - \langle Tr W^{-1}(x, y) W^{-1}(y, x) \rangle + 2 \langle (\sigma(x) - \langle \sigma \rangle) (\sigma(y) - \langle \sigma \rangle) \rangle$$

where

$$\sigma(x) \equiv Tr W^{-1}(x, x) = \bar{\psi}(x) \psi(x)$$



Disconnected diagram



- Vacuum contribution

# Lattice QCD simulations of

- ◆ There have been many Lattice Simulations of scalar without the disconnected diagram; “Valence Sigma”  $\sigma_V$
- ◆ deTar and Kogut
  - Phy.Rev. D36, (1987) 2828.
  - Screening masses
- ◆ Kim and Ohta
  - hep-lat/9609023, hep-lat/9712014
  - KS fermions,  $\beta=6.5$ ,  $48^3 \times 64$
  - $a=0.054\text{fm}$ ,  $48a=2.6\text{fm}$ ,

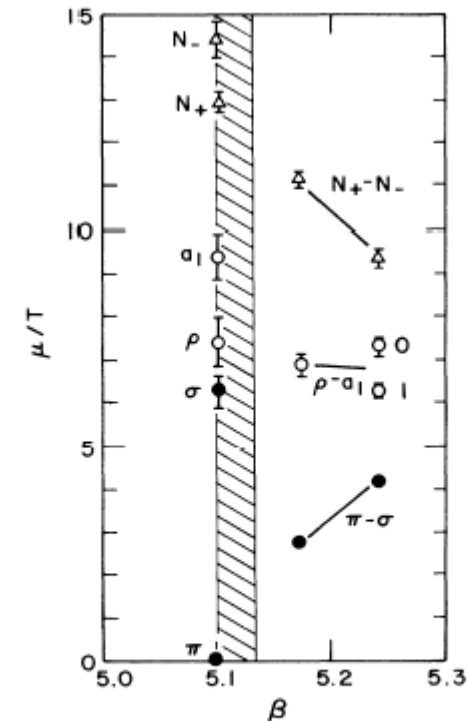
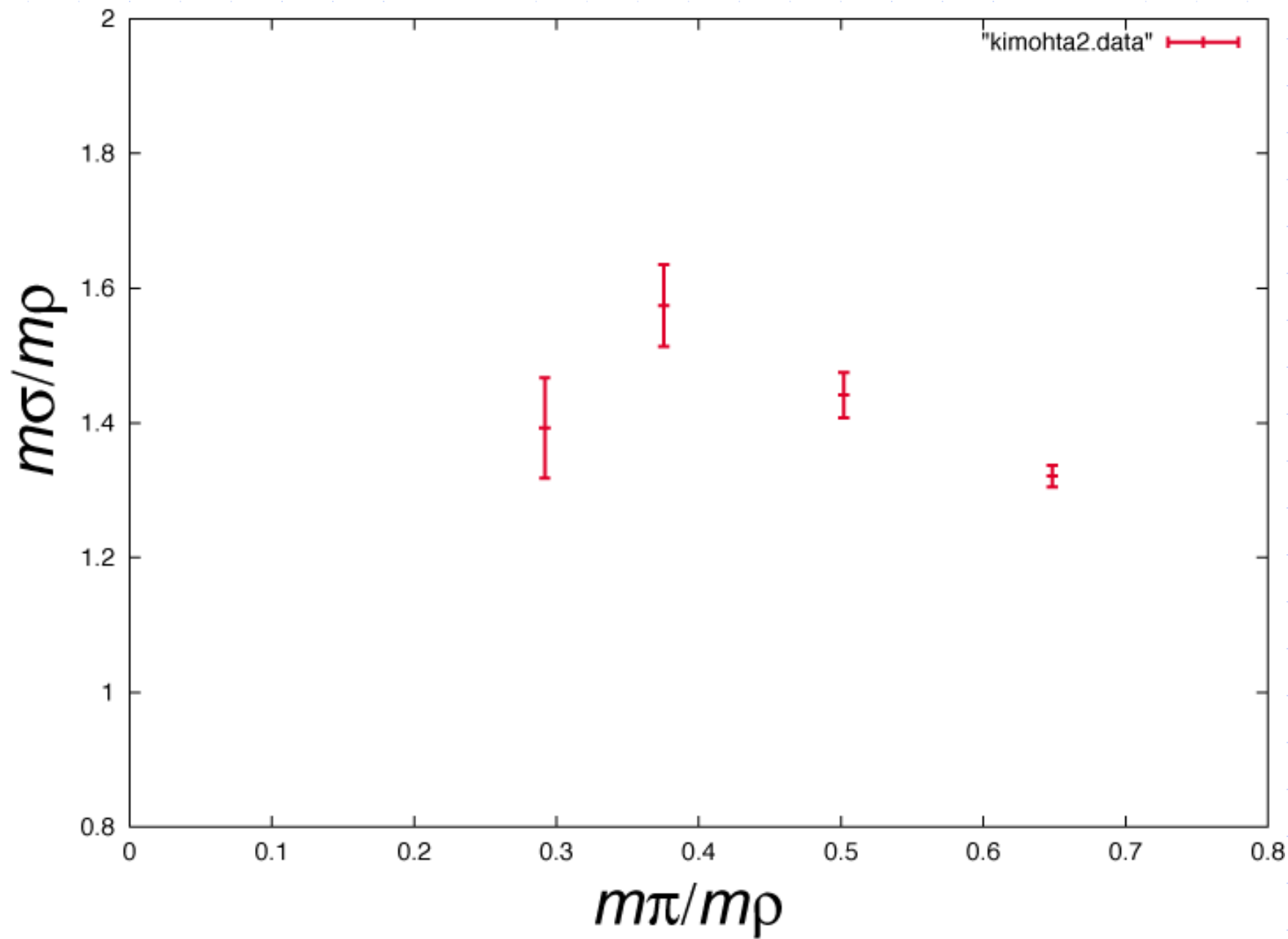


FIG. 10. Screening masses, expressed in units of the temperature, as extrapolated to the chiral limit, for the  $\pi$ -,  $\sigma$ -,  $\rho$ -, and  $\pi_1$ -meson plasmon modes, and the lowest even-parity ( $N_+$ ) and



# Lattice QCD simulations of (cont'd)

## ◆ W. Lee and D. Weingarten

- Phys. Rev. **D61** (1999) 014015
- Quench
- Mixing of Glue-ball and

## ◆ UKQCD C.McNeile and C.Michael

- Phys. Rev. **D63** (2001) 114503
- Full QCD

## ◆ Alford and Jaffe, Nucl.Phys. B578 (2000)367.

- Quench

- $\sigma = qq\bar{q}\bar{q}$        $E(q\bar{q}q\bar{q}) < E(q\bar{q} + q\bar{q})$

W. Lee and D. Weingarten,  
Phys. Rev. D **61** (1999) 014015

Mixing of  $q\bar{q}$  and glueball (  $I=0, J^{PC}=0^{++}$  )

Quenched approximation

Wilson fermion

Plaquette gauge action

$f_0(1710) \cdots$  lightest scalar glueball ( 73.8 (9.5)% )

$f_0(1500) \cdots s\bar{s}$  quarkonium ( 98.4 (1.4)% )

$f_0(1390) \cdots n\bar{n}$  quarkonium ( main )

n stands for  $u\bar{u} - d\bar{d} / \sqrt{2}$

# Lee and Weingarten (cont'd)

$$\begin{pmatrix} m_g & E(\mu_s) & \sqrt{2r}E(\mu_s) \\ E(\mu_s) & m_\sigma(\mu_s) & 0 \\ \sqrt{2r}E(\mu_s) & 0 & m_\sigma(\mu_n) \end{pmatrix} \rightarrow \begin{pmatrix} f_0(1710) & & \\ & f_0(1500) & \\ & & f_0(1390) \end{pmatrix}$$

Input:

$f_0(1710)$  1697(4)MeV  
 $f_0(1500)$  1505(9)MeV  
 $f_0(1390)$  1404(24)MeV  
 $m_\sigma(\mu_n)$  1470(25)MeV

$r \equiv E(\mu_n) / E(\mu_s)$  1.198(72)

(Only r is given by Lattice.)

Output:

$m_g$  1622(29) MeV  
 $m_\sigma(\mu_s)$  1514(11) MeV  
 $E(\mu_s)$  64(13) MeV

Lattice:

$m_g$  1654(47) MeV (World Average)  
 $m_\sigma(\mu_s)$  1322(42) MeV  
 $E(\mu_s)$  43(31) MeV



# C.McNeile and C.Michael (UKQCD), Phys. Rev. D63 (2001) 114503

## Mixing of the Iso-singlet scalar ( $I=0, J^{PC}=0^{++}$ ) and Glueball

Mass with Full QCD  
 $\ll$  mass with quench

$M < M$

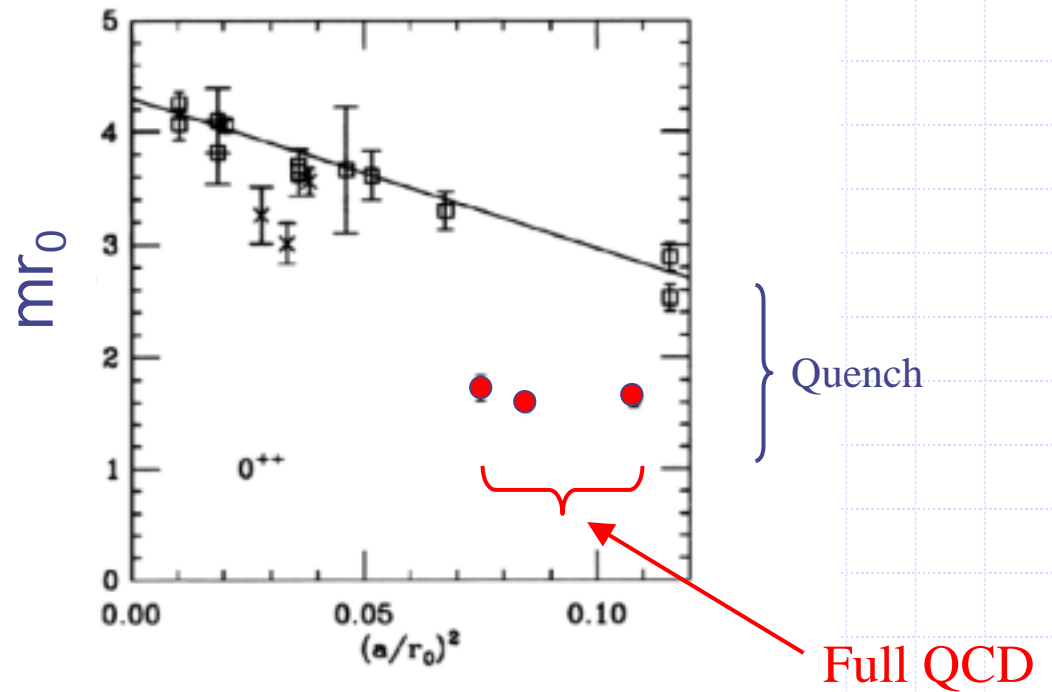


FIG. 5. The scalar mass versus  $a^2$ . The quenched results [12,10,13,14] are for the scalar glueball and are shown by boxes. The results from  $N_f=2$  flavors of sea quark are from glueballs [15] (crosses from SESAM) and the lightest flavor singlet scalar we find here (circles).

$r_0$ : Sommer factor  $r_0^{-1} = 394 \pm 20 \text{ MeV}$

# Lattice QCD simulations of - current going projects -

## ◆ Riken-Columbia-Brookhaven

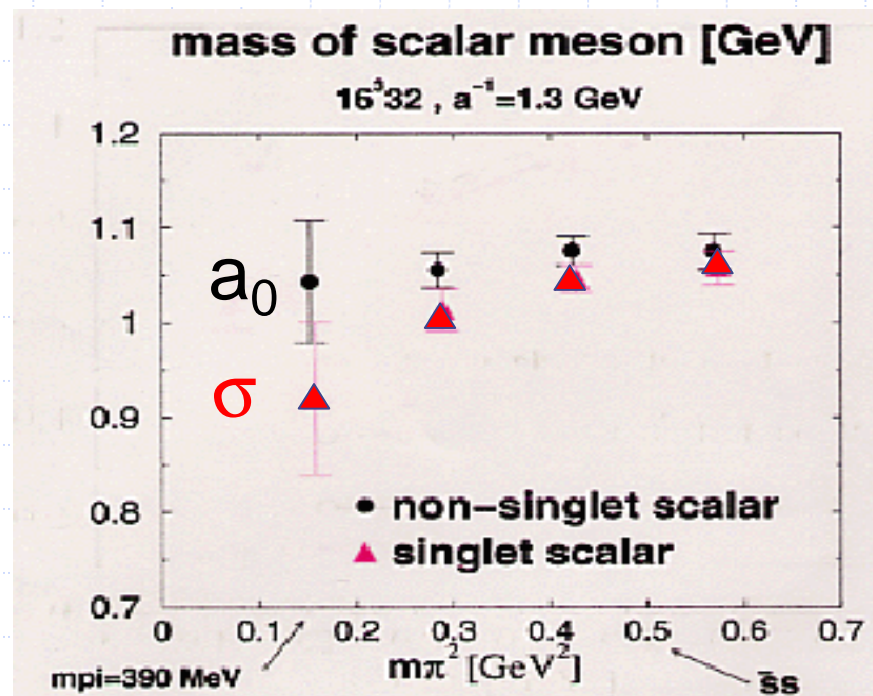
- Domain Wall Fermions
- Quench
- hep-lat/0209132 (Lattice02 Proceedings)

## ◆ Scalar Collaboration

- Wilson Fermions
- Full QCD  
■ hep-lat/0210012 (Lattice02 Proceedings)

# Riken-Brookhaven-Columbia

- ◆ Domain-wall fermions: Good Chiral nature
- ◆ Quench: Check the sickness of the quench calculations by quenched chiral perturbation theory.



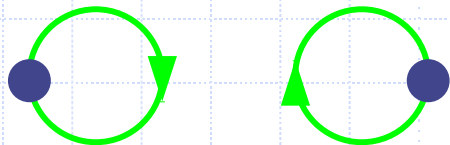
# Details of our Calculation (1)

Wilson Fermions (2 flavors)

Plaquette Gauge Action

**Full QCD** Update by Hybrid Monte Carlo (SX5 at RCNP)

Disconnected Part by  $Z_2$  Noise Method (SR8000 at KEK)



# Details of our Calculation (2)

## - Simulation parameters

Lattice size :  $8^3 \times 16$

= 4.8

= 0.1846, 0.1874, 0.1891

well established by CP-PACS,

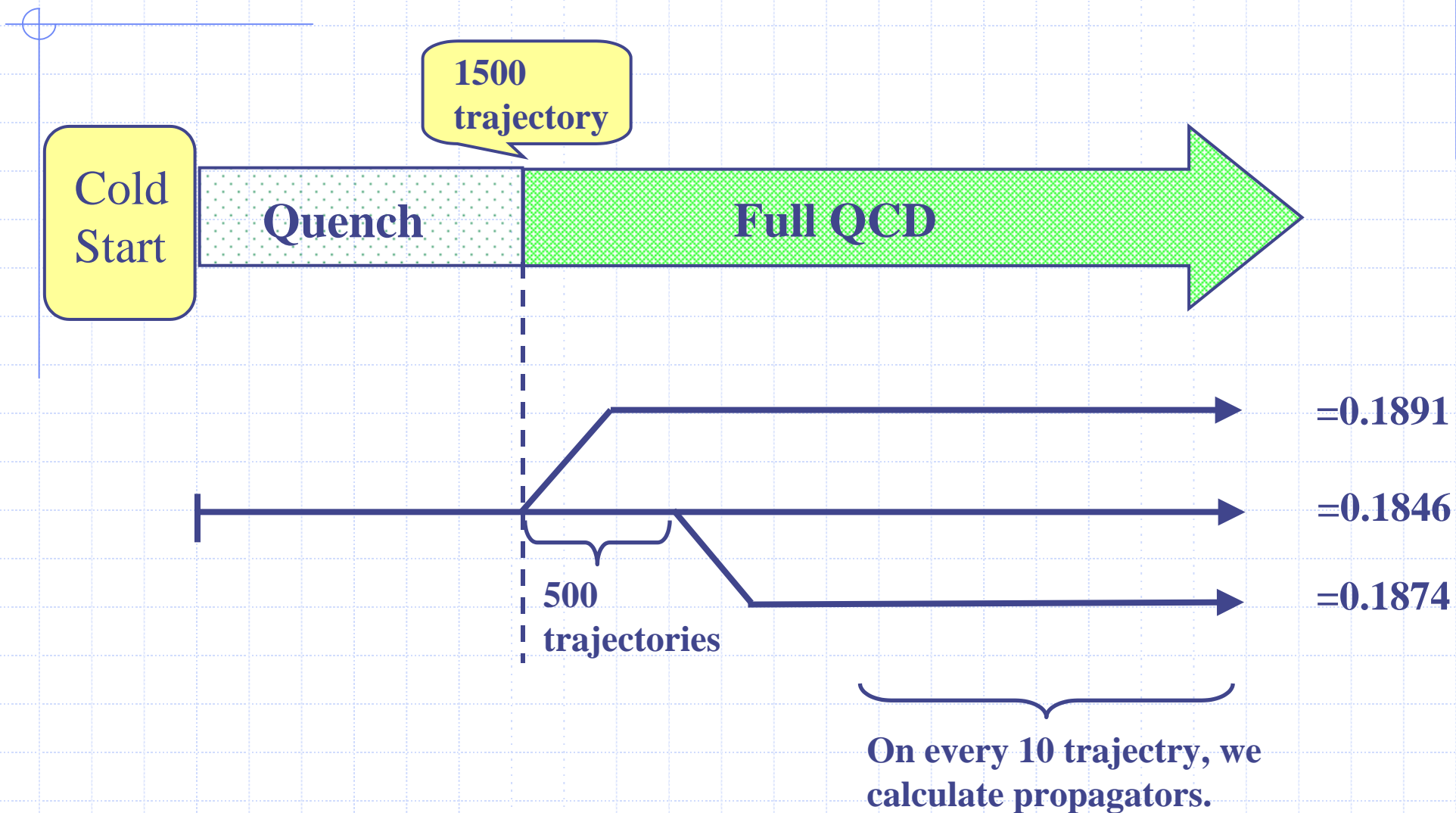
$a = 0.197(2) \text{ fm}$ ,  $c = 0.19286(14)$

( CP - PACS, Phys. Rev. D60(1999)114508 )

Wilson Fermions & Plaquette gauge action

Number of the Z2 noise = 1000 , 500

# Details of our Calculation (3)



# Details of our Calculation (4)

[ = 0.1846]

1470 configurations from 720th trajectory

[ = 0.1874]

970 configurations from 710th trajectory

[ = 0.1891]

400 configurations from 500th trajectory

Separation between configurations  
are 10 trajectories

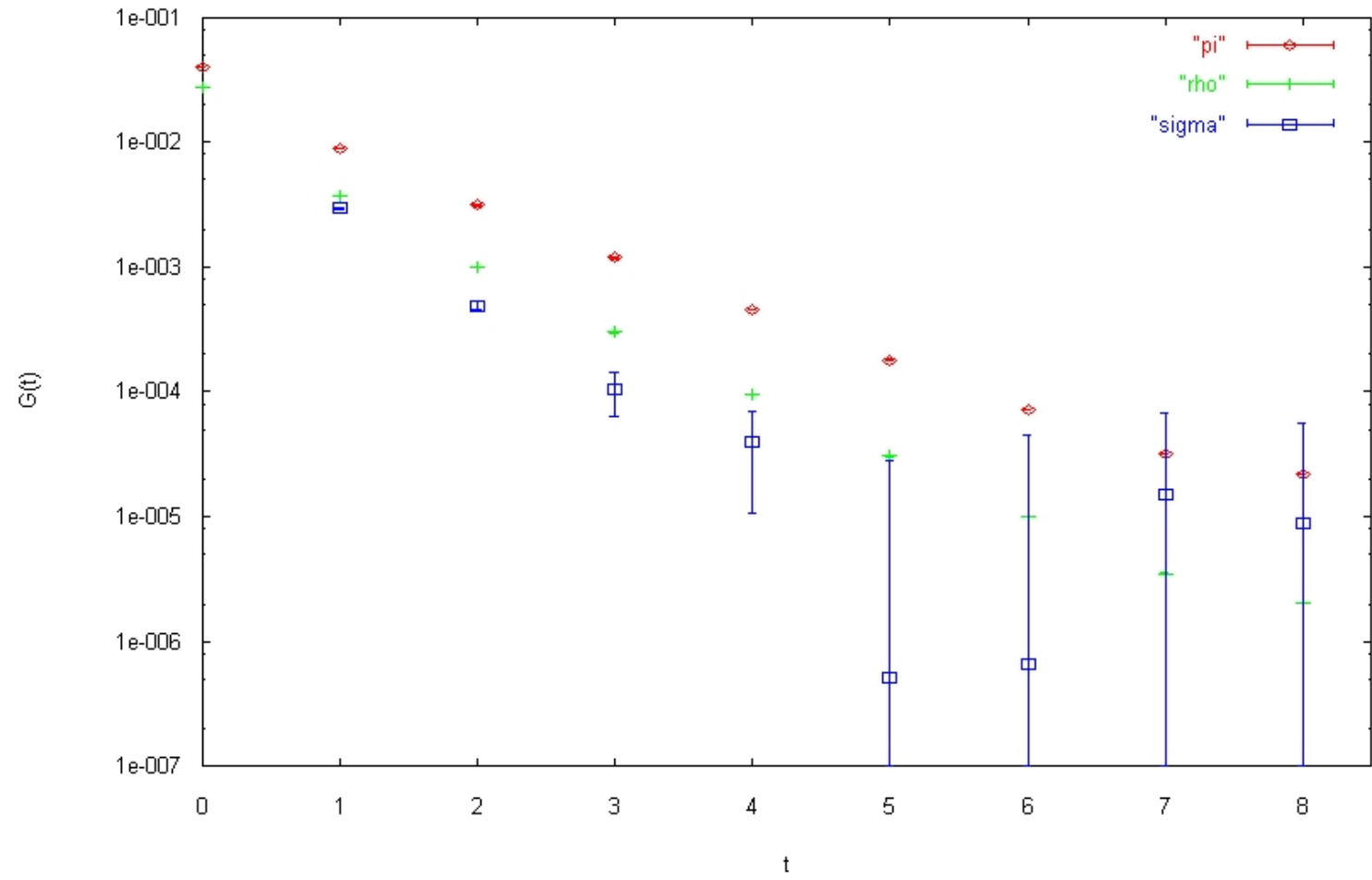
## Details of our Calculation (5)

	$m_{\pi} / m_{\rho}$ (Our Results)	$m_{\pi} / m_{\rho}$ (CP-PACS)
0.1846	$0.825 \pm 0.001$	$0.8291 \pm 0.0012$
0.1874	$0.760 \pm 0.002$	$0.7715 \pm 0.0017$
0.1891	$0.692 \pm 0.005$	$0.7026 \pm 0.0032$



# mesons ( $\kappa = 0.1846$ )

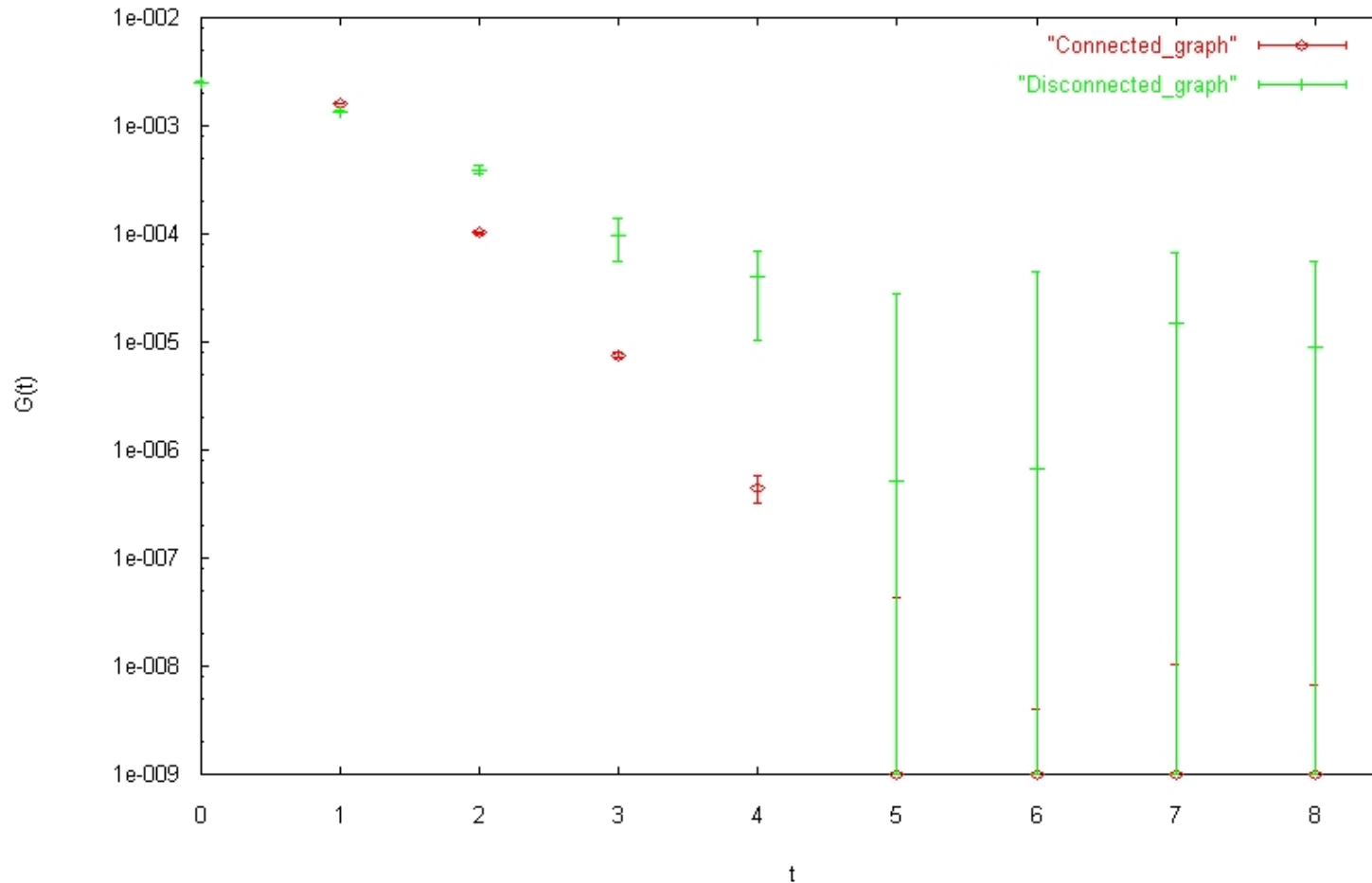
$\kappa = 0.1846$  ( 1470 configurations )



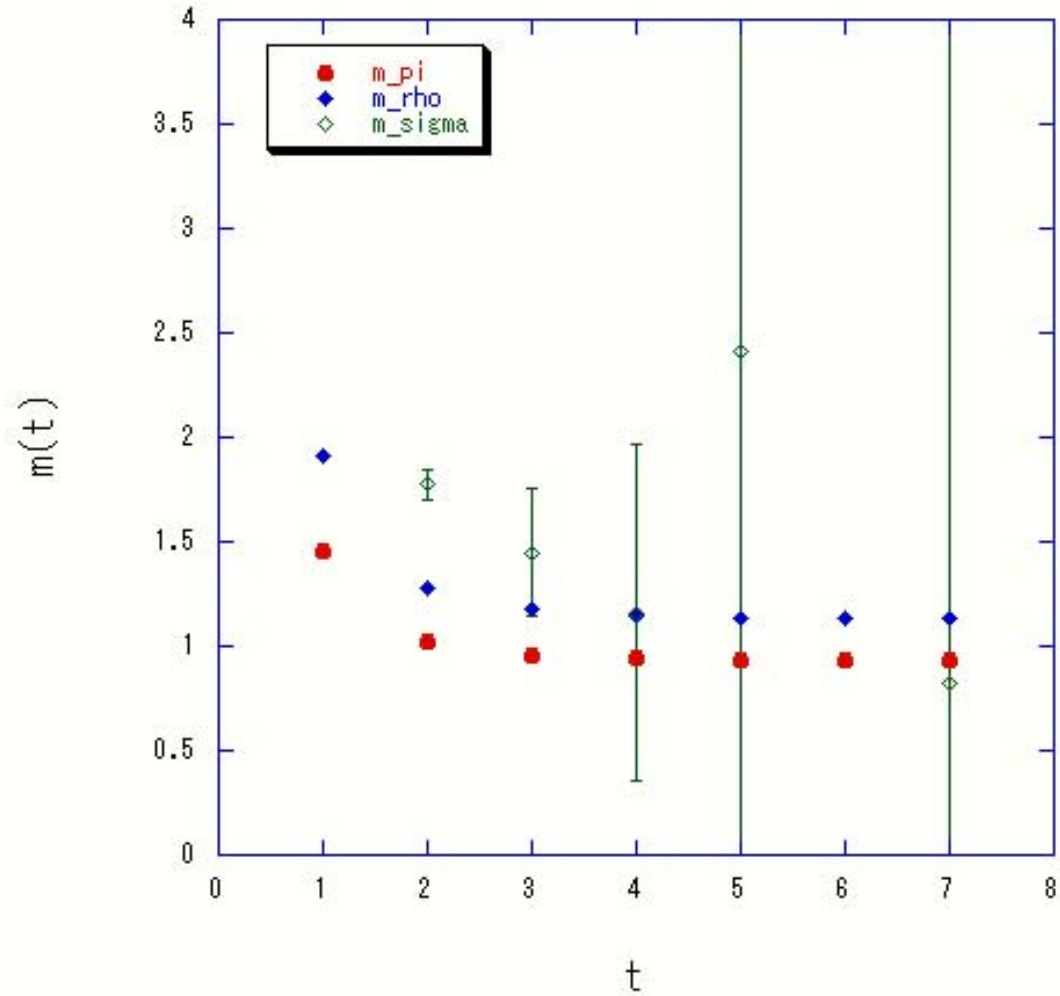
# meson propagators

## Connected and Disconnected Parts ( $\kappa = 0.1846$ )

$\kappa = 0.1846$  ( 1470 configurations )

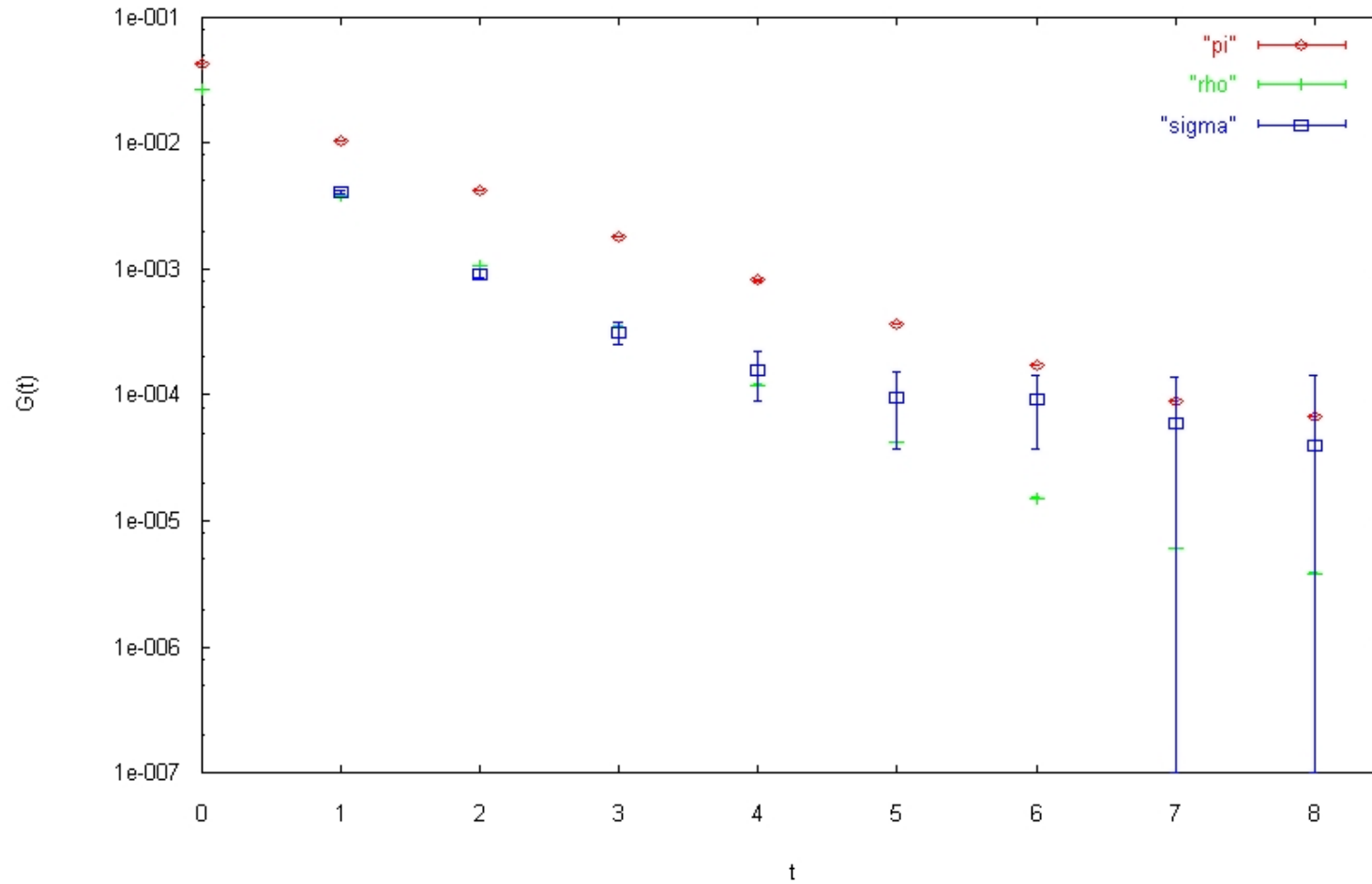


# Effective mass ( $\epsilon = 0.1846$ )



# mesons ( $\kappa = 0.1874$ )

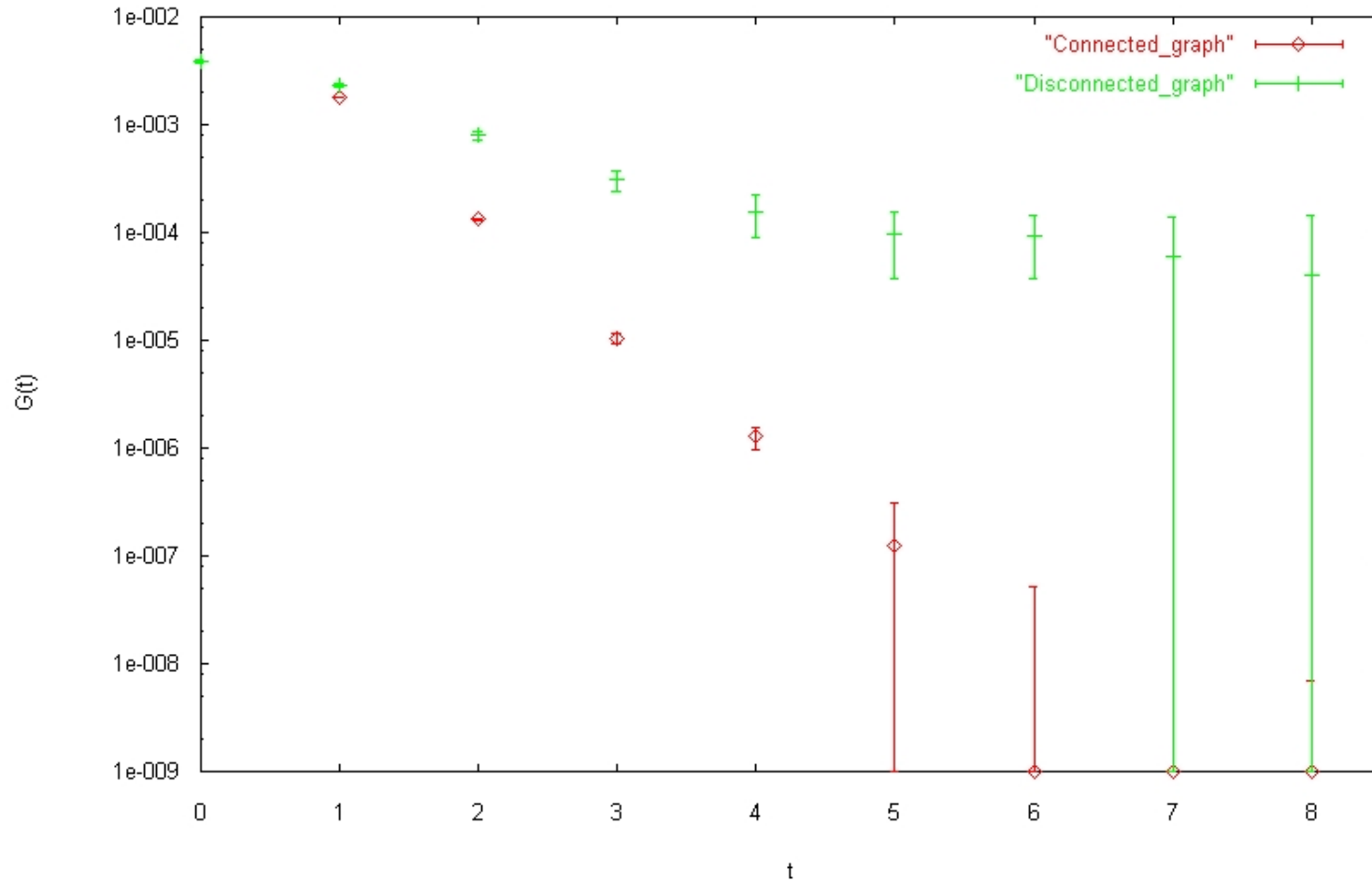
$\kappa = 0.1874$  (970 configurations)



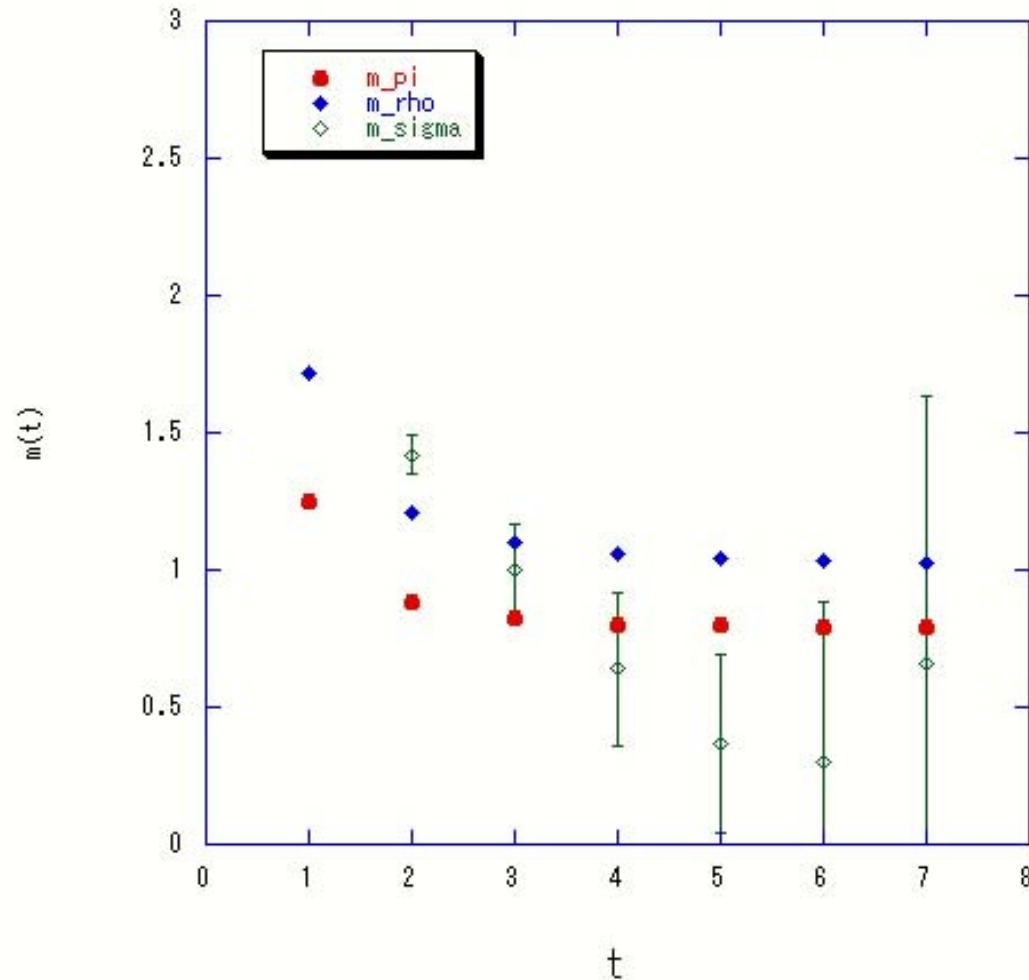
# meson propagators

## Connected and Disconnected Parts ( $\kappa = 0.1874$ )

$\kappa = 0.1874$  ( 970 configurations )

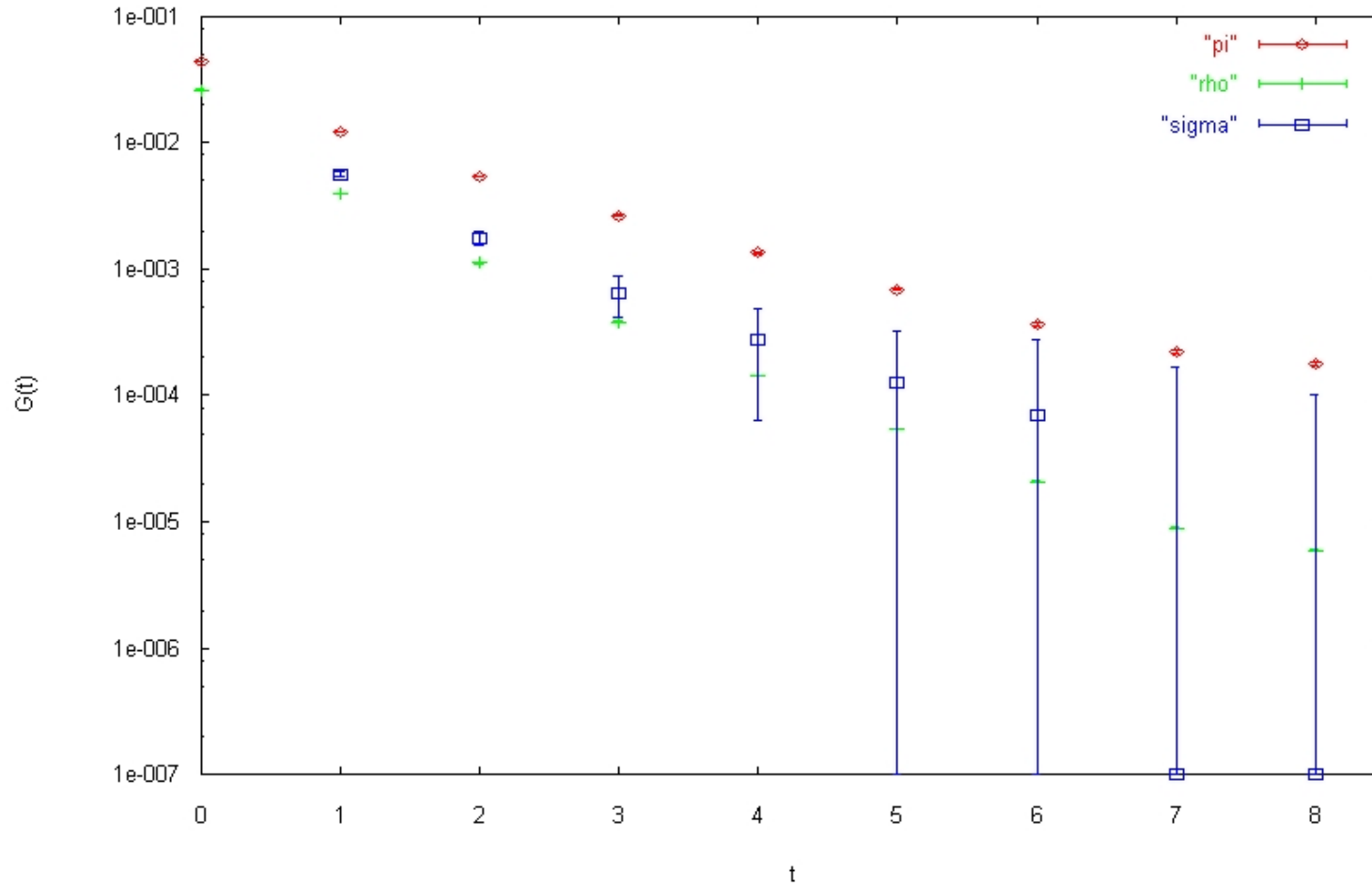


# Effective mass ( $=0.1874$ )



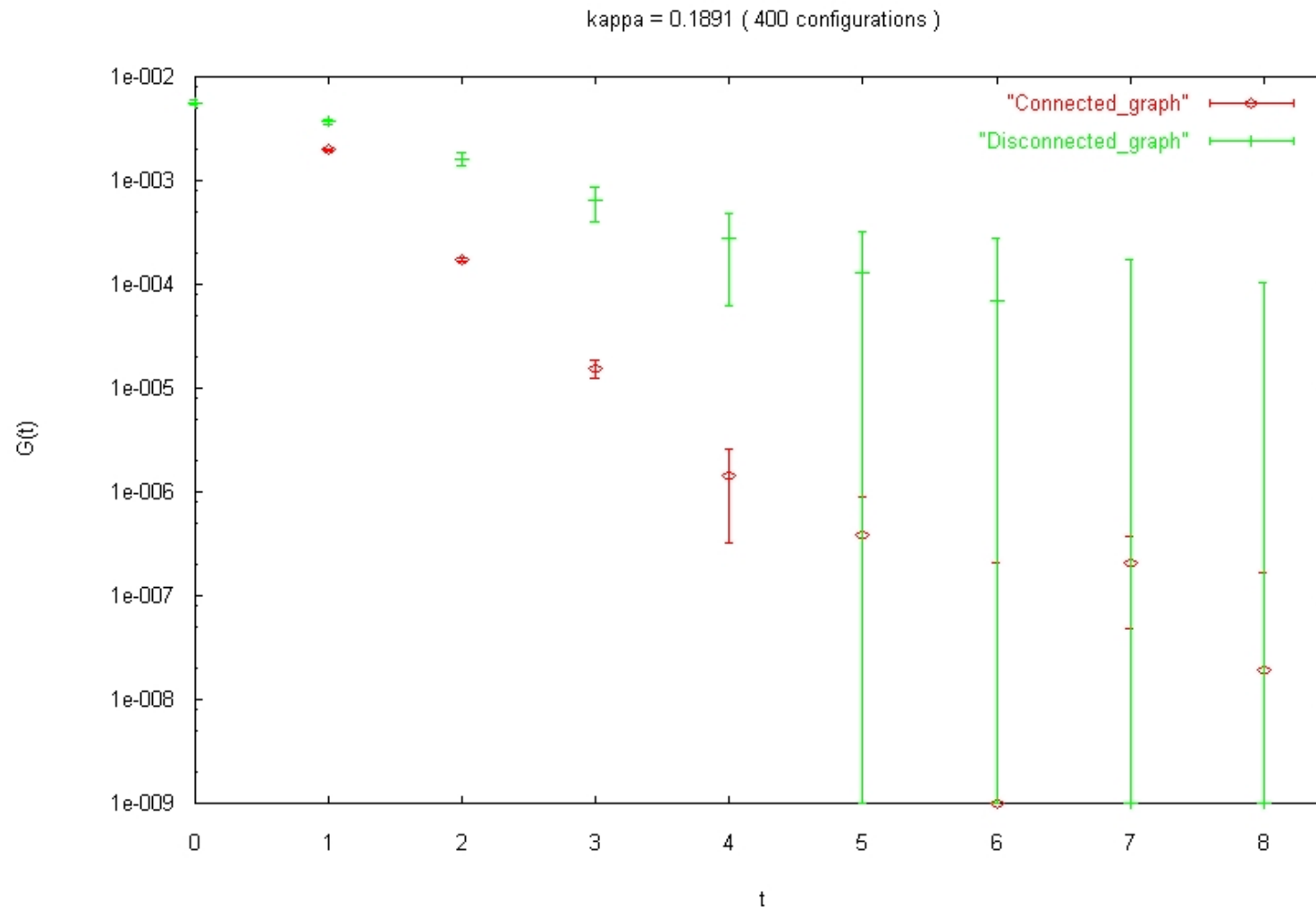
# mesons ( $\kappa = 0.1891$ )

$\kappa = 0.1891$  ( 400 configurations )



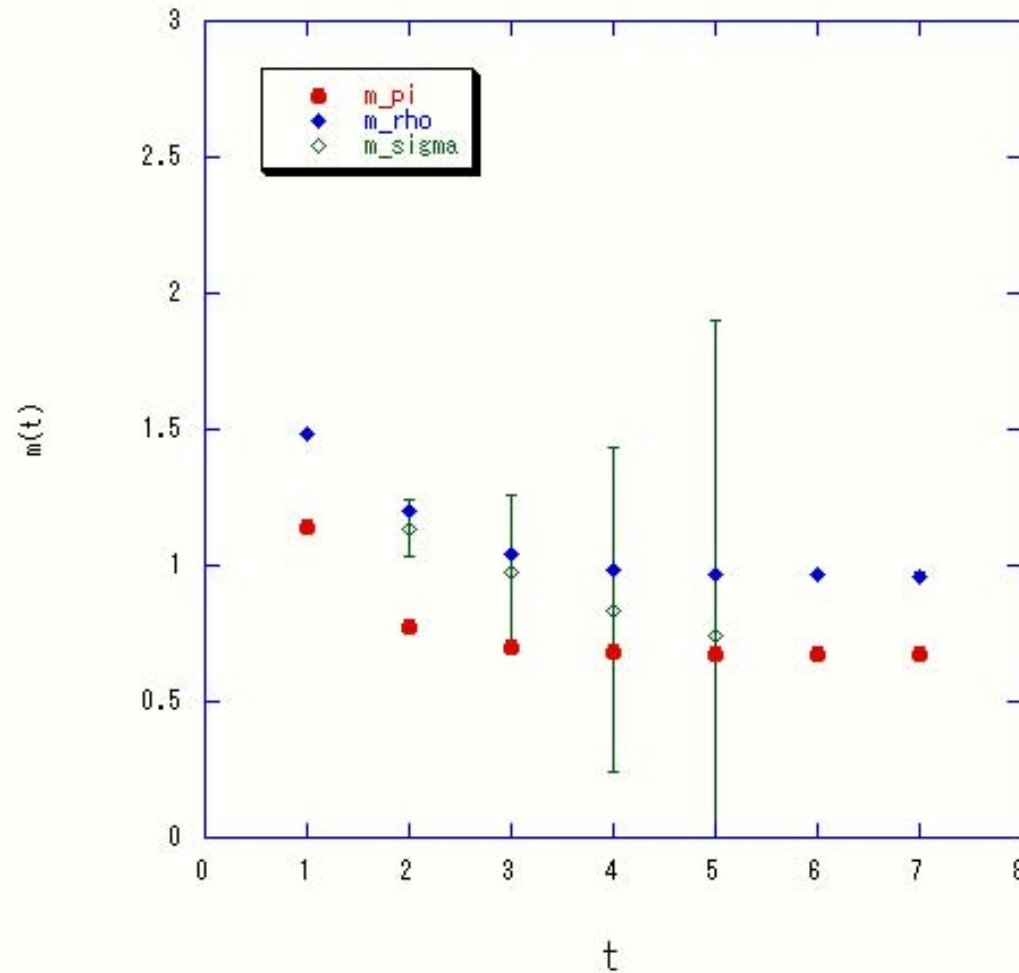
# meson propagators

## Connected and Disconnected Parts ( $\kappa = 0.1891$ )





# Effective mass ( $\approx 0.1891$ )



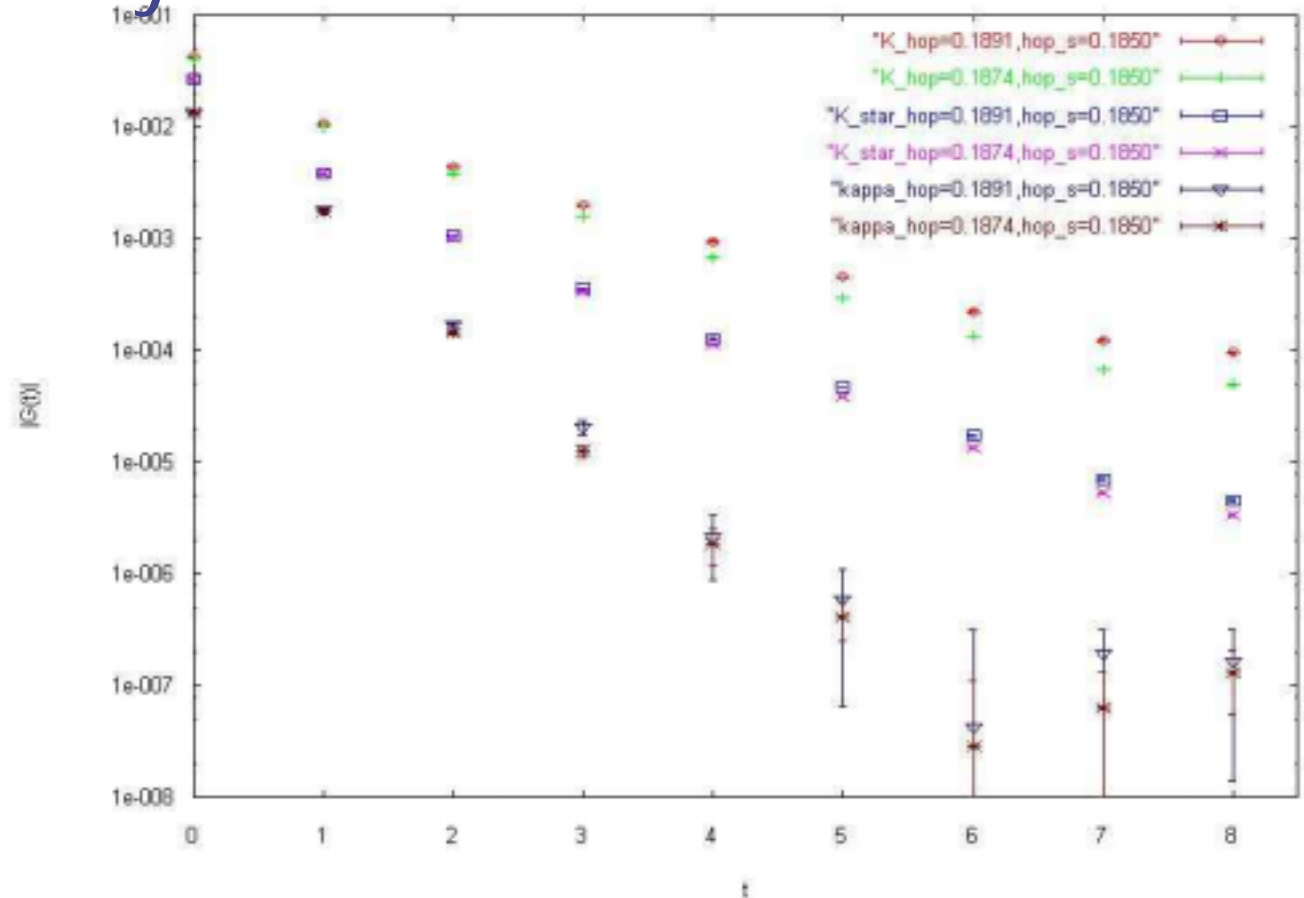
$$m_{\sigma} \approx m_{\rho}$$

# $\kappa$ meson

◆ strange scalar

$$\kappa^+ = \bar{\psi}_s \psi_u$$

◆ Connected only



# Summary

- ◆ Although  $\sigma$  propagators are noisy and we need high statistics, present data suggest that  $\sigma$  appears as a pole of QCD.
- ◆ Disconnected diagram dominates at large  $t$ .
- ◆  $m_\sigma \approx m_\rho$
- ◆ Analysis of  $\kappa$  will come soon.
- ◆ Lattice QCD Study of Scalar mesons was crazy before, but is now recognized as a meaningful work. In one year, it becomes matured and produces reliable data on the scalar mesons.